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PREDICTION OF NORMAL FORCE, PITCHING MOMENT, AND YAWING FORCE ON BODIES OF REVOLUTION AT ANGLES OF ATTACK UP TO 50 DEGREES USING A CONCENTRATED VORTEX FLOW-FIELD MODEL

M < 0.9

BY

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24 OCTOBER 1973

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less than 0.9. For models with smaller nose fineness ratios, and at higher Mach numbers, accuracy deteriorates and side force predictions are too high.

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By direction

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INTRODUCTION

As a body of revolution is exposed to an increasing angle of attack, the flow on its leeward side separates and rolls up to form vortices. Under these conditions, normal force and pitching moment can no longer be predicted by slender body theory. For relatively small angles of attack ($\alpha < \tan^{-1}(4.5/l)$) two symmetric vortices form as shown in Figure 1a. At larger angles of attack wake vortices grow unsymmetrically and break away or are shed from the leeward side, as illustrated in Figure 1b. The resulting flow-field asymmetries induce a yawing force and moment.

The analysis of bodies of revolution at angles of attack has been approached by many workers with varying degrees of rigor. Simple engineering calculations can be used to determine normal forces and pitching moments (Refs. (1)-(3)), but extensive numerical computations must be employed if detailed flow-field information is desired (Refs. (4), (5)). The objective of this report is to predict the yawing force component. Since this force arises from asymmetries in the vortex wake, the appropriate level of analysis must be sufficiently rigorous to provide an approximate wake description. Engineering level calculations do not generally provide this description and, therefore, are not amenable to these ends. While detail flow-field calculations can, in principle, achieve this purpose, they have only been applied up to very moderate angles of attack (10-15 degrees), too low for yawing forces to occur. Semi-empirical methods, such as those by Bryson (Ref. (6)), Schindel (Ref. (7)), and Angelucci (Ref. (8)), provide an approximate wake description, but assume a symmetric vortex wake and are not capable of predicting yawing forces. Also, normal forces and pitching moments can only be calculated up to moderate angles of attack (20-30 degrees).

The theory developed in this paper extends the model originally proposed by Bryson (Ref. (6)) by removing the symmetric wake assumption. This allows the side forces to be predicted, and increases angle of attack for which the theory is valid to about 50 degrees. In the current theory, the flow field is approximated using the impulsive flow analogy. This analogy assumes that the crossflow field is swept down the length of the body at the rate $U_{c\infty}\alpha$. At each axial station the flow field is taken to be analogous to flow about a cylinder in crossflow, whose radius is equal to the body radius at that axial location. The developing crossflow field is thus similar to flow about an impulsively started cylinder in crossflow. The development of the viscous flow field is modeled by superimposing point vortices on the potential flow solution for flow about a cylinder. Analysis is limited to ogive cylinders with sharp or blunt nose tips.

Using experimental data for a body with a laminar boundary layer ($2.5 < Re/ft \times 10^{-6} < 5.6$) and covering a range in Mach number from 0.5 to 1.1, free parameters inherent in the model are determined.

Calculated force and moment values are compared to experiment and computed vortex paths are compared to those determined from schlieren photographs. This type of approach allows conclusions to be drawn concerning the validity of the current problem formulation.

SYMBOLS

C_{DC}	local drag coefficient, $\frac{\pi}{2} \frac{r}{\sin^2 \alpha} \frac{dC_N}{dx}$
C_M	pitching moment coefficient, $\frac{M_y}{Dq_\infty \pi r^2}$, referenced to model nose
C_M'	yawing moment coefficient, $\frac{M_z}{Dq_\infty \pi r^2}$
C_N	normal force coefficient, $\frac{F_N}{q_\infty \pi r^2}$
C_Y	yawing force coefficient, $\frac{F_y}{q_\infty \pi r^2}$
D	diameter of body of revolution
F_N	normal force
F_Y	yawing force
k	number of vortices in wake
L	body length
L_N	nose length
λ	body fineness ratio
λ_n	nose fineness ratio
M	Mach number
M_C	crossflow Mach number, $M \sin \alpha$
M_y	pitching moment
M_z	yawing moment
n	frequency for shedding vortices of like sign

q_∞	free-stream dynamic pressure
R_n	nose tip radius of curvature
r	local body radius
r_0	maximum body radius
Δr	difference in the radial location of the two initial vortices at which the first vortex will be shed.
r_n	the distance from the separation point, measured along the feed sheet at which nascent vortices are introduced
Re	Reynolds number
Δr_I	perturbation in the radial location of the first two vortices
S	Strouhal number
t	time
U	free-stream velocity
$W(\zeta)$	crossflow plane velocity at ζ , $\frac{d\bar{\phi}(\zeta)}{d\zeta}$
w	dimensionless crossflow plane velocity, $W/Usin\alpha$
w_j	dimensionless crossflow plane velocity at the location of the j th vortex. The influence of the j th vortex is neglected.
x	distance along missile body (see Fig. 1)
Δx	x distance between points where successive vortices are shed.
y, z	crossflow plane coordinates (see Fig. 2)
y_j, z_j	y and z coordinates of the j th vortex
α	angle of attack
Γ	vortex circulation
δ	rate of radius change as a function of x , $\frac{dr}{dx}$
ζ	complex crossflow plane coordinate, $y + iz$
ζ_{01}, ζ_{02}	crossflow plane separation points (see Fig. 2)

θ_{01}, θ_{02}	crossflow plane separation angles (see Fig. 2)
λ	dimensionless circulation, $\frac{\Gamma}{2\pi r \sin \alpha U}$
λ_n	dimensionless circulation of newly introduced vortice
$\Delta\lambda_I$	initial vortex strength
ϕ	crossflow velocity potential
ϕ_N	nose roll angle
χ	dimensionless crossflow velocity potential, $\phi/Us \sin \alpha$

LITERATURE SURVEY

The complete analysis of a body of revolution at angle of attack is a three-dimensional problem involving the solutions to the Navier-Stokes equations. This problem may be approached numerically in its entirety requiring an extensive computational effort. For relatively small angles of attack (12 degrees), numerical solutions have been obtained for cones using an approximate set of Navier-Stokes equations (Ref. (4)).

The device most commonly used to simplify the problem formulation is the impulsive flow analogy. This reduces the complete problem to a two-dimensional time varying analysis of a viscous fluid. By this analogy, the crossflow field is equivalent to that about an impulsively started, expanding cylinder, where $t = x/U \cos \alpha$. The expanding cylinder represents the increase in cross-sectional radius which occurs along the length of the body. Often the additional assumption is made that the flow about an impulsively started cylinder is incompressible. Throughout the remainder of this report the term impulsive flow analogy will be assumed to imply this assumption.

On the most rigorous level, the impulsive flow analogy allows the crossflow drag to be deduced from a numerical solution to the Navier-Stokes equations for flow over an impulsively started cylinder in crossflow. Several examples of this type of procedure are available for cylinders of constant radius (Refs. (9), (10)). For direct application to a body at angle of attack Walitt and Trulio (Ref. (5)) have carried out this analysis with a cylinder of varying radius in compressible flow. However, the computational requirements are prohibitive and solutions have only been obtained for angles of attack of 15 degrees.

A less rigorous analysis of impulsively started flow over a cylinder involves superposition of a large number of point vortices on the potential solution for incompressible flow about a cylinder. Typically, several hundred discrete vortices may be used. After formation, each point vortex is assumed to drift with the local fluid velocity. The schemes for predicting rate of vorticity

formation differ extensively. Although many such studies have been carried out, most are restricted to bodies of fixed cross-sectional dimensions (Refs. (11), (12) and (13)). However, Angelucci (Ref. (8)) has made a direct application to bodies at angle of attack, assuming a symmetric vortex wake. Good agreement is obtained for normal forces and moments.

A somewhat simpler method uses a single vortex to describe each wake vortex. In a procedure originally proposed by Bryson (Ref. (6)), vortex motion is determined by a balance of forces on a vortex and its feeding sheet. The model was originally limited to circular bodies with a wake pattern containing two symmetric vortices. This approach has been extended by Schindel (Ref. (7)) to bodies of elliptic cross section using conformal mapping techniques. In addition, changes in empirical parameters allow laminar and turbulent boundary effects to be accounted for. Results compare well with experimental data for normal forces. However, agreement of the pitching-moment calculation with experimental data is not as good.

Range of validity for the results is $\alpha < \tan^{-1}(5/\ell)$. Davis (Ref. (14)) has extended Bryson's methods to a cylinder in crossflow with an asymmetric wake. But this formulation is valid only for cylinders of fixed radius and, therefore, cannot be directly applied to a body at angle of attack.

The most approximate theories available consider forces on a body at angle of attack to be made up of a viscous and an inviscid contribution which are additive. Initial attempts along these lines were by H. J. Allen (Ref. (1)). The inviscid force expression applied is similar to that obtained from slender body theory. The viscous contribution at each axial station is computed by applying the steady-state drag force for a cylinder in crossflow. This report was followed by that of Kelly (Ref. (2)) who attempted to improve the method by using drag data of an impulsively started cylinder in crossflow. His results were somewhat better than those of Allen, giving a better prediction of the pitching moment. Kelly's method has been refined by others, most recently by Thomson (Ref. (15)) to include angles of attack up to 90 degrees. In this work advantage is taken of impulsive flow data obtained by Sarpkaya (Ref. (16)). Correction tables allow Mach number, end effects, nose shapes, and the transition from laminar to turbulent flow to be taken into consideration. In addition, an experimental wake description is developed which allows side forces to be estimated. In a slightly different vein, L. H. Jorgensen (Ref. (3)) in an extension of Allen's work, presents a method of estimating normal forces on bodies at an angle of attack varying from 0 to 180 degrees and Mach numbers varying from subsonic to hypersonic.

BRYSON'S METHOD FOR IMPULSIVELY STARTED CYLINDERS
IN INCOMPRESSIBLE FLOW

As pointed out above, our analysis is essentially that of an impulsively started cylinder in incompressible, viscous flow. Such flow can be simulated by superimposing point vortices on the potential solution for flow about a cylinder. In the present approach, a single point vortex is used to simulate each wake vortex. To satisfy the boundary conditions on a cylinder, each vortex in the wake is required to have a corresponding image vortex. The resulting configuration is shown in Figure 2 and the complex velocity potential is given as follows:

$$\begin{aligned} \phi(\zeta) = & U \left[-\left(\zeta - \frac{r^2}{\zeta}\right) i - \frac{i}{2\pi} \sum_{j=1}^k \Gamma_j \ln \left\{ \frac{(\zeta - \zeta_j)}{(\zeta - r^2/\zeta_j)} \right\} \right. \\ & \left. + \frac{r}{\tan \alpha} \frac{dr}{d\zeta} \ln(\zeta) \right] \end{aligned} \quad (1)$$

The term, $-(\zeta - \frac{r^2}{\zeta})i U$, is the velocity potential solution for flow over a cylinder. Each of the logarithmic terms within the summation represent the velocity potential for a vortex and its image within the cylinder. The remaining term, which is the expression for a source, takes into account the changing cylinder radius.

The equations governing the motion of the vortices remain to be determined. Each logarithmic term within the summation in Equation (1) is a double valued function and has a branch line which connects the described vortex to its image. The portion of the branch line of physical importance is the part which crosses the flow field. It is considered to extend from the vortex location to the nearer of the points, ζ_{01} and ζ_{02} , which are on the cylinder surface (see Fig. 2).

These two points correspond to the position where flow is assumed to separate from the cylinder surface. The branch line of a growing vortex ($d\Gamma/dt \neq 0$) is known as a feeding sheet. Across it, a pressure jump of $\frac{\rho(d\Gamma/dt)}{\text{unit length}}$ occurs, which can be determined from the momentum conservation equation. The total force on the feeding sheet is $i\rho(\zeta_k - \zeta_{0k})\frac{d\Gamma}{dt}$. This force on the feeding sheet is balanced by a lift force on the vortex of $-i\rho\Gamma(\frac{d\zeta_k}{dt} - w_k)$ given by the Kutta-Joukowski theorem. Here $(\frac{d\zeta_k}{dt} - w_k)$ is the vortex velocity with

respect to the local fluid velocity. Thus, the motion of each vortex is governed by the equation

$$\frac{d\zeta_j}{dt} + \frac{(\zeta_j - \zeta_{0j})}{r_j} \frac{d\Gamma_j}{dt} = w_j \quad (2)$$

The vortex circulations are determined by applying the Kutta conditions at the crossflow separation points:

$$W(\zeta_{01}) = W(\zeta_{02}) = 0 \quad (3)$$

This allows the circulation of two different vortices at a time to be specified. This does not preclude the existence of other vortices in the flow whose circulations are known.

Left undetermined by the Bryson method are the crossflow separation points, ζ_{01} and ζ_{02} , and boundary conditions in t . The boundary conditions in t include the points where vortices are introduced and shed. Also, the location of vortices in the crossflow plane and their strengths must be specified when they are introduced.

In applying the Bryson method the boundary conditions in t must be specified in such a manner that the resulting wake conforms to observations. When the wake behind the cylinder is assumed to become asymmetric, the following type of procedure appears to be required. Initially, two vortices are introduced. Using the governing equations the growth and the motion of these vortices are traced numerically. At some point, to be empirically determined, the circulation of one of these vortices is fixed and this vortex is considered to be shed. A new vortex is then introduced on the same side of the cylinder that the shed vortex originally emanated from. This shedding process is repeated at later points in time, vortices being shed from alternate sides of the cylinder.

APPLICATION OF THE BRYSON METHOD TO A MISSILE AT ANGLE OF ATTACK

The impulsive flow analogy implies the relations $x = U \cos \alpha t$ and $W(\infty) = U \sin \alpha t$. Using these expressions and the definitions $\Gamma = 2\pi U r \sin \alpha \lambda$ and $w_j = w_j U \sin \alpha t$, Equation (2) becomes

$$\frac{d\zeta_j}{dx} + \frac{(\zeta_j - \zeta_{0j})}{\lambda_j} \frac{d\lambda_j}{dx} = \tan \alpha w_j - \frac{(\zeta_j - \zeta_{0j})}{r} \frac{dr}{dx}; \quad j=1,2 \quad \text{growing vortices}$$

$$\frac{d\zeta_j}{dx} = \tan \alpha w_j; \quad j=3,4, \dots, K \quad \text{shed vortices}$$

The Kutta condition for a cylinder whose radius may be expanding becomes:

$$\frac{1}{U_{\infty} \tan \alpha} \frac{d\bar{\phi}(\bar{z}_{01})}{d\bar{z}} = \omega(\bar{z}_{01}) = \frac{r}{\tan \alpha} \frac{dr}{dz} \frac{1}{\bar{z}_{01}} \quad (5)$$

$$\frac{1}{U_{\infty} \tan \alpha} \frac{d\bar{\phi}(\bar{z}_{02})}{d\bar{z}} = \omega(\bar{z}_{02}) = \frac{r}{\tan \alpha} \frac{dr}{dz} \frac{1}{\bar{z}_{02}}$$

Substituting Equation (1) into Equations (5), it becomes evident that λ_1 and λ_2 are determined by solving two simultaneous equations. It is also evident in this formulation that satisfaction of the inviscid boundary condition on the cylinder is sufficient to ensure that both λ_1 and λ_2 are real.

In solving Equations (4) subject to conditions (5), λ_1 and λ_2 have been eliminated from (4) reducing the number of equations involved. The actual derivation is shown in Appendix A.

The formulated method is applied in the following manner. Near the nose of the missile two nascent vortices are introduced. The x coordinate of introduction is determined by satisfying the relation $\frac{2}{\lambda} \sin \theta \tan \alpha = \tan \delta$. This relationship moves the point of vortex introduction toward the nose of the model with increasing angle of attack. In cases where this relation places the starting point in front of the nose, the problem is automatically started at the nose. This relationship was derived by Bryson (Ref. (6)) by considering a cone. The slope δ represents the half angle of the largest cone for which vortex separation occurs. The two vortices are each assumed to be located along the directions $\theta_{01} + \pi/3$ and $\theta_{02} - \pi/3$ from their respective separation points δ_{01} and δ_{02} . This criterion is also due to Bryson and was determined by the analysis of the trajectory of a vortex near the crossflow separation point.

Asymmetry is introduced at this point by slightly perturbing the radial location* of the two vortices. The initial strengths of the two vortices, as well as the magnitude of the perturbation applied, are free parameters which are empirically determined. They will be discussed in more detail later.

As the numerical solution to the governing equations progresses, the asymmetry in the problem grows. When the difference in radial location of the two vortices reaches a specified value, the outermost vortex is shed and a new vortex is introduced on the same

*the radial location of the j th vortex is defined to be

$$\sqrt{y_j^2 + z_j^2}$$

side of the body. In order to introduce a new vortex whose circulation is other than zero, the existing vortex pattern must be perturbed. This is accomplished by reducing the circulation of the shed vortex a few percent. It can be shown that this procedure will always lead to a new vortex with a circulation of the correct sign. The new vortex is assumed to be located along the feeding sheet of the vortex to be shed. Its initial strength is uniquely determined by its crossflow field location and the decrease in circulation ascribed to the shed vortex.

Subsequent vortices are shed at intervals specified by the Strouhal frequency. Initial conditions for succeeding vortices are identical to those just described.

Application of the method described above requires specification of the following free constants:

1. Initial vortex strength (λ_I)
2. Initial radial perturbation (Δr_p)
3. The difference in radial location of the two initial vortices at which the first vortex is shed (Δr)
4. The distance from the separation point, measured along the feeding sheet of the shed vortex, at which a nascent vortex is introduced (r_n)
5. Circulation reduction of the shed vortex ($\Delta \lambda_n$)
6. Crossflow separation angles (θ_{01} and θ_{02})
7. Strouhal number (S)

The Strouhal number is determined experimentally using schlieren photographs of a tangent ogive cylinder at angle of attack. Although measured values showed a large degree of fluctuation, an average value of .17 was established.

The crossflow separation angle can also be determined in the wind tunnel, however, use of measured values leads to unstable numerical results. This may be due to the approximate nature of the flow field model. Simulation of a wake vortex using a single point vortex may not be giving an accurate flow field description in the vicinity of the separation points. To circumvent this problem, the lead of Bryson is followed and the separation angle is set near 45 degrees.

There is no direct way of readily determining the remaining parameters. However, it should be pointed out that the numerical solution is much less sensitive to these remaining parameters. Their actual values were established by a comparison with experimental measurements of normal forces and schlieren pictures of the leeward flow field.

NORMAL AND YAWING FORCES

The forces acting on any section of a body can be determined by enclosing it in a cylindrical control volume and balancing the momentum flux through the volume. Such a procedure is adequately discussed elsewhere (Ref. (17)) and results in the following expression:

$$C_N + i C_Y = 2 \sin \alpha \cos \alpha \left(\frac{r}{r_0} \right)^2 \left[1 + 2 \sum_{j=1}^k \lambda_j \left\{ \frac{r_j}{r} - \frac{r}{r_j} \right\} \right]$$

The first term in the above equation reduces to the slender body theory result of 2α at small angles of attack. The second term represents the crossflow drag due to the presence of the vortices, or the viscous contribution.

EVALUATION OF THE PROBLEM FORMULATION

The validity of the current formulation hinges on several points which are difficult to resolve in a conclusive manner. First and foremost, is the validity of the impulsive flow analogy. This analogy cannot be justified in its entirety. However, for inviscid flow about a slender body at small angle of attack, it is a consequence of slender body theory. At low Mach numbers where the flow is approximately incompressible, the restriction on small angle of attack can be relaxed. This does not, of course, shed any light on the validity of the impulse flow analogy for fluid flows in which viscous effects are important. All that can really be said in these cases is that the observed development of a vortex wake on the leeward side of a body at angle of attack bears at least a qualitative relation to the wake which develops behind a cylinder in crossflow (Ref. (18)).

Another point of question in the current problem formulation is the practice of simulating viscous flow by superimposing point vortices on inviscid, potential flow solutions. Even though this is an ad-hoc problem formulation, previous success in studies for cylinders in crossflow using a large number of point vortices indicate that this manner of approach provides an accurate wake description - in one case, good enough to predict the shedding frequency (Ref. (11)). The use of a single-point vortex to describe each wake vortex, as is done in this report, is more questionable. Such an approach is only capable of describing the gross flow-field features. Pressure calculations around the circumference of a cylinder using this type of description show an unrealistic peak near each point vortex. Despite this shortcoming, previous studies using this approach indicate that the integrated circumferential pressure distribution gives a good approximation of the normal forces at low angles of attack (Refs. (6), (7)).

The present approach is thus the extension of other studies which have proved to be accurate at lower angles of attack. Free parameters available in the current model allow adjustments to be made for effects which are not directly accounted for.

NUMERICAL CALCULATIONS

The governing equations are integrated using an Adams-Moulton predictor-corrector scheme which adjusts step size based on estimated error. This routine is started using the Runge-Kutta technique. The adjustable step size feature is extremely important in reducing running time. Small step sizes are required near the model nose where the local radius is small and at points at which new vortices are introduced. A listing of the program is provided in Appendix B.

Correct use of the outlined method requires the determination of a set of free parameter values that not only give a correct qualitative description of the wake, but also produce realistic forces and moments. There does not appear to be a unique set of parameters that best satisfy these requirements. However, certain sets of parameter values produce unacceptable results in that the computed crossflow wake does not develop into a Karman vortex street type of pattern. In some instances newly introduced vortices may move with diminishing circulation toward the cylinder surface. In any event, unless an orderly wake is formed, point vortices eventually approach one another resulting in numerical instability.

Numerical results were initially calculated for an impulsively started cylinder in crossflow. Suitable results were obtained using the set of parameters shown in Table 1. As indicated in this table, calculations were carried out assuming a symmetric wake containing two vortices, an unsymmetric wake still containing only a vortex pair, and an asymmetric wake in which vortices were periodically shed. In Figure 3 the calculated drag curves are plotted as a function of dimensionless time. Also shown in this graph is the experimental drag curve determined on an impulsively started cylinder

Table 1
FREE PARAMETER VALUES

Cylinder in Crossflow

Parameter	Symmetric Vortex Pair	Asymmetric Vortex Pair no shedding	Asymmetric wake with vortex shedding	Missile at angle of attack Mach no. taken into account
s	-	-	.1666	.1666
θ_{01}, θ_{02}	40° $44.6^\circ - 4.6 e^{-7(\frac{u_t}{r} - 10.2)}$	$\begin{cases} \text{if } \frac{u_t}{r} < 10.2 \\ \text{if } \frac{u_t}{r} \geq 10.2 \end{cases}$	$42^\circ + 7.5(F)^\circ$	
λ_I	.08	.08	.08	.04
$\Delta r_p/D$	0.0	.006	.006	.006
$\Delta r_s/D$	-	∞	.1	.1 + .05(F)
$\Delta \lambda_n$	-	-	-2.5%	-2.5%
r_n/D			.075	.075

$$F = - .835 + 1.67M - .244 \sin\left[\frac{\pi}{6}(M-.5)\right] \quad .5 \leq M \leq 1.1$$

in crossflow by Sarpkaya (Ref. (16)). Only the calculation which provides for vortex shedding is comparable to experimental results over wide ranges in dimensionless time. The accompanying wake development is illustrated in Figure 4.

In our initial attempt to apply this method to a body at angle of attack, the same set of free parameter values used for the impulsively started cylinder was assumed. Results were compared to experimental tests on a tangent ogive cylinder ($l = 12$ and $l_n = 4$). These tests are further described in Appendix C. The only adjustment to the set of free parameter values concerned the crossflow separation angles. As can be seen in Table 1 for the cylinder in crossflow, the separation angles are decreased at the dimensionless time $\frac{U_t}{r} = 10.2$. In application to the ogive cylinder at angle of attack, the point in the computation where this decrease occurs must be adjusted, since the change in radial dimensions at the missile nose disturbs the time sequence of the flow development. It is found that the point at which $\frac{U_t}{r} = 10.2$ occurs two-thirds of the distance in time between the point where the first and second vortices are shed. This criterion is used to determine where θ_{01} and θ_{02} will be decreased. As suggested by Schindel (Ref. (7)), the smooth transition from one separation angle to another is achieved by solving the equation:

$$\frac{d\theta_{0j}}{dx} = \frac{.7 \tan \alpha}{r} (44^\circ - \theta_{0j}) ; \quad j = 1, 2 \quad (6)$$

and subject to the boundary condition: $x \rightarrow \infty$ as $\theta_{0j} \rightarrow 48^\circ$. In the special case of a cylinder in crossflow, r is constant and the solution to 16 reduces to that shown in Table 1 for a cylinder in crossflow.

The calculated forces and pitching moments are compared to experimental results in Figures 5 and 6. In each of these figures experimental data for a range in Mach number from .5 to 1.1 are shown. The computed normal force coefficient (Fig. 5) compares well to the experimental data being closest to the curve generated at Mach .9. The calculated pitching moment coefficient (Fig. 6) is slightly high, matching best with the Mach 1.1 curve. The predicted yawing force coefficient, shown in Fig. 7, is somewhat larger than those measured and is closest to the Mach .5 curve.

The current theory is independent of Mach number. However, as can be seen from experimental data in Figures 5, 6, and 7, force and moment data is Mach number dependent. In order to include a Mach number effect in the theory, the set of free constants is assumed to be a weak function of Mach number. The relation between Mach number and free constant value is determined empirically by adjusting the set of constants to give the measured normal force curve at each Mach number. In addition, the sets of

free parameter values are modified somewhat to obtain better correspondence between observed and calculated vortex trails. The final set of constants is shown in the right-hand column of Table 1. These constant values yield stable vortex streets up to angles of attack near 50 degrees. The upper limit in each case is dependent on the actual nose configuration. Generally, the upper limit on angles of attack increases with decreasing nose fineness.

Computations using the Mach number dependent set of constants are compared to experiment in Figures 8 through 19. The close agreement between experimental and computed normal force coefficients are shown in Figure 8 while the pitching moment results are compared in Figure 9. The computed yawing force coefficient is compared to experimental measurements at Mach .5, .7, .9, and 1.1, in Figures 10, 11, 12, and 13, respectively. In each of these graphs four different yawing force coefficient curves are shown. They were generated by rolling the model nose. From these graphs it is evident that the yawing force coefficient is reasonably well predicted for Mach numbers less than .9. For larger Mach numbers the measured magnitudes of the yaw force coefficients decreases. This trend is not followed by the computed results. At the highest Mach number tested, 1.1, the yawing force coefficient is over-estimated by a factor of four.

In Figures 14 through 19, comparisons between the computed and observed vortex wakes are shown for angles of attack varying from 20 degrees to 45 degrees and Mach numbers of .5 and .9. It is evident from these pictures that the wake is qualitatively well reproduced. More consistent agreement is obtained at lower Mach numbers. At an angle of attack of 45 degrees and a Mach number of .9, the wake width is clearly over-estimated.

In order to assess the general applicability of this model, a comparison is made between calculated and experimental results for differently shaped ogive cylinders. Results at Mach .5 are shown in Figures 20 through 22. Good agreement is obtained for normal force. Although the pitching moment coefficient comparison is not as good as the normal force comparison, the experimental trend is followed. Figure 22 reveals that the decreasing yawing force coefficient observed on blunter bodies is not predicted by the theory.

Comparison between experimental and calculated results on differently shaped bodies at Mach 1.1 (Figures 23 through 25) shows that general trends found in experimental data for normal force and pitching moments are followed. However, agreement is not as close as in the case of lower Mach numbers. Also the side forces (Fig. 26) are consistently over-estimated and do not follow experimental trends.

SUMMARY AND CONCLUSIONS

The flow field surrounding an ogive cylinder at high angles of attack is simulated using an extension of the method initially proposed by Bryson (Ref. (6)). This method employs the impulsive flow analogy and models each wake vortex using a single point vortex. Free parameters inherent in the problem formulation are initially determined by comparison to drag force coefficient data on an impulsively started cylinder in crossflow.

Final parameter values are obtained by matching the normal force measured on a tangent ogive cylinder ($l_n = 4.0$, $l = 12$) over a range in Mach number of .5 to 1.1. These parameter values are made weakly Mach number dependent to reflect the experimental data. The resulting pitching moments compare well in trend with experiment. The calculated yaw force coefficient is of the correct magnitude at Mach numbers less than .9, but is too large at higher Mach numbers.

At a Mach number of .5 calculated and experimental results are compared for ogive cylinders with different nose fineness ratios. Predicted normal force and pitching moment coefficients are in good agreement with experimental results. On ogive cylinders with small slenderness ratios ($l_n \lesssim 3$) yawing force calculations over-estimated the experimental results.

The same type of comparison is carried out at a Mach number of 1.1. Agreement between calculated normal force and pitching moment coefficients was not as good as at Mach .5. The yawing force coefficient was consistently lower than predicted.

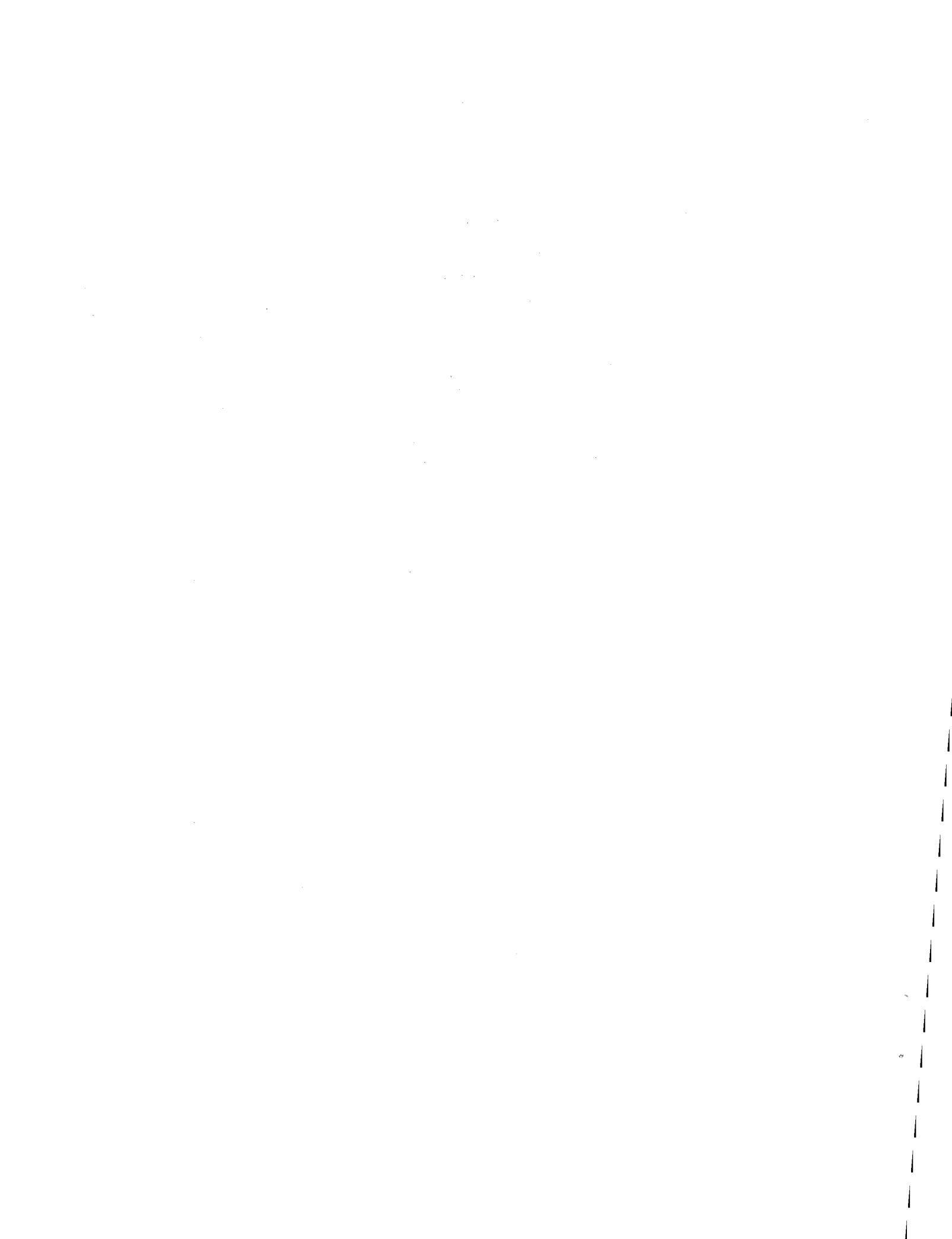
The current results indicate that on slender bodies and at low Mach numbers, the impulsive flow analogy combined with the Bryson method can be used to predict normal force, pitching moment and yawing force. For nose fineness ratios less than 4 and at Mach numbers near unity, accuracy deteriorates markedly and side forces are over-estimated.

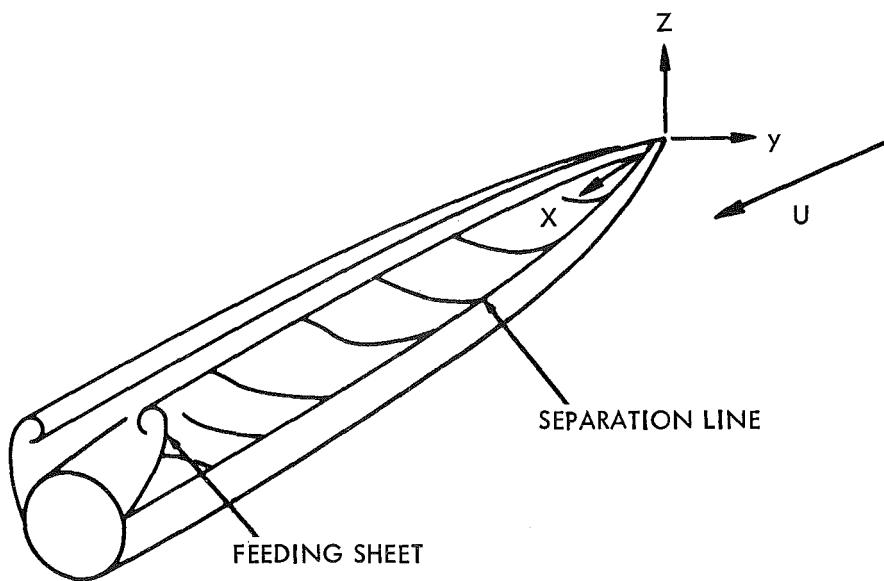
Apparently the separation parameters are sensitive to Mach number in the transonic range. A more thorough investigation of this effect would be required to improve the accuracy of side force predictions in this regime.

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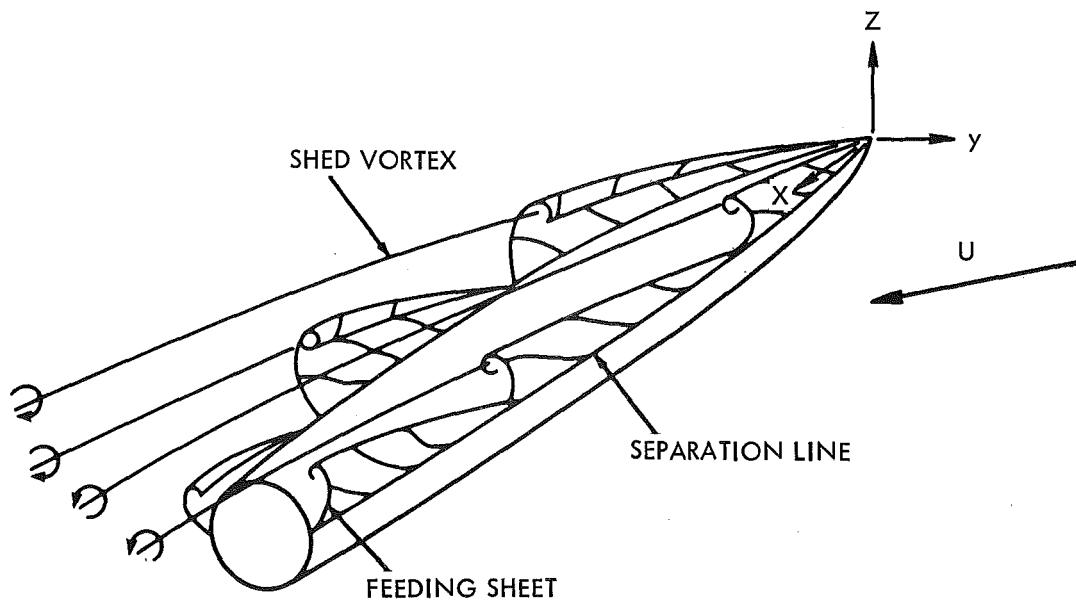
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(a) SYMMETRICAL VORTEX SHEDDING



(b) ASYMMETRICAL VORTEX SHEDDING

FIG. 1 VORTEX FORMATION ON A BODY OF REVOLUTION AT ANGLE OF A ATTACK

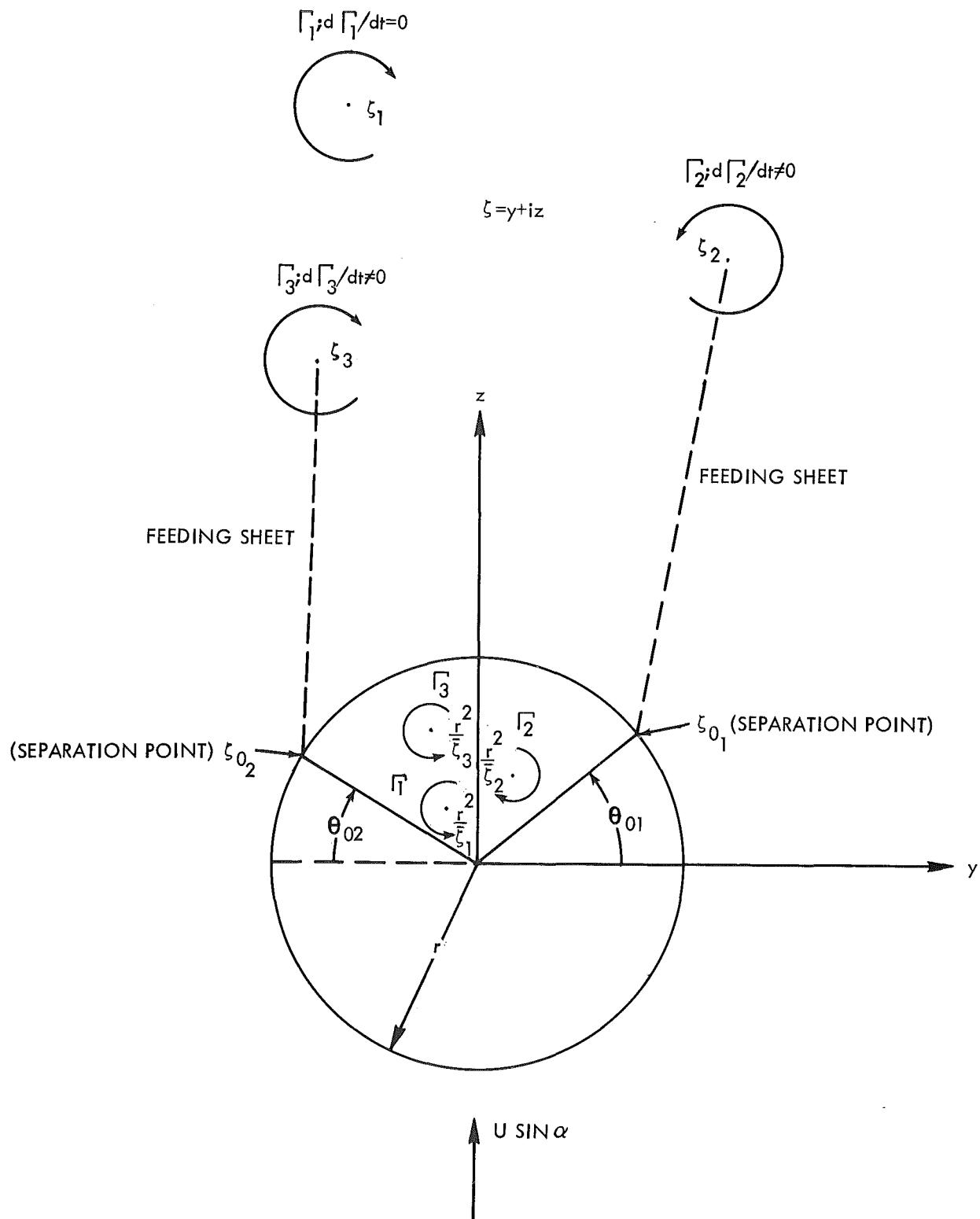


FIG. 2 DIAGRAM OF THE CROSSFLOW PLANE.

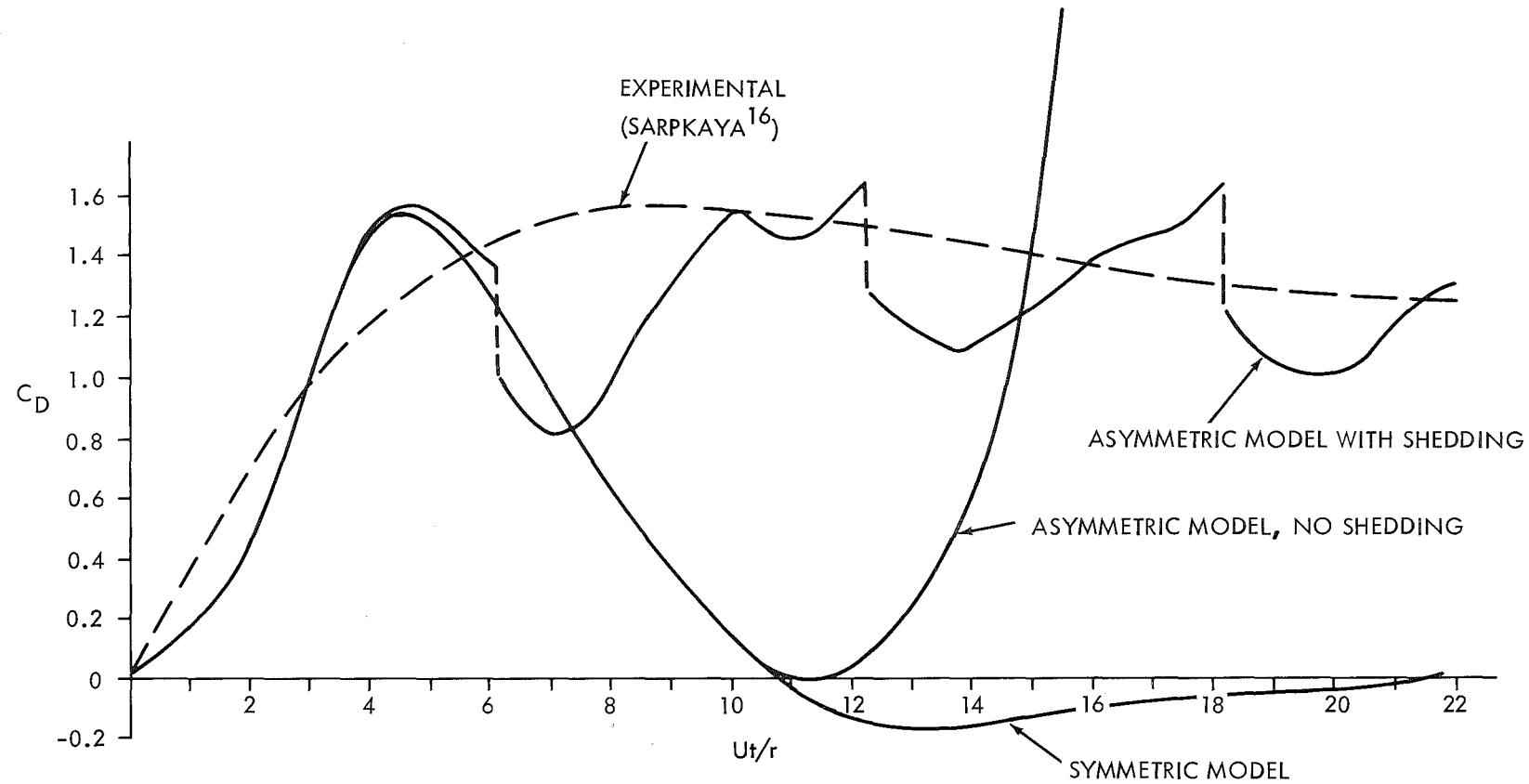


FIG. 3 DRAG FORCE COEFFICIENT AS A FUNCTION OF DIMENSIONLESS TIME FOR AN IMPULSIVELY STARTED CYLINDER IN CROSSFLOW.

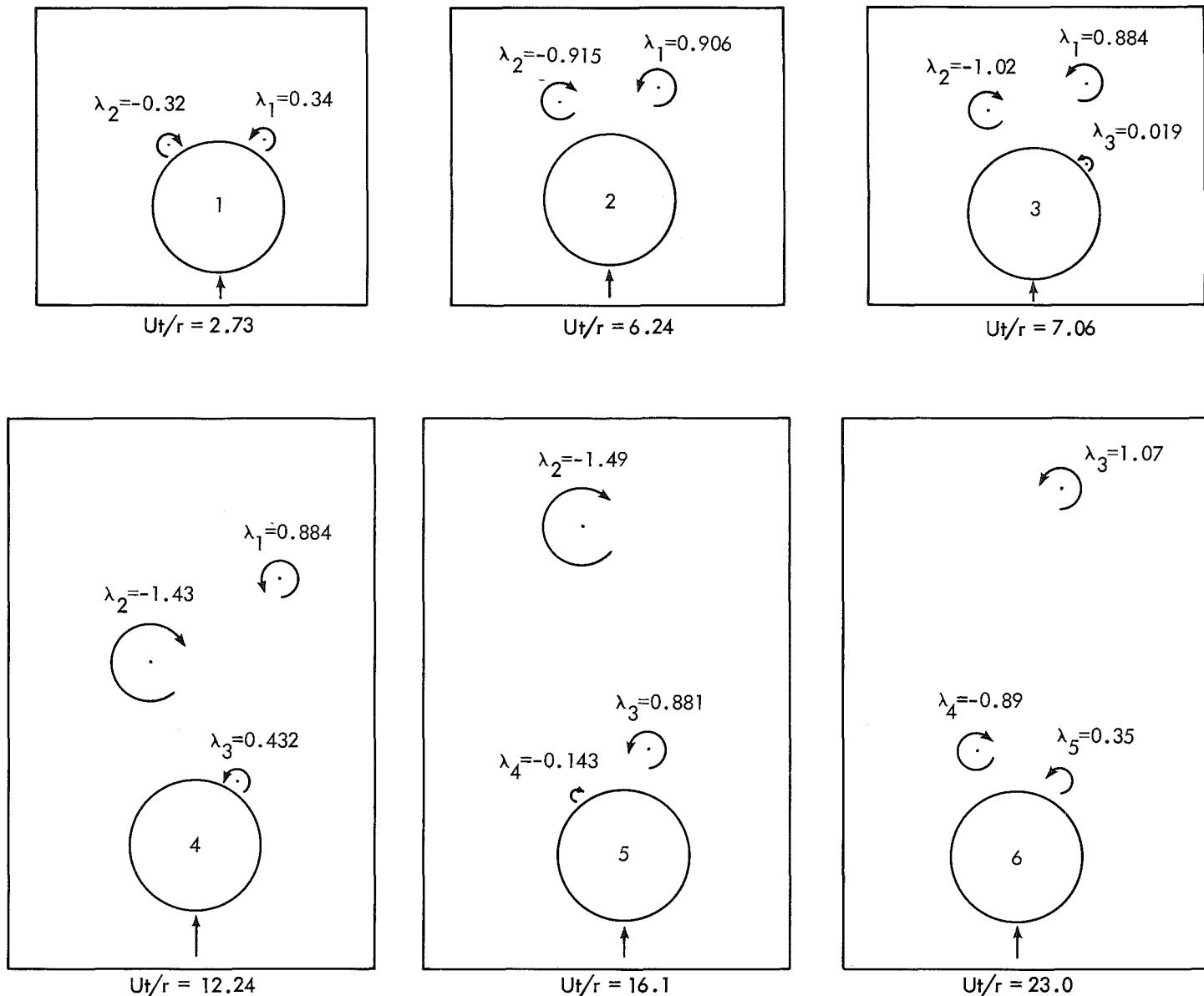


FIG. 4 DEVELOPMENT OF THE FLOW FIELD BEHIND AN IMPULSIVELY STARTED CYLINDER
 IN CROSSFLOW.

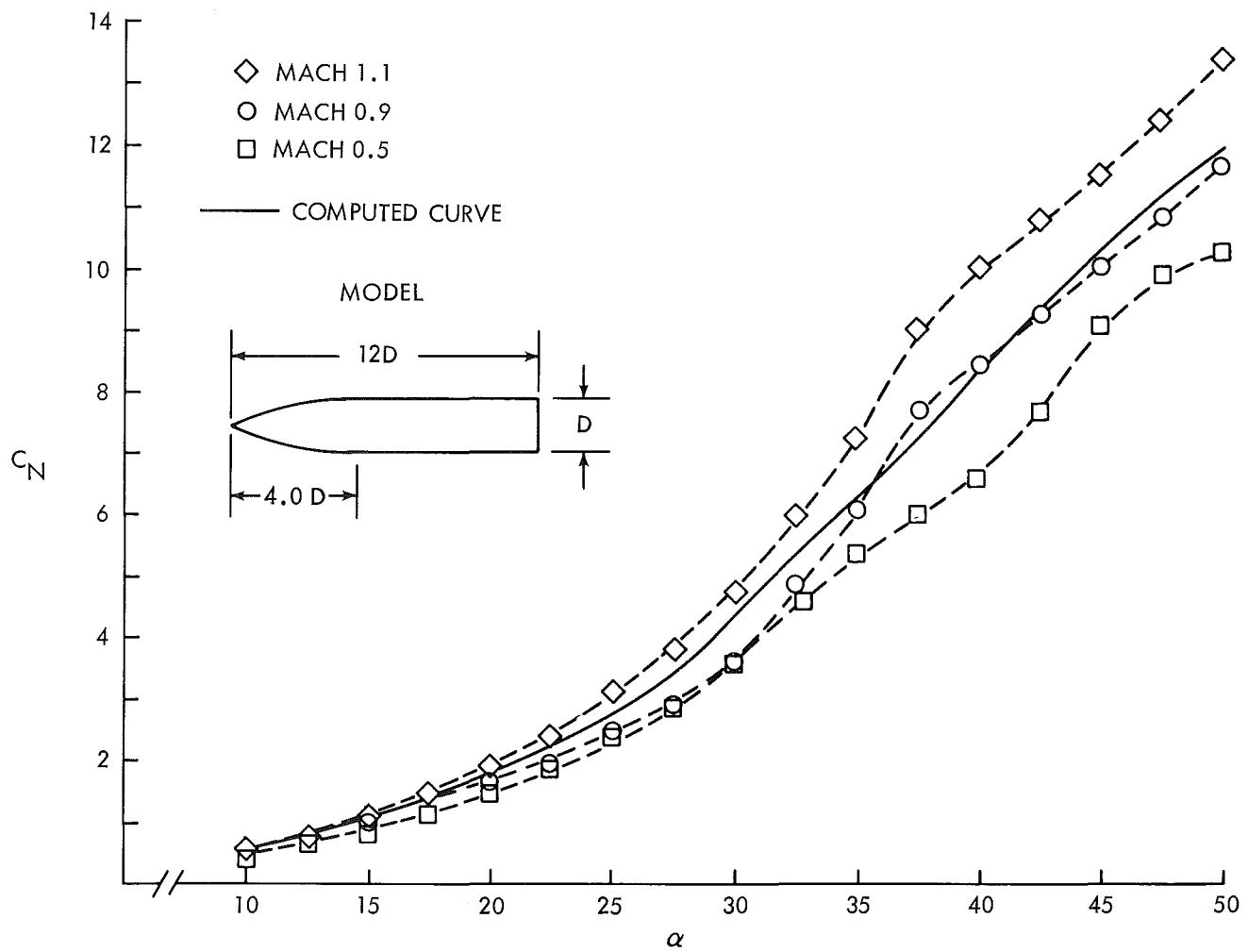


FIG. 5 NORMAL FORCE COEFFICIENT AS A FUNCTION OF ANGLE OF ATTACK. COMPUTED CURVE BASED ON FREE PARAMETER VALUES USED FOR A CYLINDER IN CROSSFLOW.

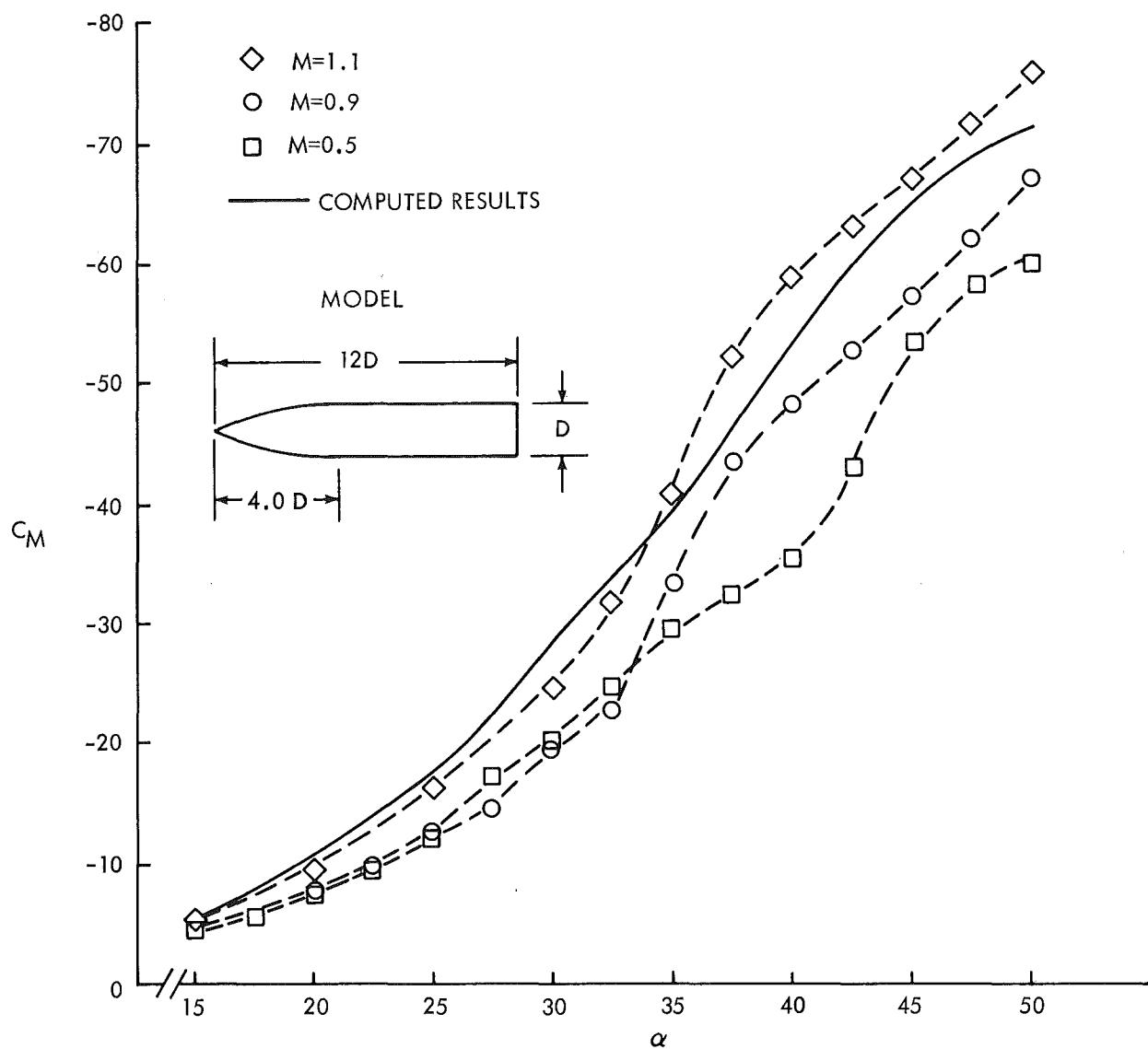


FIG. 6 PITCHING MOMENT COEFFICIENT AS A FUNCTION OF ANGLE OF ATTACK.
COMPUTED CURVE BASED ON FREE PARAMETER VALUES USED FOR A
CYLINDER IN CROSSFLOW.

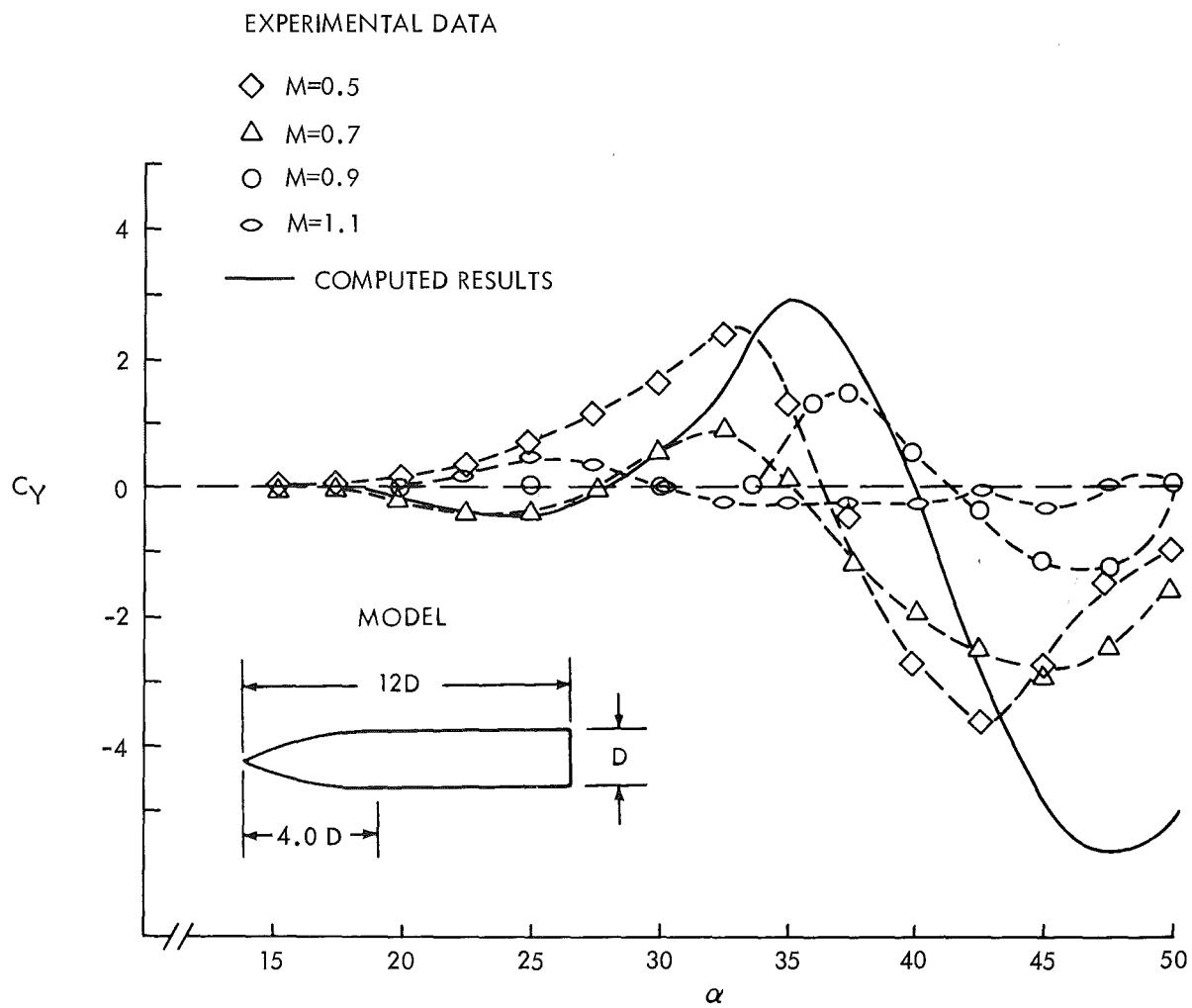


FIG. 7 YAWING FORCE COEFFICIENT AS A FUNCTION OF ANGLE OF ATTACK. COMPUTED CURVE BASED ON THE FREE PARAMETER VALUES USED FOR A CYLINDER IN CROSSFLOW.

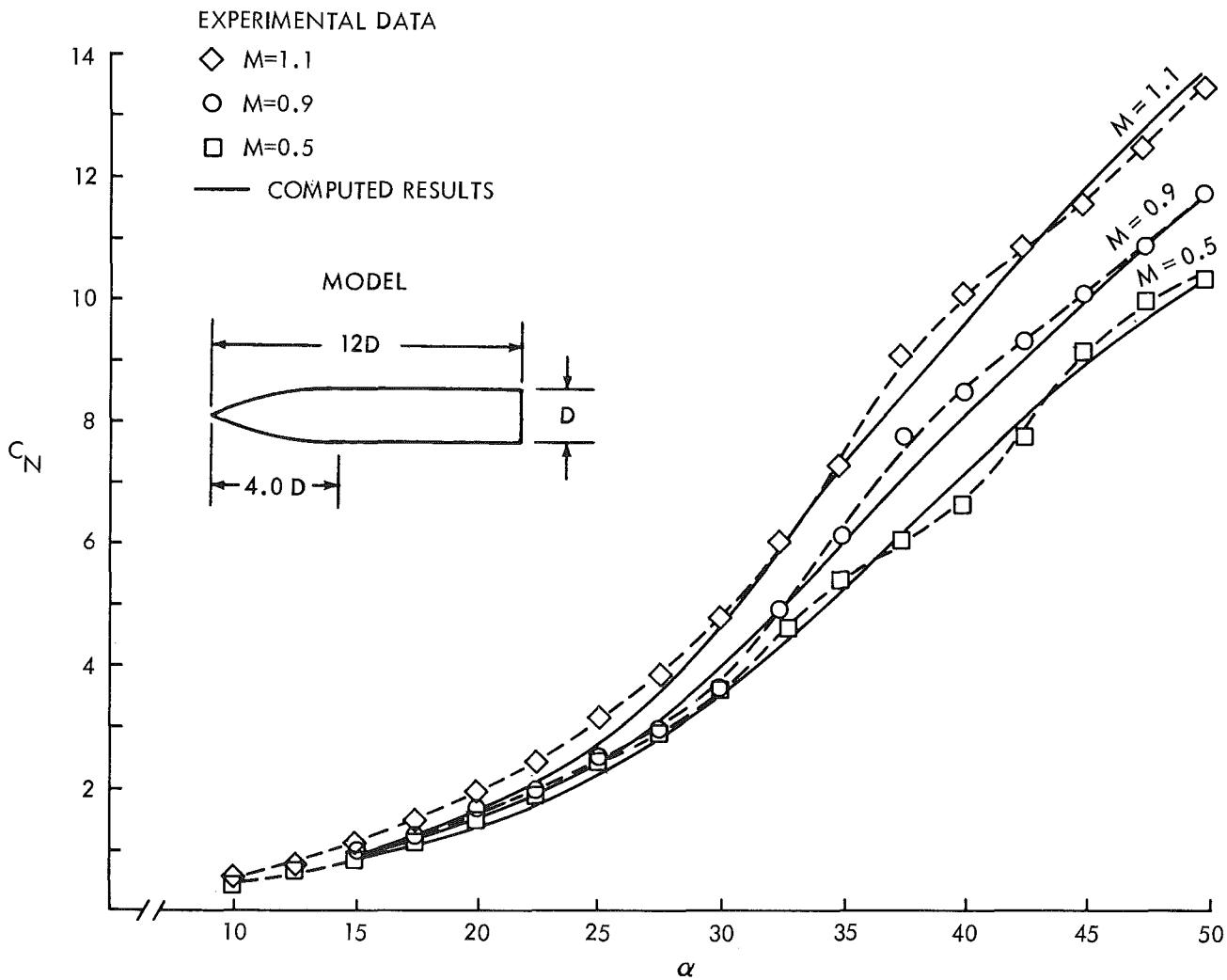


FIG. 8 NORMAL FORCE COEFFICIENT AS A FUNCTION OF ANGLE OF ATTACK.
COMPUTED CURVES ARE BASED ON A SET OF PARAMETER VALUES WHICH
ARE MACH NUMBER DEPENDENT.

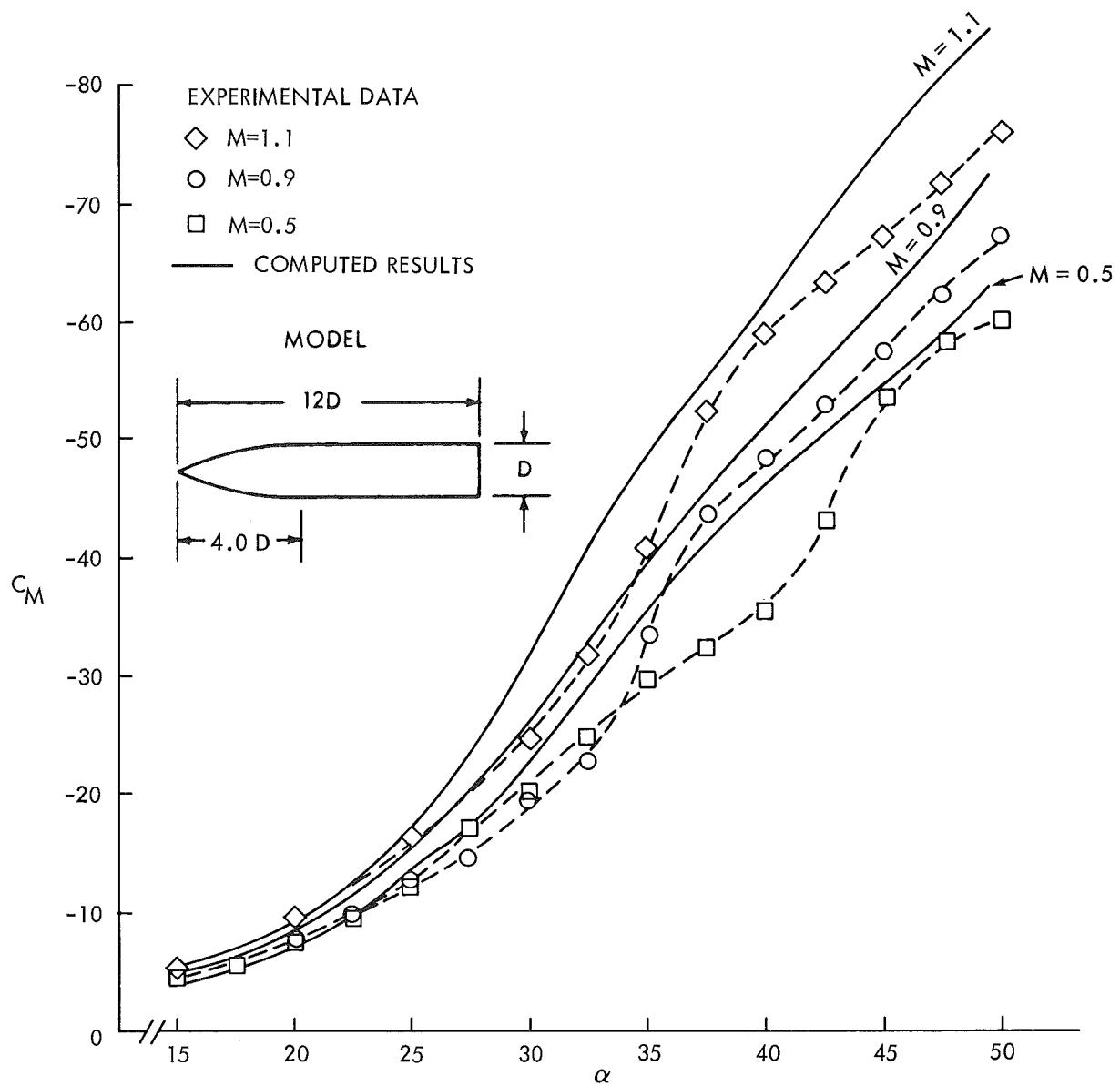


FIG. 9 PITCHING MOMENT COEFFICIENT AS A FUNCTION OF ANGLE OF ATTACK. COMPUTED CURVES ARE BASED ON A SET OF PARAMETER VALUES WHICH ARE MACH NUMBER DEPENDENT.

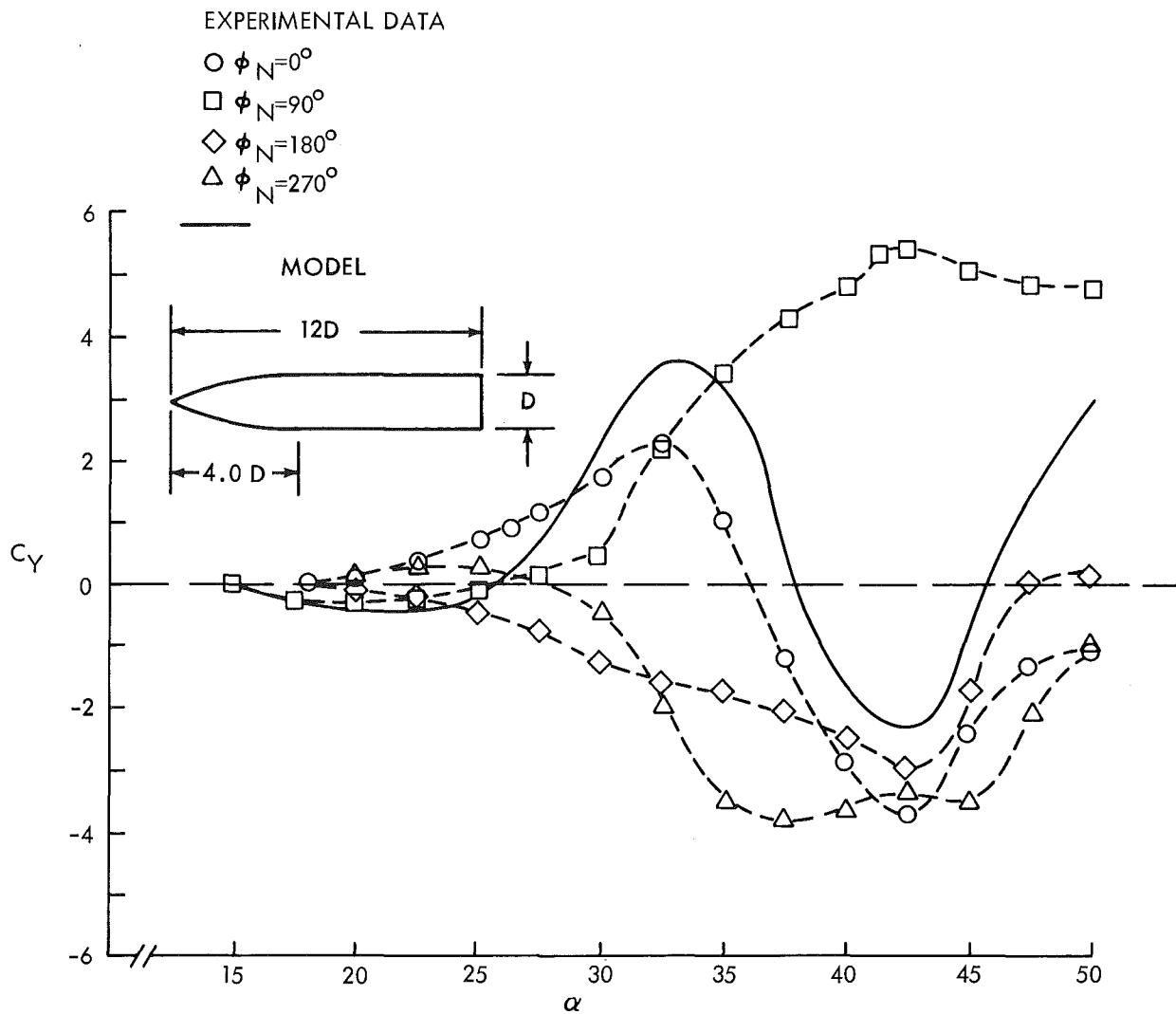


FIG. 10 YAW FORCE COEFFICIENT AS A FUNCTION OF ANGLE OF ATTACK FOR $M=0.5$. COMPUTED CURVES BASED ON FREE PARAMETER VALUES WHICH ARE MACH NUMBER DEPENDENT.

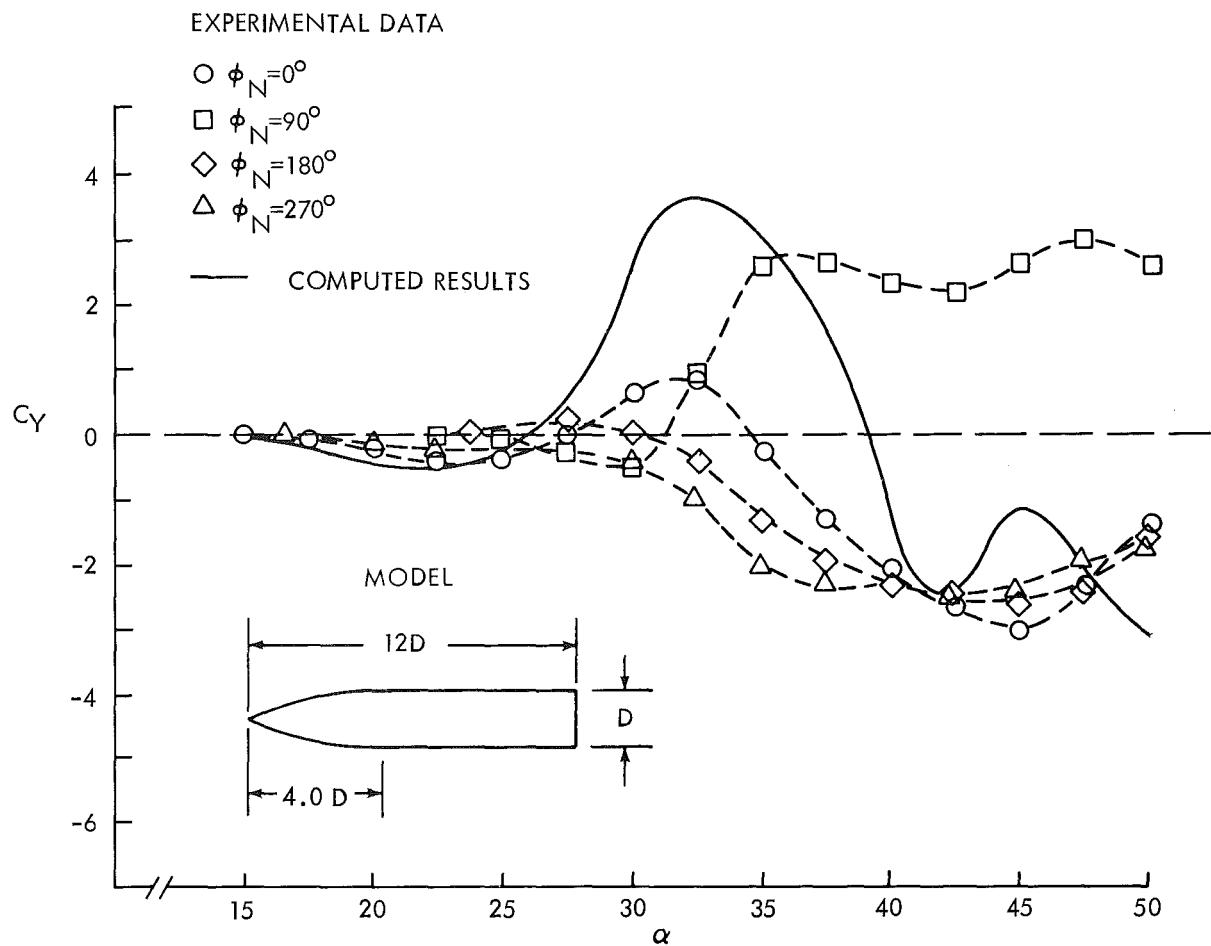


FIG. 11 YAW FORCE COEFFICIENT AS A FUNCTION OF ANGLE OF ATTACK FOR $M=0.7$. COMPUTED CURVE BASED ON FREE PARAMETER VALUES WHICH ARE MACH NUMBER DEPENDENT.

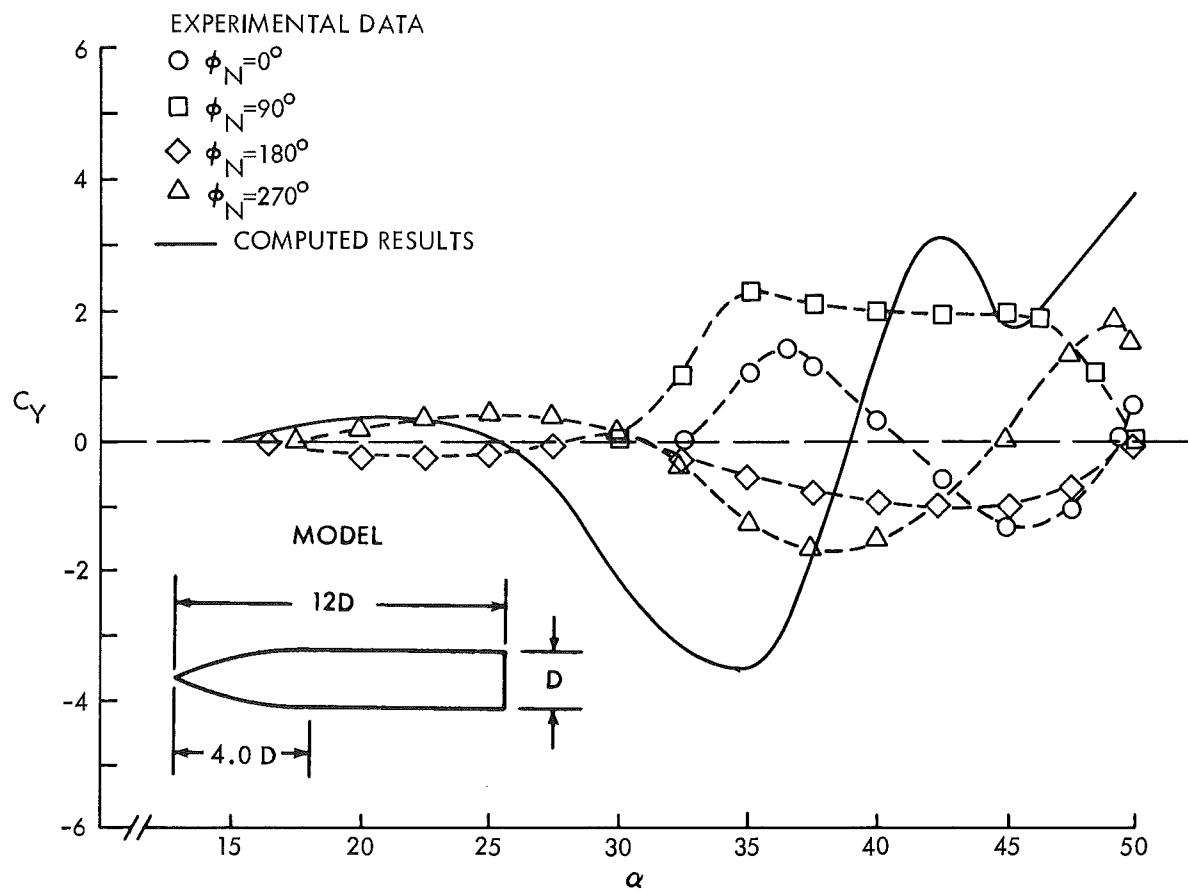


FIG. 12 YAW FORCE COEFFICIENT AS A FUNCTION OF ANGLE OF ATTACK FOR $M=0.9$. COMPUTED CURVE BASED ON FREE PARAMETER VALUES WHICH ARE MACH NUMBER DEPENDENT.

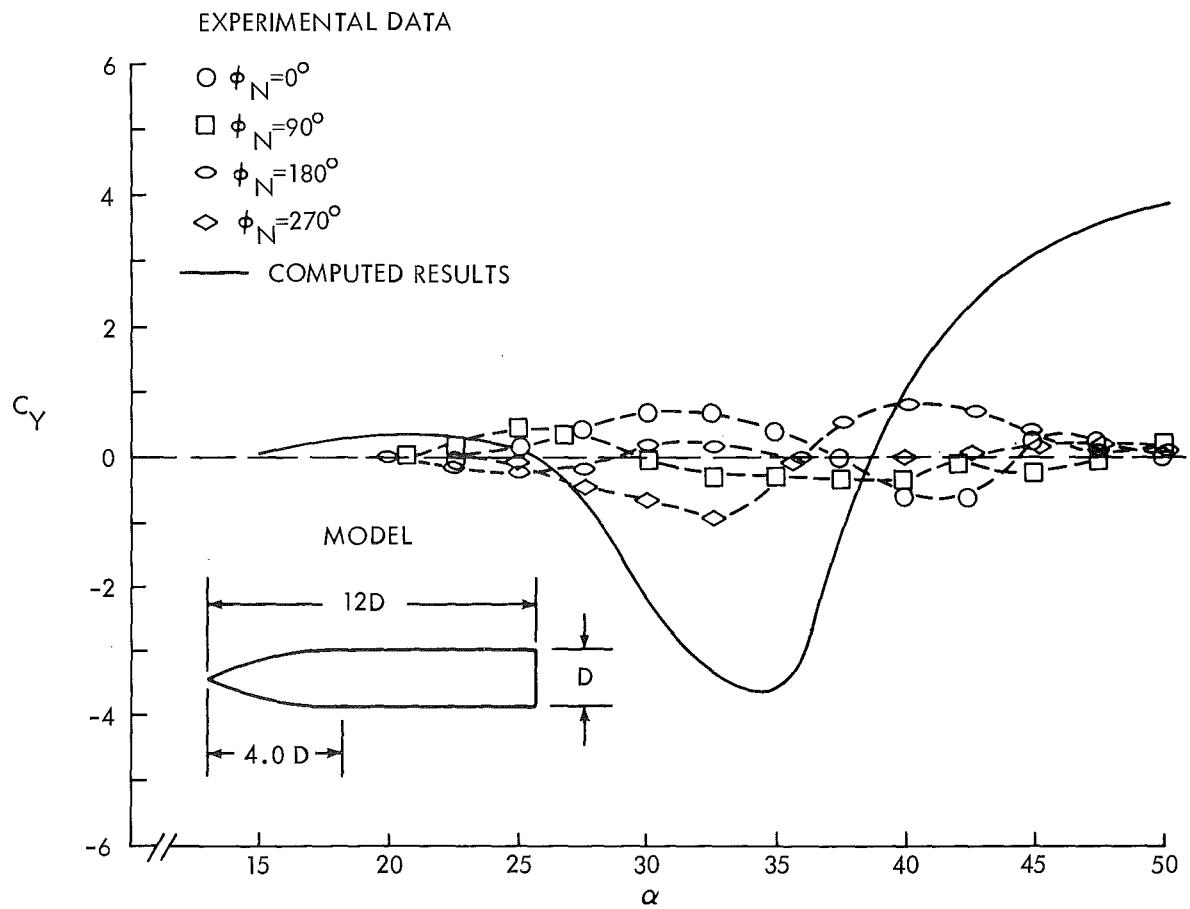


FIG. 13 YAW FORCE COEFFICIENT AS A FUNCTION OF ANGLE OF ATTACK FOR $M=1.1$. COMPUTED CURVE BASED ON FREE PARAMETER VALUES WHICH ARE MACH NUMBER DEPENDENT.

III III II CORRIDOR IN WHICH VORTICES ARE OBSERVED.

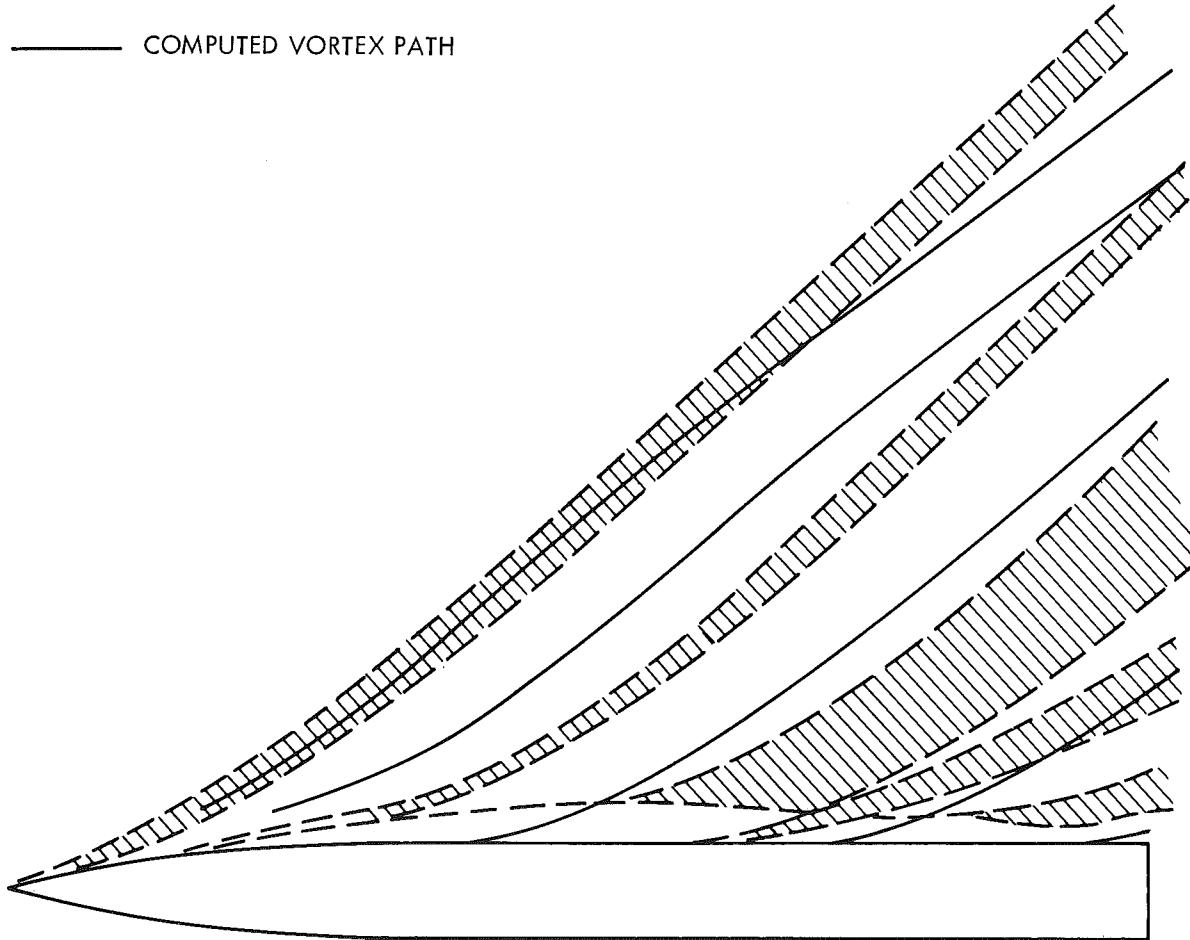


FIG. 14 COMPARISON OF OBSERVED AND COMPUTED VORTEX PATHS. $\alpha = 45^\circ$ MACH=0.5

□□□□ CORRIDOR IN WHICH VORTICES ARE OBSERVED.
— COMPUTED VORTEX PATH.

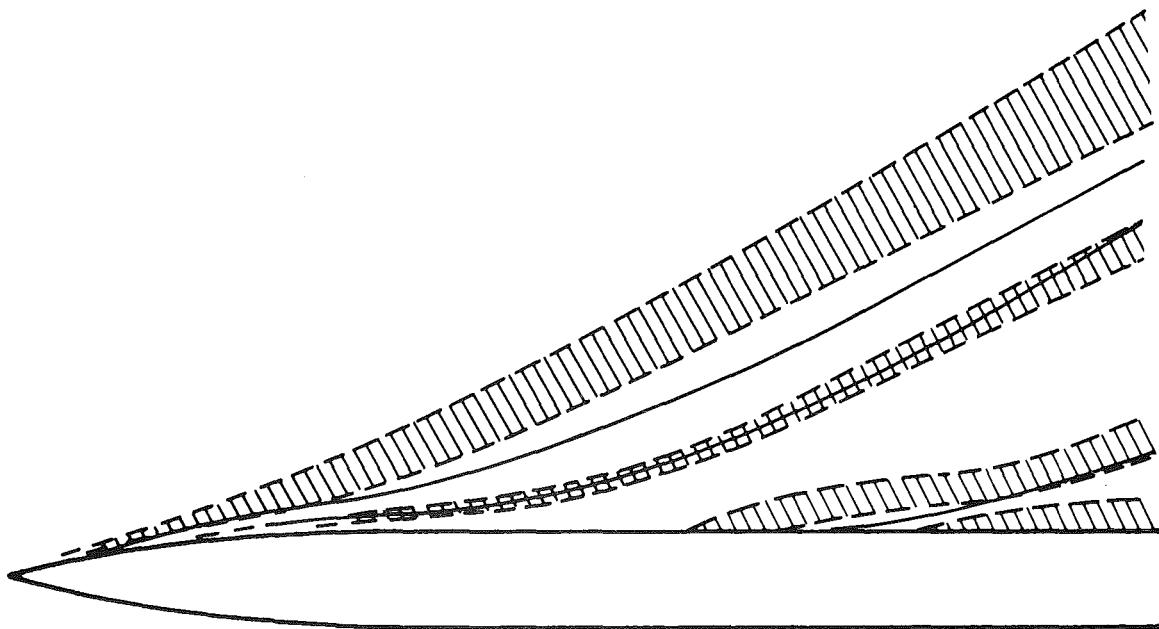


FIG. 15 COMPARISON OF OBSERVED AND COMPUTED VORTEX PATHS.
 $\alpha = 35^\circ$ MACH = 0.5

— — — OBSERVED VORTEX PATH.
— COMPUTED VORTEX PATH.



FIG. 16 COMPARISON OF OBSERVED AND COMPUTED VORTEX PATHS.
 $\alpha = 20^\circ$ MACH = 0.5

□□□□ CORRIDOR IN WHICH VORTICES ARE OBSERVED.

— COMPUTED VORTEX PATH.

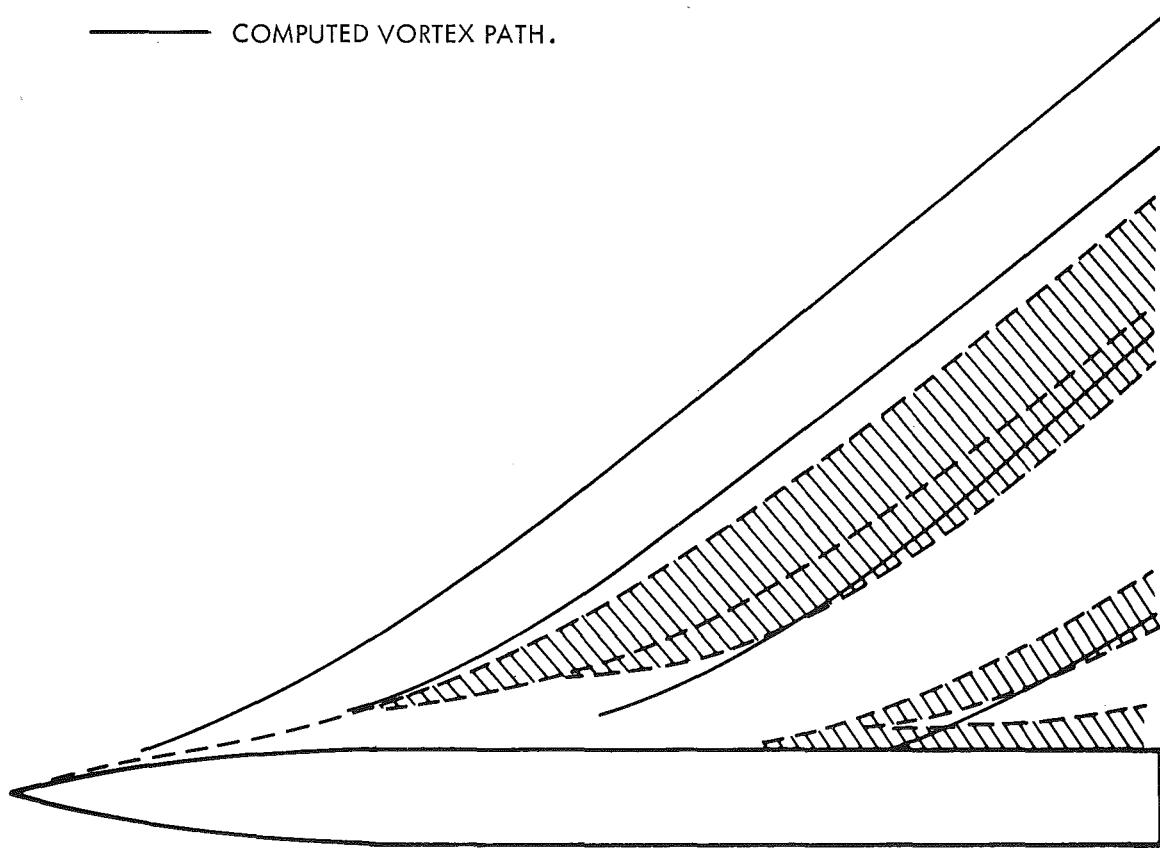


FIG. 17 COMPARISON OF OBSERVED AND COMPUTED VORTEX PATHS.
 $\alpha = 45^\circ$ MACH = 0.9

□□□□□ CORRIDOR IN WHICH VORTICES ARE OBSERVED.

— COMPUTED VORTEX PATH.

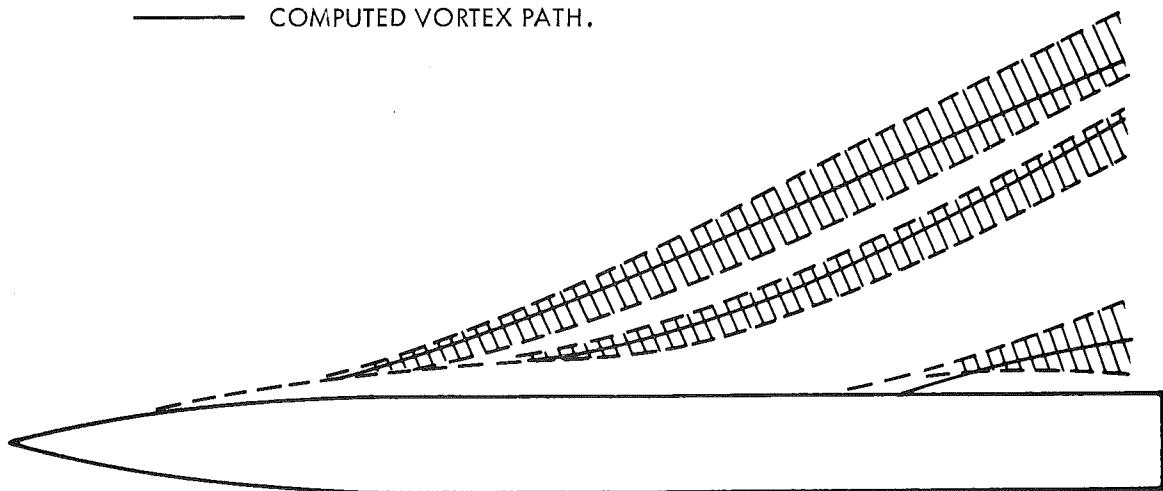


FIG. 18 COMPARISON OF OBSERVED AND COMPUTED VORTEX PATHS.

$\alpha = 35^\circ$ MACH = 0.9

— - - - OBSERVED VORTEX PATH.

— COMPUTED VORTEX PATH.

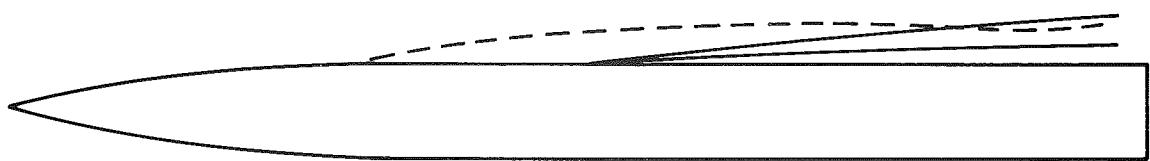


FIG. 19 COMPARISON OF OBSERVED AND COMPUTED VORTEX PATHS.

$\alpha = 20^\circ$ MACH = 0.9

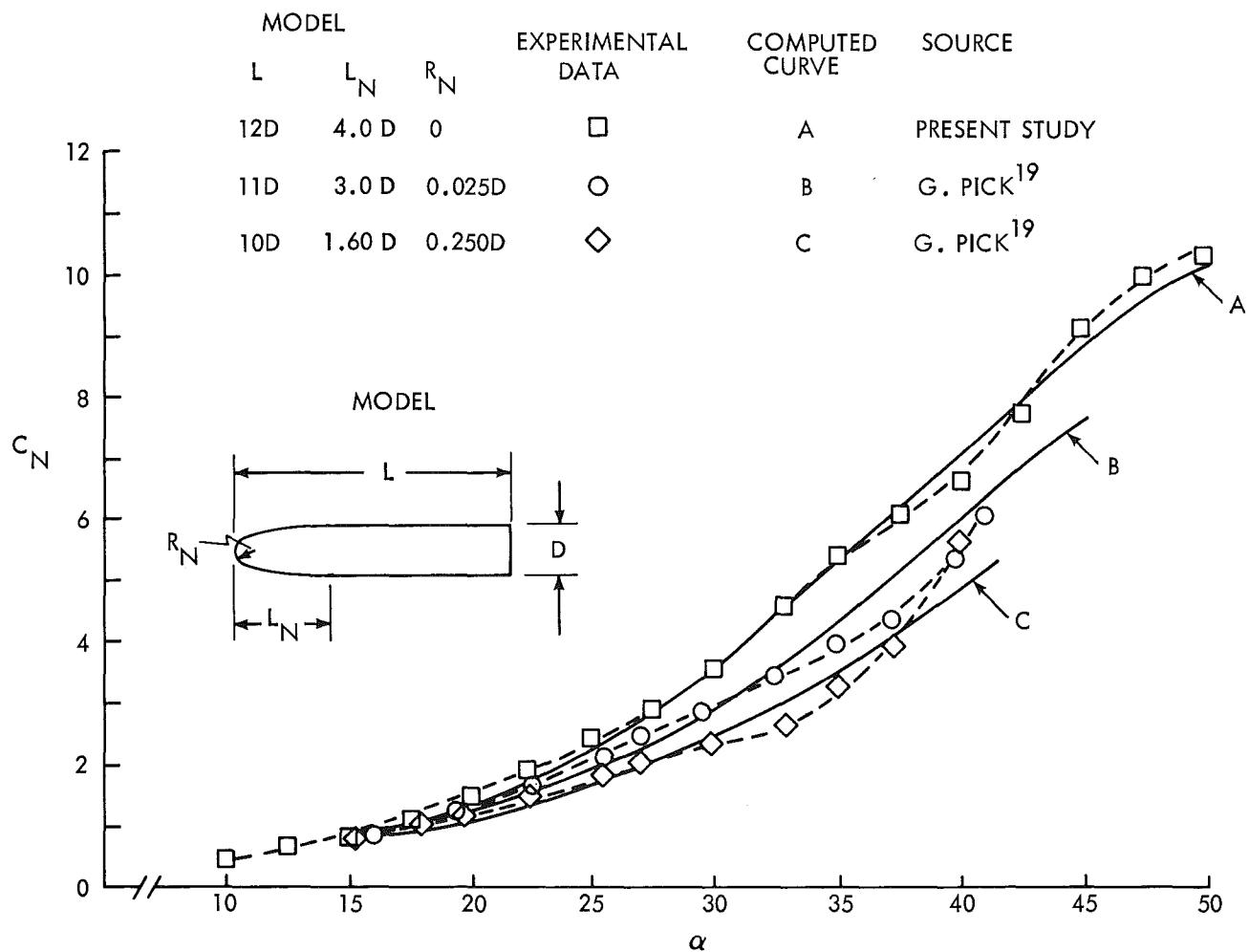


FIG. 20 NORMAL FORCE COEFFICIENT AS A FUNCTION OF ANGLE OF ATTACK FOR SEVERAL DIFFERENT MODELS, $M = 0.5$

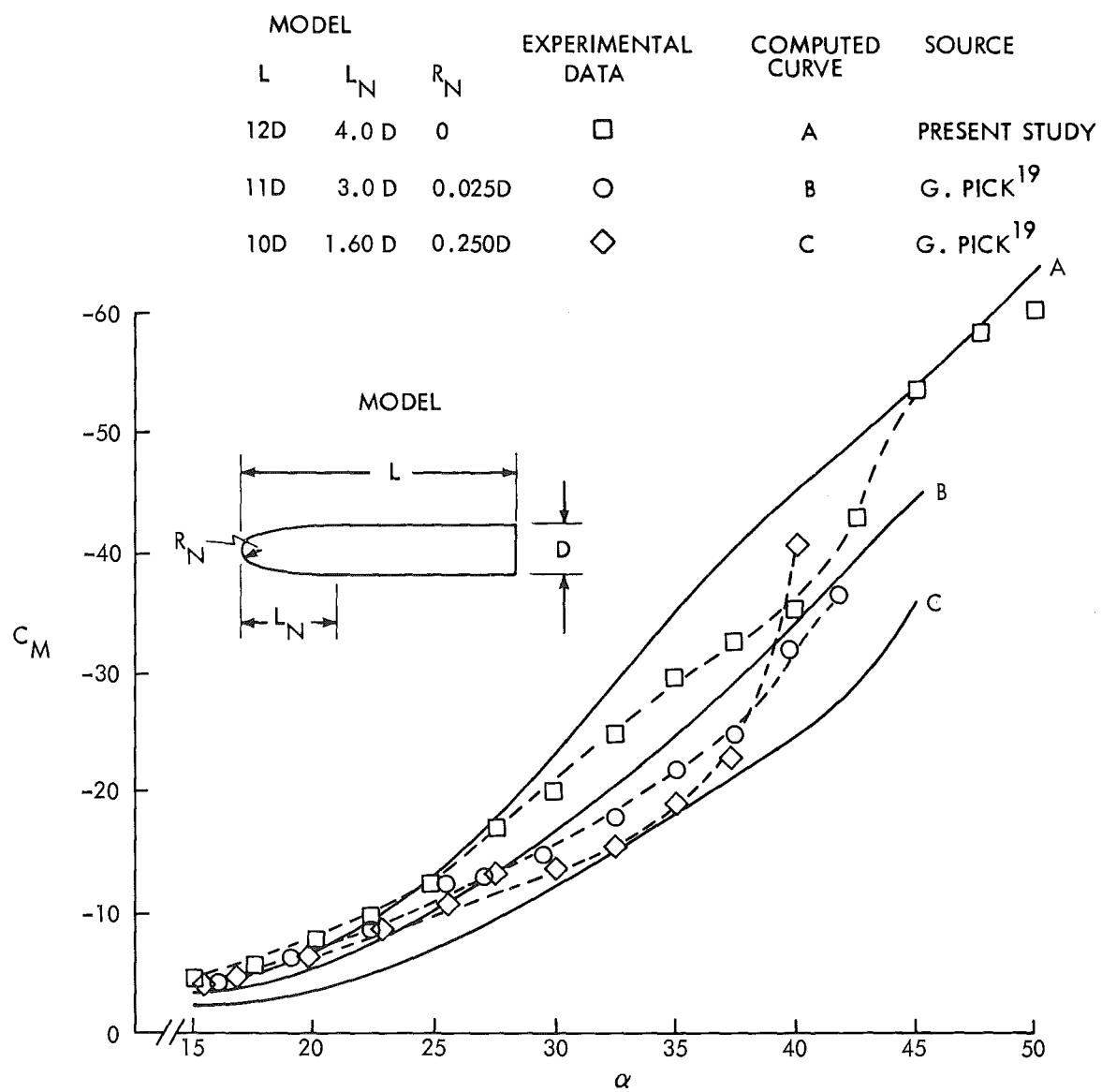


FIG. 21 PITCHING MOMENT COEFFICIENT AS A FUNCTION OF ANGLE OF ATTACK FOR SEVERAL DIFFERENT MODELS, $M = 0.5$.

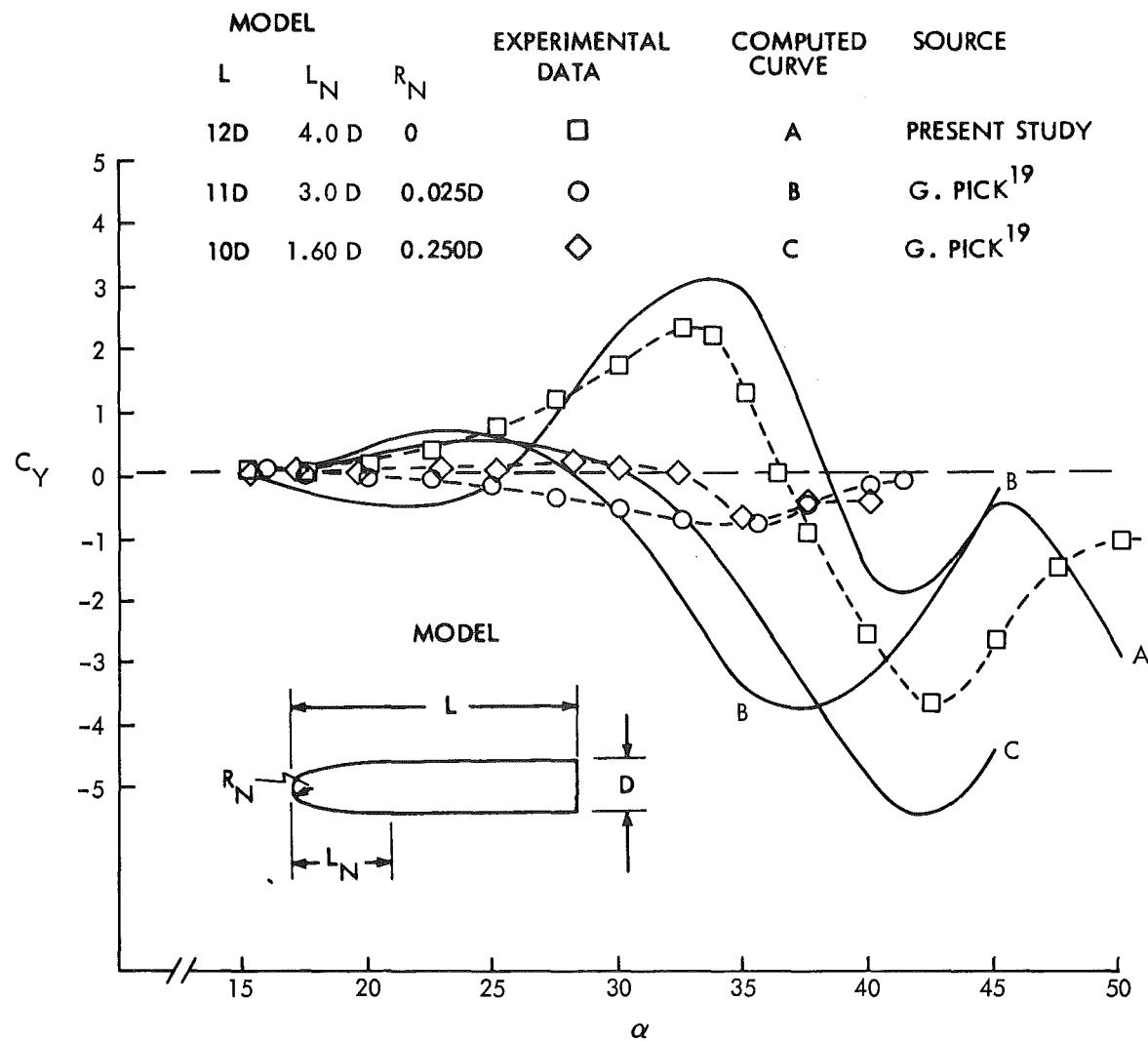


FIG. 22 YAW FORCE COEFFICIENT AS A FUNCTION OF ANGLE OF ATTACK FOR SEVERAL DIFFERENT MODELS, $M = 0.5$.

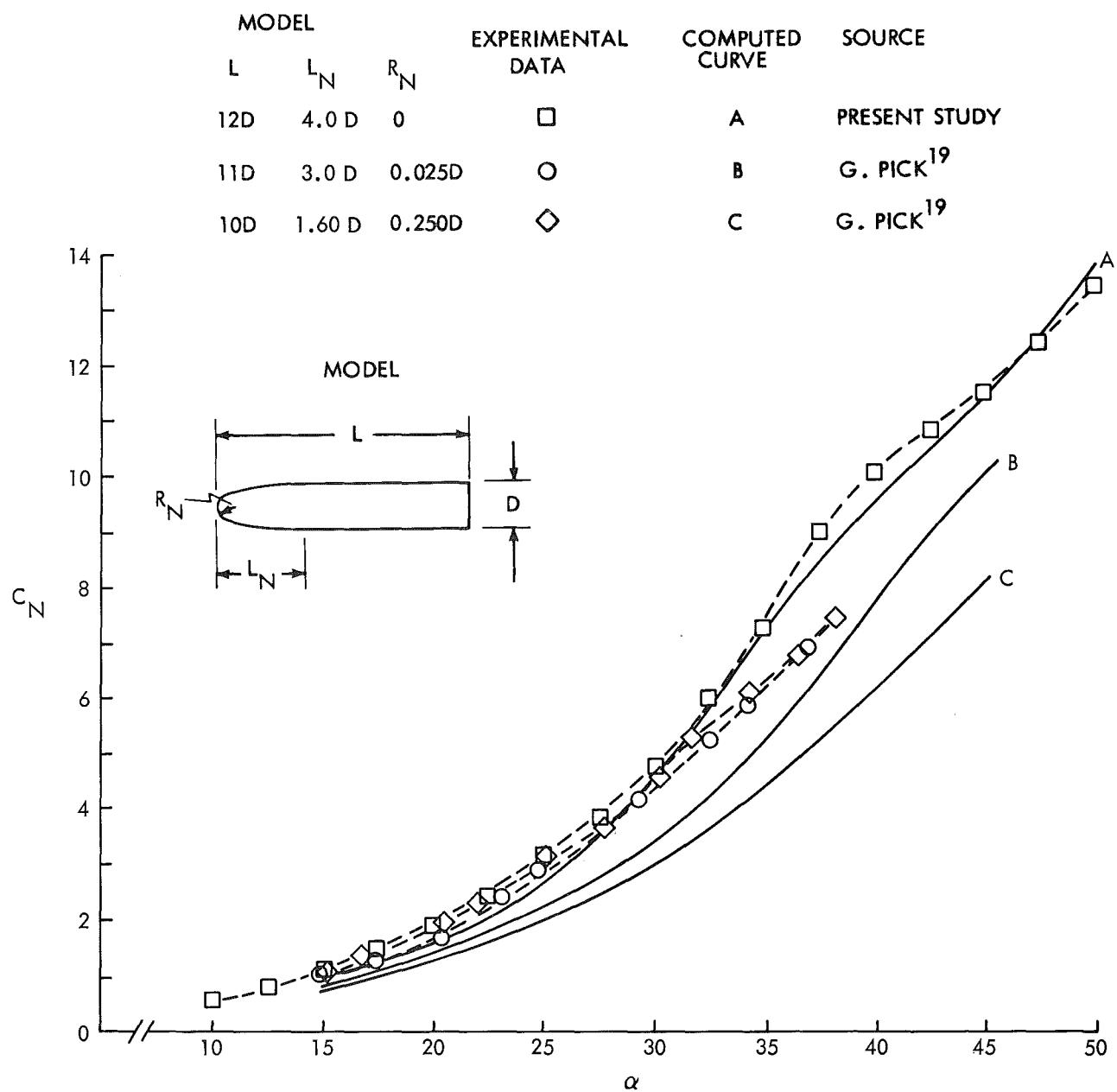


FIG. 23 NORMAL FORCE COEFFICIENT AS A FUNCTION OF ATTACK
FOR SEVERAL DIFFERENT MODELS, $M = 1.1$.

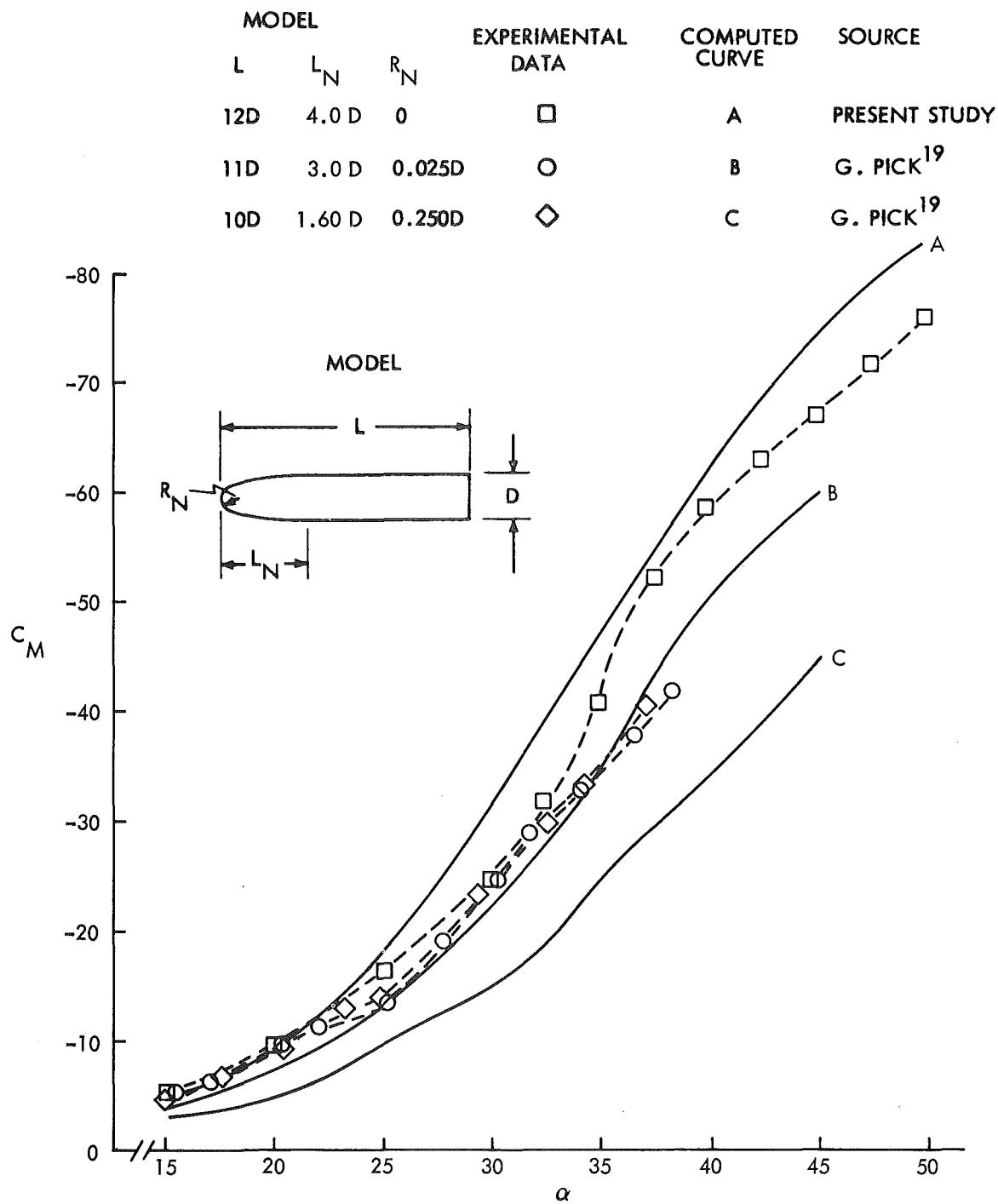


FIG. 24 PITCHING MOMENT COEFFICIENT AS A FUNCTION OF ANGLE OF ATTACK FOR SEVERAL DIFFERENT MODELS, $M = 1.1$

MODEL			EXPERIMENTAL DATA	COMPUTED CURVE	SOURCE
L	L_N	R_N			
12D	4.0 D	0	□	A	PRESENT STUDY
11D	3.0 D	0.025D	○	B	G. PICK ¹⁹
10D	1.60 D	0.250D	◇	C	G. PICK ¹⁹

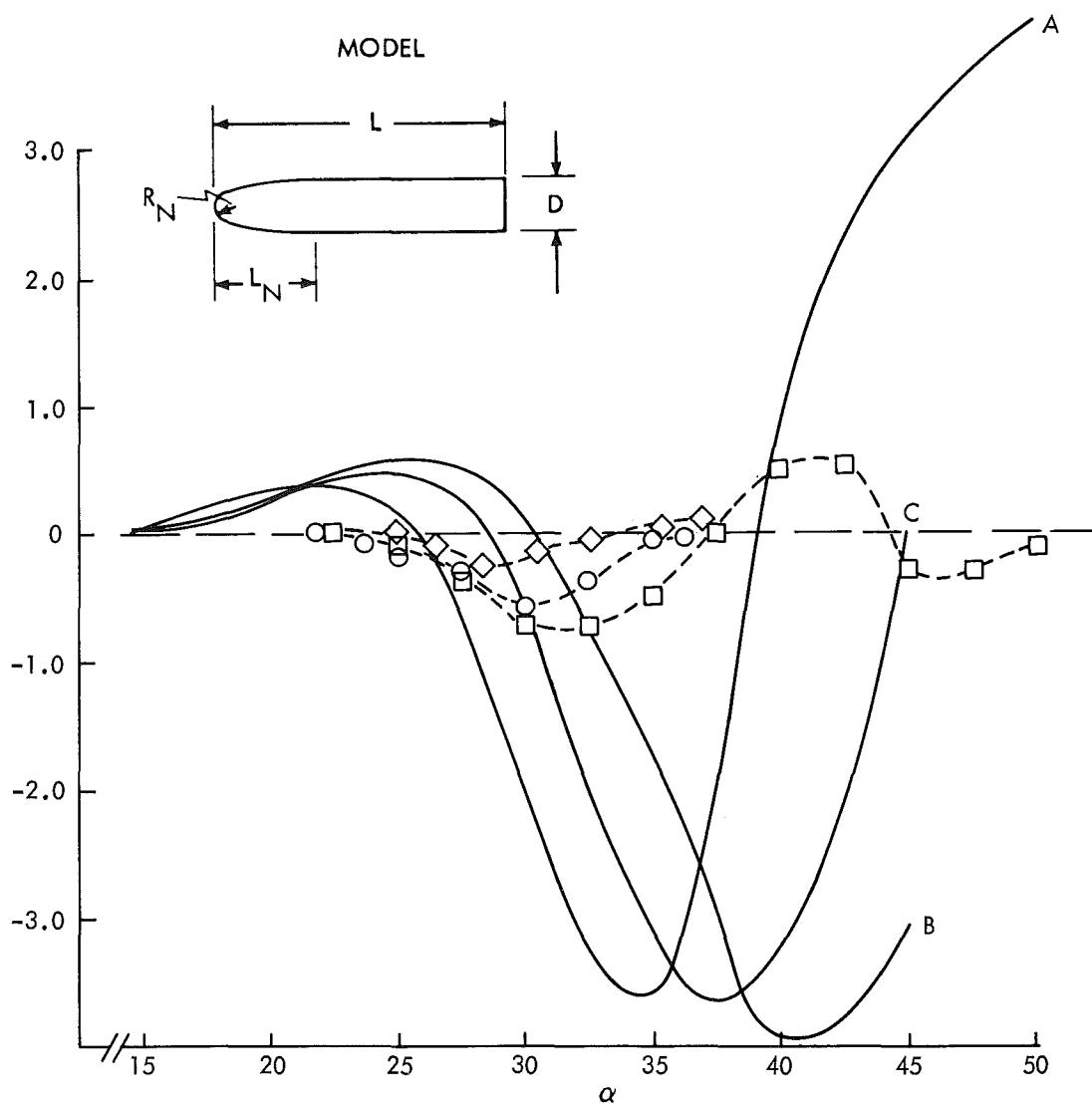
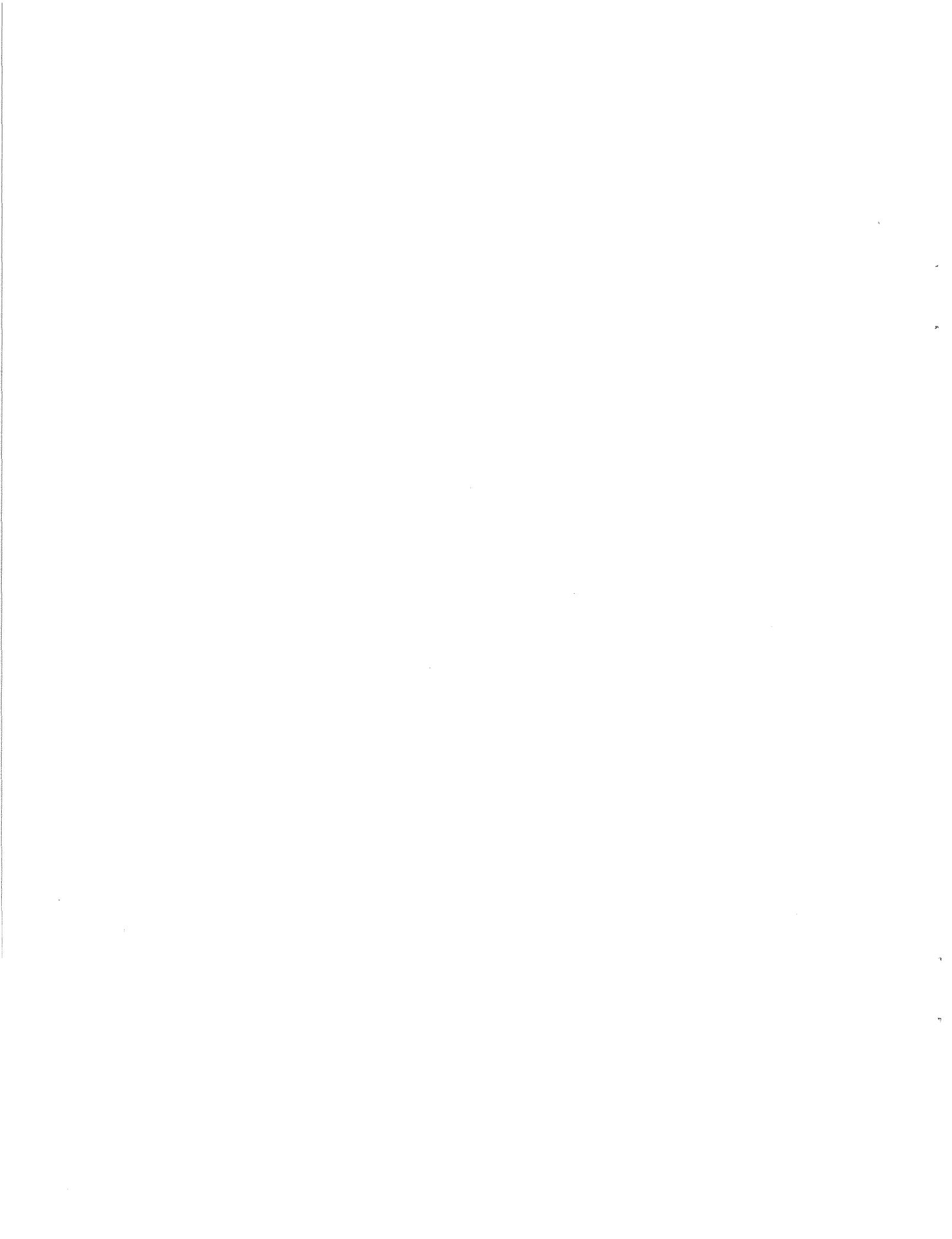


FIG. 25 YAW FORCE COEFFICIENT AS A FUNCTION OF ANGLE OF ATTACK FOR SEVERAL DIFFERENT MODELS, $M = 1.1$



APPENDIX A

DERIVATION OF PROGRAMMED EQUATIONS

The governing equations are:

$$\left. \begin{aligned} \frac{ds_1}{dx} + \frac{(s_1 - s_{01})}{\lambda_1} \frac{d\lambda_1}{dx} &= \tan \alpha \omega_1 - \frac{(s_1 - s_{01})}{r} \frac{dr}{dx} \\ \frac{ds_2}{dx} + \frac{(s_2 - s_{02})}{\lambda_2} \frac{d\lambda_2}{dx} &= \tan \alpha \omega_2 - \frac{(s_2 - s_{02})}{r} \frac{dr}{dx} \\ \frac{ds_j}{dx} = \tan \alpha \omega_j \quad j = 3, 4, \dots, k & \end{aligned} \right\} \begin{matrix} \text{growing} \\ \text{vortices} \end{matrix}$$

shed vortices (A-1)

In order to define λ_1 and λ_2 the above equation must be augmented by the constraint that the fluid velocity be zero at the stagnation points relative to the changing radius:

$$\omega(s_{01}) = \frac{r}{\tan \alpha} \frac{dr}{dx} \frac{1}{s_{01}} \quad \omega(s_{02}) = \frac{r}{\tan \alpha} \frac{dr}{dx} \frac{1}{s_{02}} \quad (A-2)$$

The dimensionless velocity potential, ϕ , is defined as follows:

$$\begin{aligned} \tilde{\phi}(s) &= -\left(s - \frac{r^2}{s}\right)^{\frac{1}{2}} - ir \sum_{j=1}^k \lambda_j \ln \left[\frac{(s - s_j)}{(s - r^2/s_j)} \right] \\ &\quad + \frac{r}{\tan \alpha} \frac{dr}{dx} \ln(s) \end{aligned} \quad (A-3)$$

Differentiating Equation (A-3), the dimensionless velocity is determined.

$$\begin{aligned} \omega &= \text{CONJUGATE} \left| -i \left(1 + \frac{r^2}{s^2}\right) - ir \sum_{j=1}^k \lambda_j \left\{ \frac{1}{(s - s_j)} - \frac{1}{(s - r^2/s_j)} \right\} \right. \\ &\quad \left. + \frac{r}{\tan \alpha} \frac{dr}{dx} \frac{1}{s} \right| \end{aligned} \quad (A-4)$$

To simplify manipulations, the following functions are defined

$$a(s, s_j) = \left\{ \frac{1}{(s - s_j)} - \frac{1}{(s - r^2/\bar{s}_j)} \right\} \quad (A-5a)$$

$$b(s) = \left(\frac{1}{r} + \frac{r}{s^2} \right) + \sum_{j=3}^k \lambda_j a(s, s_j) \quad (A-5b)$$

Substituting Equations (A-5) into Equation (A-4):

$$\omega = -i r \text{ CONJUGATE} \left\{ b(s) + \lambda_1 a(s, s_1) + \lambda_2 a(s, s_2) + \frac{i}{\tan \alpha} \frac{dr}{dz} \frac{1}{s} \right\} \quad (A-6)$$

Substituting Equation (A-6) into constraints (A-2):

$$\begin{aligned} \lambda_1 a(s_{01}, s_1) + \lambda_2 a(s_{01}, s_2) &= -b(s_{01}) \\ \lambda_1 a(s_{02}, s_1) + \lambda_2 a(s_{02}, s_2) &= -b(s_{02}) \end{aligned} \quad (A-7)$$

Solving the above:

$$\begin{aligned} \lambda_1 &= \frac{b(s_{02})a(s_{01}, s_2) - b(s_{01})a(s_{02}, s_2)}{a(s_{01}, s_1)a(s_{02}, s_2) - a(s_{02}, s_1)a(s_{01}, s_2)} \equiv \frac{A1}{C} \\ \lambda_2 &= \frac{b(s_{01})a(s_{02}, s_1) - b(s_{02})a(s_{01}, s_1)}{a(s_{01}, s_1)a(s_{02}, s_2) - a(s_{02}, s_1)a(s_{01}, s_2)} \equiv \frac{B1}{C} \end{aligned} \quad (A-8)$$

Differentiating the above and substituting into the governing equations produces differential equations for z_j and y_j . The algebra involved is simplified by defining the following functions.

$$T(s, s_j) = \left[\frac{1}{(s - r^2/\bar{s}_j)^2} - \frac{1}{(s - s_j)^2} \right]$$

$$E(s, s_j) = (s - s_j)^{-2}$$

$$H(s, s_j) = \frac{-2r}{\bar{s}_j(s - r^2/\bar{s}_j)^2} \frac{dr}{dz}$$

$$X_F(s, s_j) = \frac{r^2}{\bar{s}_j^2(s - r^2/\bar{s}_j)^2}$$

$$G(\delta) = -\frac{2r}{\delta^3} + \sum_{j=3}^K \lambda_j T(\delta, \delta_j)$$

$$P(\delta) = \frac{dr}{dz} \left(\frac{1}{\delta^2} - \frac{1}{r^2} \right) + \sum_{j=3}^K H(\delta, \delta_j) \lambda_j$$

$$C1 = [-\alpha(\delta_{02}, \delta_2) E(\delta_{01}, \delta_1) + \alpha(\delta_{01}, \delta_2) E(\delta_{02}, \delta_1)] / C$$

$$C2 = [-\alpha(\delta_{02}, \delta_2) XF(\delta_{01}, \delta_1) + \alpha(\delta_{01}, \delta_2) XF(\delta_{02}, \delta_1)] / C$$

$$C3 = [-\alpha(\delta_{01}, \delta_1) E(\delta_{02}, \delta_2) + \alpha(\delta_{02}, \delta_1) E(\delta_{01}, \delta_2)] / C$$

$$C4 = [-\alpha(\delta_{01}, \delta_1) XF(\delta_{02}, \delta_2) + \alpha(\delta_{02}, \delta_1) XF(\delta_{01}, \delta_2)] / C$$

$$C5 = [-\alpha(\delta_{02}, \delta_2) T(\delta_{01}, \delta_1) + \alpha(\delta_{01}, \delta_2) T(\delta_{02}, \delta_1)] / C$$

$$C6 = [-\alpha(\delta_{01}, \delta_1) T(\delta_{02}, \delta_2) + \alpha(\delta_{02}, \delta_1) T(\delta_{01}, \delta_2)] / C$$

$$C7 = [\alpha(\delta_{02}, \delta_1) H(\delta_{01}, \delta_2) + \alpha(\delta_{01}, \delta_2) H(\delta_{02}, \delta_1) \\ - \alpha(\delta_{01}, \delta_1) H(\delta_{02}, \delta_2) - \alpha(\delta_{02}, \delta_1) H(\delta_{01}, \delta_2)] / C$$

$$K11 = C1(\delta_1 - \delta_{01}) + 1.$$

$$K12 = C2(\delta_1 - \delta_{01})$$

$$K13 = \left[C3 - \frac{E(\delta_{02}, \delta_2) b(\delta_{01})}{A1} + \frac{b(\delta_{02}) E(\delta_{01}, \delta_2)}{A1} \right] (\delta_1 - \delta_{01})$$

$$K14 = \left[C4 - \frac{XF(\delta_{02}, \delta_2) b(\delta_{01})}{A1} + \frac{b(\delta_{02}) E(\delta_{01}, \delta_2)}{A1} \right] (\delta_1 - \delta_{01})$$

$$K15 = \left[C5 - \frac{\alpha(\delta_{02}, \delta_2) G(\delta_{01})}{A1} + \frac{b(\delta_{02}) T(\delta_{01}, \delta_2)}{A1} \right] (\delta_1 - \delta_{01})$$

$$K16 = \left[C6 - \alpha(\delta_{01}, \delta_2) G(\delta_{02}) + \frac{T(\delta_{02}, \delta_2) B(\delta_{01})}{A1} \right] (\delta_1 - \delta_{01})$$

$$K L(\delta_j) = [-\alpha(\delta_{02}, \delta_2) E(\delta_{01}, \delta_j) + \alpha(\delta_{01}, \delta_2) E(\delta_{02}, \delta_j)] (\delta_1 - \delta_{01}) / A1$$

$$K M(\delta_j) = [-\alpha(\delta_{02}, \delta_2) XF(\delta_{01}, \delta_j) + \alpha(\delta_{01}, \delta_2) XF(\delta_{02}, \delta_j)] (\delta_1 - \delta_{01}) / A1$$

$$K19 = [C7 - \{b(S_{01})H(S_{02}, S_2) + a(S_{02}, S_2)P(S_{01}) - b(S_{02})H(S_{01}, S_2) - a(S_{01}, S_2)P(S_{02})\} / B1] (S_1 - S_{01})$$

$$K21 = [C1 - \frac{b(S_{02})E(S_{01}, S_1) + b(S_{01})E(S_{02}, S_1)}{B1}] (S_2 - S_{02})$$

$$K22 = [C2 - \frac{b(S_{02})XF(S_{01}, S_1)}{B1} + \frac{b(S_{01})XF(S_{02}, S_1)}{B1}] (S_2 - S_{02})$$

$$K23 = C3 (S_2 - S_{02}) + 1.$$

$$K24 = C4 (S_2 - S_{02})$$

$$K25 = [C5 - \frac{b(S_{02})T(S_{01}, S_2)}{B1} + \frac{a(S_{02}, S_1)G(S_{01})}{B1}] (S_2 - S_{02})$$

$$K26 = [C6 - \frac{b(S_{01})T(S_{02}, S_1)}{B1} + \frac{a(S_{01}, S_2)G(S_{02})}{B1}] (S_2 - S_{02})$$

$$KN(S_j) = [-E(S_{02}, S_j) a(S_{01}, S_j) + E(S_{01}, S_j) a(S_{02}, S_1)] (S_2 - S_{02}) / B1$$

$$KO(S_j) = [-XF(S_{02}, S_j) a(S_{01}, S_j) + XF(S_{01}, S_j) a(S_{02}, S_1)] (S_2 - S_{02}) / B1$$

$$K29 = [C7 - \{b(S_{02})H(S_{01}, S_1) + a(S_{01}, S_1)P(S_{02}) - b(S_{01})H(S_{02}, S_1) - a(S_{02}, S_1)P(S_{01})\} / B1] (S_2 - S_{02})$$

Using the above, explicit expressions for all derivatives can be obtained in the following form:

$$\frac{dy_j}{dx} = \text{Real} [\omega_j] \quad (A-9a)$$

$j = 3, 4 \dots k$

$$\frac{dz_j}{dx} = \text{Imag.} [\omega_j] \quad (A-9b)$$

$$\left\{ \begin{array}{l} \frac{du_1}{dx} \\ \frac{dz_1}{dx} \\ \frac{dy_2}{dx} \\ \frac{dz_2}{dx} \end{array} \right\} = [D]^{-1} \{Q\} \quad (A-9c)$$

The matrices D and Q are defined as follows:

$$D(1,1) = \text{Real}[K_{11} + K_{12}]$$

$$D(1,2) = \text{Imag.}[K_{12} - K_{11}]$$

$$D(1,3) = \text{Real}[K_{13} + K_{14}]$$

$$D(1,4) = \text{Imag.}[K_{14} - K_{13}]$$

$$D(2,1) = \text{Real}[K_{11} + K_{12}]$$

$$D(2,2) = \text{Imag.}[K_{11} - K_{12}]$$

$$D(2,3) = \text{Real}[K_{13} + K_{14}]$$

$$D(2,4) = \text{Imag.}[K_{13} - K_{14}]$$

$$D(3,1) = \text{Real}[K_{21} + K_{22}]$$

$$D(3,2) = \text{Imag.}[K_{22} - K_{21}]$$

$$D(3,3) = \text{Real}[K_{23} + K_{24}]$$

$$D(3,4) = \text{Imag.}[K_{24} - K_{23}]$$

$$D(4,1) = \text{Real}[K_{21} + K_{22}]$$

$$D(4,2) = \text{Imag.}[K_{21} - K_{22}]$$

$$D(4,3) = \text{Real}[K_{23} + K_{24}]$$

$$D(4,4) = \text{Imag.}[K_{23} - K_{24}]$$

$$Q(1) = \text{Real}(L_1)$$

$$Q(2) = \text{Imag}(L_1)$$

$$Q(3) = \text{Real}(L_2)$$

$$Q(4) = \text{Imag}(L_2)$$

APPENDIX B

COMPUTER PROGRAM FOR A BODY OF REVOLUTION AT
ANGLE OF ATTACK UP TO 50 DEGREES

The program listing and a sample run are included in this section. The program was developed on the CDC 6400. Running time is on the order of a minute. The sample run is for a sharp nose tangent ogive body of revolution at an angle of attack of 45 degrees.

The program is written to describe an ogive body of revolution with a sharp or blunt nose. If a sharp nose is to be specified, variables listed by the caption "nose tip dimensions" can be neglected. Allowable geometries for the blunted nose tip are restricted to ogive shapes of which a circle is a special case.

To use the program the following input is required.

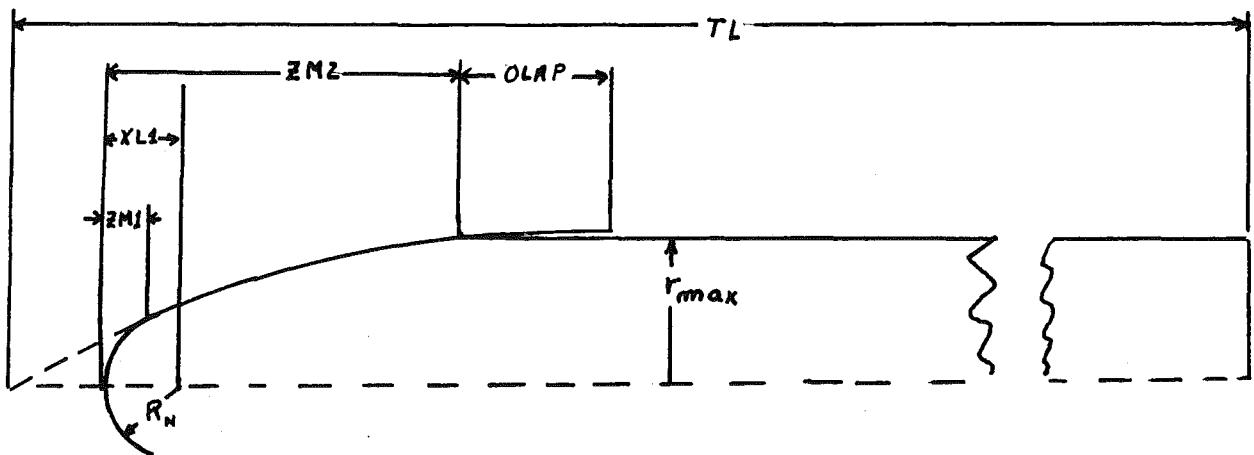
<u>Card</u>	<u>Variables</u>	<u>Format</u>
1	N	I5
2	A	F10.0
3	XL, r_{max} , ZM2, TL, OLAP, XLL, R_N , ZML	7F10.0
4	XTR, AKT, AKL, λ_I , XST, S, $\Delta r_p/D$	7F10.0
5	M, NE, NGB, G, Y(N+2)	3I5, 2F10.0
6	$\Delta r_s/D$, GS, r_n/D , $\Delta \lambda_n$, NNE	4F10.0, I5

where:

N	- number of cases to be run	
A	- angle of attack (degrees)	
XL		
r_{max}	ogive dimensions	
ZM2		
TL	total missile length	
OLAP		these parameters are defined in the accom- panying figure
XLL		
R_N	nose tip dimensions	
ZML		

XTR	fractional distance from missile nose to the point where transition from one separation angle to another takes place
AKT	final separation angle (radians)
AKL	initial separation angle (radians)
λ_I	initial vortex strength. If not specified, the value of .04 is assumed.
XST	nearest point to missile nose that integration is allowed to start. If not specified, the value of .1D is assumed.
S	Strouhal number
Δr_p	initial radial perturbation
M	number of integration steps between printing
NE	error bound. Step size is adjusted to keep estimated error between 10^{-NE} and $10^{-(NE+3)}$.
NGB	NGB-1 is the number of consecutive steps to be printed at the start or restart of integration.
G	initial step size
$Y(N+2)$	applies when M=0. It is the approximate distance in x between printing.
Δr_s	first vortex is shed when the radial location of the first two vortices, multiplied by r_{max} , differs by more than this amount.
r_n	this parameter, multiplied by r , represents the distance between the flow separation point, ζ_0 , and a newly introduced vortex. This distance is taken along the feeding sheet of the shed vortex.
$\Delta\lambda$	percent decrease in the circulation of a shed vortex
NNE	error bound in effect after first vortex is shed

Values for the parameters listed on this page which are used in the current study are found in the right-hand column of Table 1.



The symbols on the output page have the following meanings.

x distance along body axis

R local body radius

THETA0 θ_{01}, θ_{02}

VORTEX NUMBER - identification number for vortex

LAM circulation strength, λ

RADIUS } polar coordinates of vortex

ANGLE } see accompanying figure

Y } cartesian coordinates of vortex

Z } see accompanying figure

CL linear portion of normal force

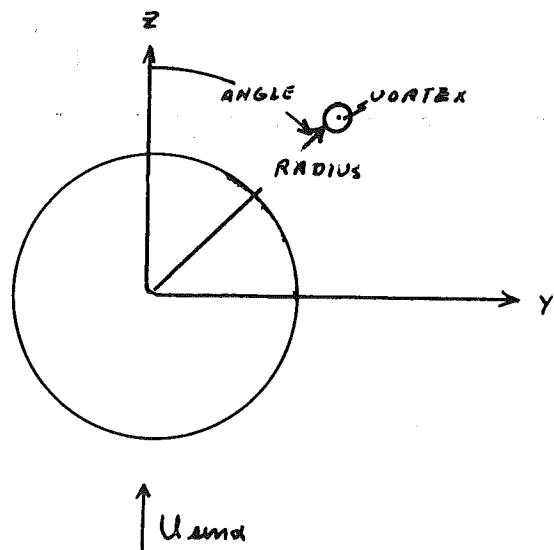
CN nonlinear portion of normal force

CML linear portion of pitching moment

CMN nonlinear portion of pitching moment

$$\left. \begin{array}{l} C_N = CL + CN \\ C_M = CML + CMN \end{array} \right\}$$

CD local normal force coefficient, $\frac{C_{DC}}{dx}$
CY yaw force coefficient, C_Y
CYM yawing moment coefficient, C_M'



PROGRAM VORT TRACE CDC 6400 FTN V3.0-P316 OPT=0 01/07/74 14.28.01. PAGE 1

	PROGRAM VORT(INPUT,OUTPUT,TAPE5=INPUT,TAPE6=OUTPUT)	VORT	2
	C.....MAIN PROGRAM	VORT	3
	COMPLEX XC	VORT	4
5	DIMENSION Y(35),D(35),XL(15),XC(15),ERROR(31)	VORT	5
	COMMON/GAMM/WRAD,ZM,FLEN,CR,TL,TSL,WRAD2,ZM2,FLEN2,CR2,OLAP	VORT	6
	COMMON/BG/N,DRT,PRAD,PRTU,TOP,XLP,PERR,STR,SHED,XAKT,XAKL,XXTR	VORT	7
	COMMON/GARB/XL,R,DR,XC,TAL,FVORT	VORT	8
	COMMON/STA/A,ANGB	VORT	9
10	COMMON/VEL/U,IE	VORT	10
	COMMON/XLIF/MC,M,CDMN,CDML,XS,CYM,CDCL,CDCN,CY	VORT	11
	COMMON/TRAN/XTR,X,AKT,AKL,ISD	VORT	12
	EXTERNAL DERIV,OUT,TERM2,TERM3	VORT	13
	CONVR=57.2957795	VORT	14
	READ(5,1005)NDS	VORT	15
15	DO 21 KT=1,NDS	VORT	16
	READ (5,1000)A	VORT	17
	WRITE(6,2000)A	VORT	18
	ZM2=0.0	VORT	19
20	READ(5,1000)FLEN,WRAD,ZM,TL,OLAP,FLEN2,WRAD2,ZM2	VORT	20
	A=A/CONVR	VORT	21
	TAL=TAN(A)	VORT	22
	ISEC=0	VORT	23
	IF(FLEN2.GT.0.0)ISEC=1	VORT	24
	WRITE(6,2001)WRAD,FLEN,ZM,TL	VORT	25
25	CR2=(WRAD2*WRAD2+FLEN2*FLEN2)/2./WRAD2	VORT	26
	IF(ISEC.EQ.1)WRITE(6,2016)WRAD2,ZM2,FLEN2	VORT	27
	CR=(WRAD*WRAD+FLEN*FLEN)/2./WRAD	VORT	28
	READ(5,1000)XTR,AKT,AKL,XLP,XST,STR,PERR	VORT	29
	XXTR=XTR	VORT	30
30	XAKT=AKT	VORT	31
	XAKL=AKL	VORT	32
	ISD=0.0	VORT	33
	IF(AKT.EQ.0.0)AKT=.9251	VORT	34
	IF(AKL.EQ.0.0)AKL=13.12	VORT	35
35	IF(AKL.LT.1.6) ISD=1	VORT	36
	IF(XLP.LE.0.0)XLP=.04	VORT	37
	IF(XST.LE.0)XST=.2*WRAD	VORT	38
	J=2	VORT	39
	IE=1	VORT	40
40	E1=.000001	VORT	41
	E2=E1	VORT	42
	DT=.04	VORT	43
	C.....DEFINING COMMON QUANTITIES	VORT	44
	ICOUNT=0	VORT	45
45	SHED=0.0	VORT	46
	IF(STR.NE.0)SHED=WRAD/STR/TAL	VORT	47
	U=1.	VORT	48
	N=5	VORT	49
	FVORT=0.0	VORT	50
	TOP=0.0	VORT	51
50	DO 10 K=1,15	VORT	52
	10 XL(K)=0.0	VORT	53
	C.....DETERMINE INITIAL INTEGRATION PARAMETERS	VORT	54
	CALL START(A,DD1,DT,E1,E2,XLP,XST)	VORT	55
55	C.....SET UP INTEGRATION	VORT	56

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NOLTR 73-209

PROGRAM VORT TRACE CDC 6400 FTN V3.0-P316 OPT=0 01/07/74 14.28.01. PAGE 2

	CALL PROFILE(R,DR,X,DD1,1)	VORT	57
	Y(1)=E1*R*COS(E2)	VORT	58
	Y(2)=E1*R*SIN(E2)	VORT	59
	Y(3)=-Y(1)	VORT	60
60	Y(4)=Y(2)	VORT	61
	Y(1)=Y(1)*(1.+PERR)	VORT	62
	Y(2)=Y(2)*(1.+PERR)	VORT	63
	Y(3)=Y(3)*(1.-PERR)	VORT	64
	Y(4)=Y(4)*(1.-PERR)	VORT	65
65	Y(5)=THETI(DD1)	VORT	66
	XS=X	VORT	67
	Y(N+1)=X	VORT	68
	Y(N+3)=-1.	VORT	69
	READ(5,1005)M,NE,NGB,G,Y(N+2)	VORT	70
70	MC=M	VORT	71
	M=1	VORT	72
	CALL REST(X,Y,J,G,L,NE,D,0,MC)	VORT	73
	CYM=0.0	VORT	74
	CY =0	VORT	75
75	CDCN=SIN(A)*COS(A)*8.*R*Y(1)*XLP*(1.-R*R/(Y(1)*Y(1)+Y(2)*Y(2)))	VORT	76
	1/WRAD/WRAD	VORT	77
	CDMN=CDCN*X/2./WRAD	VORT	78
	CDCL=2.*SIN(A)*COS(A)*R*R/WRAD/WRAD	VORT	79
	X0=CR*CR-FLEN*FLEN	VORT	80
80	X2=(X0+2.*FLEN*X-X*X)**.5	VORT	81
	X0=X0**.5	VORT	82
	CDML=2.*SIN(A)*COS(A)*((WRAD-CR)*((X+FLEN)*X2/2.-FLEN/2.*X0-CR*CR/	VORT	83
	12.*(ASIN((X-FLEN)/CR)-ASIN(-FLEN/CR)))+X*X*(FLEN/2.-X/3.))/WRAD/WR	VORT	84
	2AD/WRAD	VORT	85
85	TOP=0.0	VORT	86
	FVORT=0.0	VORT	87
18	CONTINUE	VORT	88
	CALL FNOL3(J,N,G,L,1,NE,X,Y,D,DERIV,TERM2,OUT,ERROR)	VORT	89
	M=1	VORT	90
90	CALL OUT(X,Y,D,ERROR,N,L,G)	VORT	91
	IF(X+.001.GE.TL)GO TO 20	VORT	92
	C..... INTRODUCING FIRST NEW VORTEX	VORT	93
	N=7	VORT	94
	IC=3	VORT	95
95	IF(Y(1)**2+Y(2)**2.GT.Y(3)**2+Y(4)**2)GO TO 25	VORT	96
	WRITE(6,2003)	VORT	97
	STOP	VORT	98
25	CONTINUE	VORT	99
	Y(10)=Y(8)	VORT	100
100	Y(9)=Y(7)	VORT	101
	Y(8)=Y(6)	VORT	102
	Y(7)=Y(5)	VORT	103
	Y(6)=Y(2)	VORT	104
	Y(5)=Y(1)	VORT	105
105	XL(3)=XL(1)*(1.+PRTU)	VORT	106
40	CALL REST(X,Y,J,G,L,NE,D,2,MC)	VORT	107
	TSL=X+SHED	VORT	108
	M=1	VORT	109
	CALL FNOL3(J,N,G,L,1,NE,X,Y,D,DERIV,TERM3,OUT,ERROR)	VORT	110
110	M=1	VORT	111

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NOLTR 74-209

PROGRAM	VORT	TRACE	CDC 6400 FTN V3.0-P316 OPT=0	01/07/74	14.28.01.	PAGE	3
		CALL OUT(X,Y,D,ERROR,N,L,H)			VORT	112	
		IF(X+.001.GE.TL)GO TO 20			VORT	113	
		C.....NEW VORTEX IS ADDED			VORT	114	
		N=N+2			VORT	115	
115		IF(N.GT.31)GO TO 20			VORT	116	
		Y(N+3)=Y(N+1)			VORT	117	
		Y(N+2)=Y(N)			VORT	118	
		Y(N+1)=Y(N-1)			VORT	119	
		Y(N)=Y(N-2)			VORT	120	
120		IC=(N-1)/2			VORT	121	
		IC2=IC-1			VORT	122	
		DO 50 IN=1,IC2			VORT	123	
		K=2*(IC-IN)			VORT	124	
		Y(K+2)=Y(K)			VORT	125	
125		Y(K+1)=Y(K-1)			VORT	126	
	50	XL(IC+1-IN)=XL(IC-IN)			VORT	127	
		XL(3)=XL(3)*(1.+PRTU)			VORT	128	
		ISTOP=1			VORT	129	
		ICOUNT=0			VORT	130	
130		GO TO 40			VORT	131	
	20	CNN=CDCL+CDCN			VORT	132	
		CMM=CDMN+CDML			VORT	133	
		CSP=CMM/CNN			VORT	134	
		CYP=0.0			VORT	135	
135		IF(DRT.NE.0.0)CYP=CYM/CY			VORT	136	
		WRITE(6,2004)			VORT	137	
		WRITE(6,2005)CNN			VORT	138	
		WRITE(6,2006)CMM			VORT	139	
		WRITE(6,2007)CSP			VORT	140	
140		WRITE(6,2008)CY			VORT	141	
		WRITE(6,2009)CYM			VORT	142	
		WRITE(6,2010)CYP			VORT	143	
	21	CONTINUE			VORT	144	
		STOP			VORT	145	
145	1000	FORMAT(8F10.0)			VORT	146	
	1001	FORMAT(6F10.0,I5)			VORT	147	
	1005	FORMAT(3I5,2F10.0)			VORT	148	
	2000	FORMAT(1H1,40X,16HANGLE OF ATTACK ,F5.2,2X,7HDEGREES)			VORT	149	
	2001	FORMAT(1H0,46X,16HOGIVE DIMENSIONS/1H ,6X,14HMAXIMUM RADIUS,6X,			VORT	150	
150		113H OGIVE LENGTH,18X,16HZ MATCHING POINT,5X,11HBODY LENGTH/1H ,			VORT	151	
		26X,F10.5,14X,F6.3,22X,F10.5,13X,F10.5)			VORT	152	
	2003	FORMAT(27H0VORTEX IDENTITIES REVERSED)			VORT	153	
	2004	FORMAT(1H0,40X,26H*****FINAL STATISTICS*****)			VORT	154	
	2005	FORMAT(1H ,37X,32HNORMAL FORCE COEFFICIENT=CL+CN= ,F10.5)			VORT	155	
155	2006	FORMAT(1H ,35X,37HPITCHING MOMENT COEFFICIENT=CML+CMN= ,F10.5)			VORT	156	
	2007	FORMAT(1H ,41X,23HCENTER OF NORMAL FORCE=,F10.5)			VORT	157	
	2008	FORMAT(1H ,42X,22HYAW FORCE COEFFICIENT=,F10.5)			VORT	158	
	2009	FORMAT(1H ,41X,23HYAW MOMENT COEFFICIENT=,F10.5)			VORT	159	
	2010	FORMAT(1H ,43X,20HCENTER OF YAW FORCE=,F15.5)			VORT	160	
160	2016	FORMAT(1H0,46X,15HNOSE DIMENSIONS/1H ,28X,6HRADIUS,5X, 14HMATCHIN			VORT	161	
		1G POINT,5X,6HLENGTH/1H ,35X,3(F12.7,3X))			VORT	162	
		END			VORT	163	

SUBROUTINE REST TRACE CDC 6400 FTN V3.0-P316 OPT=0 01/07/74 14.28.01. PAGE 1

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SUBROUTINE REST(X,Y,J,G,L,NE,D,IT,M) REST 2
COMMON/GARB/XL,R,DR,XC,TAL,FVORT REST 3
COMMON/BG/N,DRT,PRAD,PRTU,TOP,XLP,PERR,S,SHED,AKT,AKL,XTR REST 4
COMMON/GAMM/WRAD REST 5
5 DIMENSION Y(1),D(1),XL(15),XLR(15),XC(15) REST 6
COMPLEX XC REST 7
COMMON/STA/A REST 8
CONVR=57.2957795 REST 9
IC=(N-1)/2 REST 10
10 IF(IT.EQ.0)GO TO 10 REST 11
NE=NNE REST 12
G=XG REST 13
ANG=Y(N) REST 14
IF((IC/2)*2.EQ.IC)ANG=3.1415926536-ANG REST 15
15 P1=COS(ANG)*R REST 16
P2=SIN(ANG)*R REST 17
PAN=ABS((Y(6)-P2)/(Y(5)-P1)) REST 18
PAN=ATAN(PAN) REST 19
IF(Y(5).LT.P1)PAN=3.1415926536-PAN REST 20
20 Y(1)=P1+COS(PAN)*PRAD*R REST 21
Y(2)=P2+SIN(PAN)*PRAD*R REST 22
RETURN REST 23
10 CONTINUE REST 24
25 READ(5,1001)DRT,XG,PRAD,PRTU,NNE REST 25
WRITE(6,2000)XLP,PERR,DRT,PRTU,PRAD,S,SHED,AKL,AKT,XTR REST 26
PRTU=PRTU/100. REST 27
WRITE(6,2005)G,M,NE,Y(N+2),NNE REST 28
DRT=DRT*2.*WRAD REST 29
RETURN REST 30
30 1001 FORMAT(4F10.0,I5) REST 31
2000 FORMAT(1H0,40X,25H*****FREE PARAMETERS*****/1H ,40X,24HINITIAL VOR REST
1TEX STRENGTH ,F10.5/1H ,38X,28HINITIAL RADIAL PERTURBATION ,F10.5/ REST 32
21H ,30X,47HRADIAL ASYMMETRY AT WHICH FIRST VORTEX IS SHED ,F10.5/ REST 33
31H ,28X,50HPER CENT DECREASE IN CIRCULATION OF SHED VORTICES ,F10. REST 34
35 45/1H ,25X,46HLOCATION FACTOR FOR NEWLY INTRODUCED VORTICES ,F10.5/ REST 35
51H ,18X,16HSTROUHAL NUMBER ,F10.5,2X,32HCORRESPONDING SHEDDING DIS REST 36
6TANCE ,F10.5/1H ,12X,27HSEPARATION ANGLE CONSTANTS ,F10.5,2X,F10.5 REST 37
7,2X,20HPOINT OF TRANSITION ,F10.5) REST 38
2005 FORMAT(1H0,42X,19HINTEGRATION OPTIONS/1H ,18HINITIAL STEP SIZE ,F1 REST 39
40 10.7,5X,15HPRINT FREQUENCY, I5,5X,12HERROR BOUND ,I5,2X,16HPRINT IN REST 40
2CREMENT ,F10.7/1H ,42X,34HERROR BOUND AFTER VORTEX SHEDDING ,I5) REST 41
END REST 42
43

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FUNCTION	THETI	TRACE	CDC 6400 FTN V3.0-P316 OPT=0	01/07/74	14.28.01.	PAGE	1
	FUNCTION THETI(DELTA)			THETI	2		
	COMMON/STA/AL			THETI	3		
	COMMON/GAMM/WRAD,ZM,FLEN,CR,TL			THETI	4		
	COMMON/TRAN/XTR,X,AKT,AKL,ISD			THETI	5		
5	CONVR=57.2957795			THETI	6		
	IF(X/TL.LT.XTR)GO TO 11			THETI	7		
	THETI=AKT			THETI	8		
	RETURN			THETI	9		
10	11 CONTINUE			THETI	10		
	IF(ISD.EQ.0)GO TO 12			THETI	11		
	THETI=AKL			THETI	12		
	RETURN			THETI	13		
15	12 CONTINUE			THETI	14		
	THETI=(AKL+CONVR*DELTA)*(3.13-(.116*AL*CONVR-1.16)**.5)			THETI	15		
	THETI=THETI/CONVR			THETI	16		
	RETURN			THETI	17		
20	10 THETI=3.13*(AKL+CONVR*DELTA)			THETI	18		
	THETI=THETI/CONVR			THETI	19		
	RETURN			THETI	20		
	END			THETI	21		

SUBROUTINE PROFILE TRACE

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        SUBROUTINE PROFILE(R,DR,Z,DEL,IC)
        COMMON/GAMM/R02,ZM2,XL2,CR2,TL,TSL,R01,ZM1,XL1,CR1,OLAP
        IF(IC.EQ.1)GO TO 4
6  CONTINUE
5    RO=R02
     CR=CR2
     ZM=ZM2+OLAP
     IF(Z.GT.ZM1)GO TO 3
     RO=R01
     CR=CR1
     ZM=XL1
     GO TO 3
4    SF=TAN(DEL)
     IF(ZM1.EQ.0.0)GO TO 5
15   Z=XL1-CR1*SF/((SF*SF+1.)**.5)
     IF(Z.LT.ZM1)GO TO 6
5    Z=ZM2+OLAP-CR2*SF/((SF*SF+1.)**.5)
     IF(Z.LT.ZM1)Z=ZM1+.001
     GO TO 6
20   3  IF(Z.LT.ZM2)GO TO 20
     DES=CR*CR-(ZM-ZM2)**2
     R=RO-CR+DES**.5
     DR=0.0
     IF(IC.EQ.1)Z=ZM2
     RETURN
25   20 DES=CR*CR-(ZM-Z)**2
     IF(DES.GE.0.0)GO TO 30
     WRITE(6,200)RO,ZM,Z,DEL,XL1,CR1
     STOP
30   30 DES=DES**.5
     R=RO-CR+DES
     DR=(ZM-Z)/DES
     RETURN
35   200 FORMAT(1H ,25HPROBABLE INPUT DATA ERROR,6E16.8)
     END

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	PROFILE	2
	PROFILE	3
	PROFILE	4
	PROFILE	5
	PROFILE	6
	PROFILE	7
	PROFILE	8
	PROFILE	9
	PROFILE	10
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	PROFILE	12
	PROFILE	13
	PROFILE	14
	PROFILE	15
	PROFILE	16
	PROFILE	17
	PROFILE	18
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	PROFILE	29
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	PROFILE	31
	PROFILE	32
	PROFILE	33
	PROFILE	34
	PROFILE	35
	PROFILE	36

SUBROUTINE	START	TRACE	CDC 6400 FTN V3.0-P316 OPT=0	01/07/74	14.28.01.	PAGE	1
		SUBROUTINE START(AL,DD1,DT,E1,E2,XLMS,XST)	START	2			
		COMMON/GAMM/WRAD,ZM,FLEN,CR,TL	START	3			
		COMMON/TRAN/XTR,X,AKT,AKL,ISH	START	4			
		C.....INTERNAL FUNCTIONS	START	5			
5		XR(R1)=((R1**.5-1./(R1**.5))**2.)/4.	START	6			
		XL(R1,T1,T0)=8.* (XR(R1)+(SIN((T1-T0)/2.))**2)*(XR(R1)	START	7			
		1+(COS((T1+T0)/2.))**2) /COS(T1)/(R1-1./R1)	START	8			
		EN1(D)=(AKL/CONVR+D)*XK2-ASIN(XK1*TAN(D))	START	9			
		CONVR=57.2957795	START	10			
10		C.....DETERMINE DELTA CRITICAL	START	11			
		TAL=TAN(AL)	START	12			
		IF(ISH.EQ.0)GO TO 3	START	13			
		D=ATAN(TAL*SIN(AKL)/1.5)	START	14			
		GO TO 10	START	15			
15		3 CONTINUE	START	16			
		XK1=1.5/TAN(AL)	START	17			
		XK2=3.13	START	18			
		IF(AL.GT..17453)XK2=3.13-(.116*AL*CONVR-1.16)**.5	START	19			
		D1=ATAN(1./XK1)	START	20			
20		IF((AKL+D1*CONVR)*XK2.LT.90.)GO TO 4	START	21			
		WRITE(6,2004)	START	22			
	2004	FORMAT(21HONO STARTING SOLUTION)	START	23			
		STOP	START	24			
25		4 CONTINUE	START	25			
		DD1=D1/3.	START	26			
		D=D1-DD1/2.	START	27			
		DP=D	START	28			
		IC=0	START	29			
30		30 ER1=EN1(D)	START	30			
		IF(ABS(ER1).LT.E1)GO TO 10	START	31			
		IC=IC+1	START	32			
		IF(IC.NE.1)GO TO 20	START	33			
		ER2=ER1	START	34			
		DP=D	START	35			
35		D=D-DD1	START	36			
		GO TO 30	START	37			
	20	DH=D	START	38			
		D=D-ER1*(DP-D)/(ER2-ER1)	START	39			
		IF(D.GT.D1)D=(D1+DH)/2.	START	40			
40		DP=DH	START	41			
		ER2=ER1	START	42			
		GO TO 30	START	43			
	10	CONTINUE	START	44			
45		C.....TRIAL AND ERROR CALCULATION FOR R1	START	45			
		CALL PROFILE(R,DR,X,D,1)	START	46			
		IF(X.LT.0.0)X=0.0	START	47			
		IF(X/TL.LT.XTR)GO TO 12	START	48			
		D=ATAN(TAL*SIN(AKT)/1.5)	START	49			
		CALL PROFILE(R,DR,X,D,1)	START	50			
50		12 IF(X.GT.XST)GO TO 31	START	51			
		13 X=XST	START	52			
		CALL PROFILE(R,DR,X,DEL,2)	START	53			
		D=ATAN(DR)	START	54			
	31	DDEG=D*CONVR	START	55			
55		D1=D	START	56			

SUBROUTINE START TRACE

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	TO=THETI(D1)	START	57
	T1=TO+.001	START	58
	IT=0	START	59
	IP2=1	START	60
60	TP=TO+1.047198	START	61
	REX=SIN(TO)-TAN(TP)*COS(TO)	START	62
	TTP=TAN(TP)	START	63
	70 FE2=FE1	START	64
	R1=REX/(SIN(T1)-TTP*COS(T1))	START	65
65	FE1=ABS(XL(R1,T1,TO))-XLMS	START	66
	IF(ABS(FE1).LT.E2) GO TO 90	START	67
	IT=IT+1	START	68
	IF(IT.NE.1) GO TO 100	START	69
	T2=T1	START	70
70	T1=T1+DT	START	71
	GO TO 70	START	72
	100 TH=T1	START	73
	T1=T1-FE1*(T2-T1)/(FE2-FE1)	START	74
	IF(FE1*FE2.LT.0.0) IP2=0	START	75
75	IF(IP2.EQ.1) T1=TH+DT	START	76
	T2=TH	START	77
	GO TO 70	START	78
	90 E1=R1	START	79
	IF(T1.GT.1.570796327) T1=3.141592654-T1	START	80
80	E2=T1	START	81
	DD1=D1	START	82
	RETURN	START	83
	END	START	84

SUBROUTINE	FNOL3	TRACE	CDC 6400 FTN V3.0-P316 OPT=0	01/07/74	14.28.01.	PAGE	1
			SUBROUTINE FNOL3(J,NN,G,L,MPR,NE,X,Y,D,DERIV,TERM,OUTPUT,ERROR)				
	C 001	J=INTEGRATION METHOD CONTROL				FNOL3	2
	C 002	NN=NUMBER OF SIMULTANEOUS DIFFERENTIAL EQUATIONS				FNOL3	3
	C 003	G=FIRST INTERVAL OF INTEGRATION				FNOL3	4
5	C 004	L=NUMBER OF Y'S GREATER THAN N TO BE PRINTED				FNOL3	5
	C 005	MPR=PRINT FREQUENCY				FNOL3	6
	C 006	NE=CONTROL FOR INTERVAL OF INTEGRATION				FNOL3	7
	C 007	X=INDEPENDENT VARIABLE				FNOL3	8
10	C 008	Y=DEPENDENT VARIABLE				FNOL3	9
	C 009	D=ARRAY CONTAINING DERIVATIVES				FNOL3	10
	C 010	DERIV=SUBROUTINE IN WHICH DIFFERENTIAL EQUATIONS ARE FOUND				FNOL3	11
	C 011	TERM=SUBROUTINE FOR TERMINATION CONDITION				FNOL3	12
	C 012	OUTPUT=SUBROUTINE FOR PRINTING OUTPUT				FNOL3	13
		NRET=TERM LOOP COUNTER				FNOL3	14
15		NPR=COUNT OF STEPS SINCE LAST PRINT				FNOL3	15
		PC=Y(N+1)=PRINT CONTROL OTHER THAN STEP COUNT				FNOL3	16
		JJ=J-2=0 FOR AM,-1 FOR RK,+1 FOR RK2				FNOL3	17
		MK=AM RK STEP COUNT				FNOL3	18
		XD=DP FORM OF X				FNOL3	19
20		C DOUBLE PRECISION CARD HAS BEEN REMOVED IN ORDER TO RUN ON 6400				FNOL3	20
		C THIS DECK RUNS ON THE 7090 WITH OR WITHOUT DOUBLE PRECISION				FNOL3	21
		C DOUBLE PRECISION XD,YD,YP,YC,YK,H,HC,XS,XDS,HS,HN,HD24,HD6				FNOL3	22
		DIMENSION C(3),Y(183),YD(183),YP(183),YC(183),D(203),DM(183,4)				FNOL3	23
		DIMENSION DK(183,4),ERROR(183),YK(183)				FNOL3	24
25		DATA EP6,EP11,M4/1.E-6,1.E-11,-4/				FNOL3	25
		DATA (C(K),K=1,3)/2*.5,1./				FNOL3	26
		DATA HMAX5/1.E35/				FNOL3	27
		NHTS=0				FNOL3	28
		N=NN				FNOL3	29
30		JJ=J-2				FNOL3	30
		H=G				FNOL3	31
		HN=H				FNOL3	32
		MK=1				FNOL3	33
		NRET=M4				FNOL3	34
35		EC=Y(N+3)				FNOL3	35
		JTEST=1				FNOL3	36
		IF (JJ .LT. 0) JTEST=4				FNOL3	37
		IF (NE.EQ.0) GO TO 15				FNOL3	38
		EUP=10.*#(-NE)				FNOL3	39
40		ELO=EUP*.001				FNOL3	40
		EM=ELO*31.6227766				FNOL3	41
	15	XD=X				FNOL3	42
		DO 20 I=1,N				FNOL3	43
		ERROR(I)=0.				FNOL3	44
45	20	YD(I)=Y(I)				FNOL3	45
		CALL DERIV(X,Y,D)				FNOL3	46
		CALL TERM (X,Y,D,F)				FNOL3	47
	C	PRINT				FNOL3	48
	50	CALL OUTPUT(X,Y,D,ERROR,N,L,H)				FNOL3	49
50		IF (NRET) 65,60,55				FNOL3	50
	60	RETURN				FNOL3	51
	55	WRITE (6,3000) HN				FNOL3	52
	3000	FORMAT(108H1EXECUTION TERMINATED BECAUSE INTERVAL OF INTEGRATION L				FNOL3	53
		1ESS THAN 1.0E -6 TIMES INDEPENDENT VARIABLE (X). H =,1PF15.7)				FNOL3	54
55		STOP				FNOL3	55
						FNOL3	56

SUBROUTINE FNOL3 TRACE CDC 6400 FTN V3.0-P316 OPT=0 01/07/74 14.28.01. PAGE 2

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    65 IF (JTEST .EQ. 5 .AND. H .NE. HN) GO TO 455 FNOL3 57
      NPR=0 FNOL3 58
      PC=Y(N+1) FNOL3 59
    100 IF (JTEST .EQ. 5 .AND. H .NE. HN) GO TO 455 FNOL3 60
      IF (JJ .GE. 0) H=HN FNOL3 61
      IF (MK .NE. 0 .OR. JJ .NE. 0) GO TO 300 FNOL3 62
    C-----THE ADAMS MOULTON METHOD FNOL3 63
    200 HD24=H/24. FNOL3 64
      JAM=0 FNOL3 65
    65 DO 210 I=1,N FNOL3 66
      YPI=(55.*DM(I,1)+37.*DM(I,2))-(59.*DM(I,3)+9.*DM(I,4)) FNOL3 67
      YP(I)=YD(I)+HD24*YPI FNOL3 68
      Y(I)=YP(I) FNOL3 69
    210 CONTINUE FNOL3 70
      X=XD+H FNOL3 71
      CALL DERIV(X,Y,DM(1,4)) FNOL3 72
    70 DO 220 I=1,N FNOL3 73
      YPI=(9.*DM(I,4)+19.*DM(I,1)+DM(I,2))-5.*DM(I,3) FNOL3 74
      YC(I)=YD(I)+HD24*YPI FNOL3 75
    75 ERROR(I)=(YP(I)-YC(I))/14. FNOL3 76
      C THIS ADDS IN A 2D CORRECTION FNOL3 77
      YC(I)=YC(I)+ERROR(I) FNOL3 78
    220 CONTINUE FNOL3 79
      IF (NE.NE.0) GO TO 700 FNOL3 80
    80 GO TO 455 FNOL3 81
    C-----THE RUNGE KUTTA METHOD FNOL3 82
    300 GO TO (301,309,308,309,303),JTEST FNOL3 83
    301 DO 302 I=1,N FNOL3 84
      YK(I)=YD(I) FNOL3 85
    85 302 CONTINUE FNOL3 86
      XDS=XD FNOL3 87
      GO TO 309 FNOL3 88
    303 DO 304 I=1,N FNOL3 89
      YK(I)=YC(I) FNOL3 90
    90 304 CONTINUE FNOL3 91
      XDS=XD+H FNOL3 92
    308 HS=H FNOL3 93
      H=2.*H FNOL3 94
      GO TO 320 FNOL3 95
    95 309 X=XD FNOL3 96
      JAM=1 FNOL3 97
      DO 310 I=1,N FNOL3 98
      Y(I)=YD(I) FNOL3 99
      DK(I,1)=D(I) FNOL3 100
    100 310 CONTINUE FNOL3 101
      IF (JTEST .LE. 2) CALL DERIV(X,Y,DK) FNOL3 102
      IF (MK .GT. 1 .OR. JTEST .GT. 1) GO TO 320 FNOL3 103
      DO 315 I=1,N FNOL3 104
      DM(I,4)=DK(I,1) FNOL3 105
    105 315 CONTINUE FNOL3 106
    320 DO 335 K=2,4 FNOL3 107
      HC=H*C(K-1) FNOL3 108
      DO 330 I=1,N FNOL3 109
      Y(I)=YD(I) + HC*DK(I,K-1) FNOL3 110
    110 330 CONTINUE FNOL3 111
  
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SUBROUTINE	FNOL3	TRACE	CDC 6400 FTN V3.0-P316 OPT=0	01/07/74	14.28.01.	PAGE
		X=XD+HC			FNOL3	112
		CALL DERIV(X,Y,DK(1,K))			FNOL3	113
	335	CONTINUE			FNOL3	114
		HD6=H/6.			FNOL3	115
115		DO 340 I=1,N			FNOL3	116
		YPI=DK(I,1)+DK(I,4)+2.* (DK(I,2)+DK(I,3))			FNOL3	117
		YC(I)=YD(I)+HD6*YPI			FNOL3	118
	340	CONTINUE			FNOL3	119
		GO TO (360,390,370,455,370),JTEST			FNOL3	120
120	360	DO 365 I=1,N			FNOL3	121
		YP(I)=YC(I)			FNOL3	122
	365	CONTINUE			FNOL3	123
		JTEST=3			FNOL3	124
		GO TO 308			FNOL3	125
125	370	DO 380 I=1,N			FNOL3	126
		YD(I)=YP(I)			FNOL3	127
		YP(I)=YC(I)			FNOL3	128
	380	CONTINUE			FNOL3	129
		H=HS			FNOL3	130
130		XD=XD+H			FNOL3	131
		JTEST=2			FNOL3	132
		IF (MK .EQ. 1) GO TO 309			FNOL3	133
		GO TO 451			FNOL3	134
	390	DO 400 I=1,N			FNOL3	135
		ERROR(I)=(YC(I)-YP(I))/15.			FNOL3	136
		YC(I)=YC(I)+ERROR(I)			FNOL3	137
		YP(I)=YC(I)			FNOL3	138
	400	CONTINUE			FNOL3	139
		JTEST=5			FNOL3	140
140		IF (NE.NE.0) GO TO 700			FNOL3	141
		C-----ACCEPT THE STEP SIZE			FNOL3	142
	450	IF (JAM .EQ. 0) GO TO 455			FNOL3	143
		IF (MK .EQ. 3 .AND. JJ .EQ. 0) GO TO 455			FNOL3	144
		IF (MK .NE. 1) GO TO 303			FNOL3	145
145		IF (JTEST .EQ. 1) GO TO 455			FNOL3	146
	451	DO 452 I=1,N			FNOL3	147
		Y(I)=YD(I)			FNOL3	148
	452	CONTINUE			FNOL3	149
		GO TO 466			FNOL3	150
150	455	DO 459 NQ=1,N			FNOL3	151
		YD(NQ)=YC(NQ)			FNOL3	152
		Y(NQ)=YD(NQ)			FNOL3	153
	459	CONTINUE			FNOL3	154
		IF (JJ .GE. 0) JTEST=1			FNOL3	155
155		IF (MK .NE. 0 .OR. JJ .NE. 0 .OR. NRET .NE. M4) GO TO 465			FNOL3	156
		DO 460 I=1,N			FNOL3	157
		DM(I,4)=DM(I,2)			FNOL3	158
		DM(I,2)=DM(I,3)			FNOL3	159
		DM(I,3)=DM(I,1)			FNOL3	160
160		460 CONTINUE			FNOL3	161
	465	XD=XD+H			FNOL3	162
	466	X=XD			FNOL3	163
		IF (MK .EQ. 3) MK=0			FNOL3	164
		CALL DERIV(X,Y,D)			FNOL3	165
165		DO 470 I=1,N			FNOL3	166

SUBROUTINE FNOL3 TRACE CDC 6400 FTN V3.0-P316 OPT=0 01/07/74 14.28.01. PAGE 4

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        DM(I,MK+1)=D(I) FNOL3 167
470  CONTINUE FNOL3 168
      FP=F FNOL3 169
      CALL TERM (X,Y,D,F) FNOL3 170
170  C-----DO YOU TERMINATE FNOL3 171
      500 IF (ABS(F) .LE. EP6) GO TO 800 FNOL3 172
          IF (FP .EQ. 0.) GO TO 510 FNOL3 173
          IF (NRET .NE. M4 .OR. F*FP .LT. EP11) GO TO 805 FNOL3 174
      510 XS=XD FNOL3 175
175  IF (MK .NE. 0 .AND. H .EQ. HN) MK=MK+1 FNOL3 176
C-----DO YOU PRINT FNOL3 177
      600 NPR=NPR+1 FNOL3 178
          IF (MPR .GT. 0 .AND. NPR .GE. MPR) GO TO 50 FNOL3 179
          IF (MPR .LE. 0 .AND. ABS(Y(N+1)-PC) .GE. Y(N+2)) GO TO 50 FNOL3 180
      GO TO 100 FNOL3 181
180  C-----DETERMINING THE STEP SIZE FNOL3 182
      700 HB = HMAX5 FNOL3 183
          DO 760 I = 1,N FNOL3 184
          Z=ABS(ERROR(I)) FNOL3 185
185  IF (YC(I) .EQ. 0.) GO TO 720 FNOL3 186
          ZZ= ABS(YC(I)) FNOL3 187
          IF (EC) 720,710,705 FNOL3 188
      705 IF (EC .GT. ZZ) ZZ=EC FNOL3 189
      710 Z=Z/ZZ FNOL3 190
190  720 IF (Z.GT.ELO.AND.Z.LT.EUP) GOTO 750 FNOL3 191
          HB = AMIN1(HB,EM/(Z+EP11)) FNOL3 192
          GOT0760 FNOL3 193
      750 HB=AMIN1(HB,1.) FNOL3 194
195  760 CONTINUE FNOL3 195
          IF (HB.EQ.1.) GOTO 450 FNOL3 196
          HN=H*HB**.2 FNOL3 197
          IF (MK .NE. 1) JTEST=1 FNOL3 198
          MK=1 FNOL3 199
200  IF (HB.LT.1.) GOTO 770 FNOL3 200
          IF (ABS(HN) .GT. ABS(4.*H)) HN=4.*H FNOL3 201
          NHTS=0 FNOL3 202
          GOTO 450 FNOL3 203
      770 HEPS=ABS(X*EP6) + EP11 FNOL3 204
          IF (ABS(HN) .LT. ABS(H/4.)) HN=H/4. FNOL3 205
          IF (ABS(HN) .GT. HEPS) GO TO 790 FNOL3 206
          NHTS = NHTS + 1 FNOL3 207
          IF (NHTS.LT.10) GOTO 780 FNOL3 208
          JTEST=1 FNOL3 209
          NRET = 1 FNOL3 210
          GOTO 50 FNOL3 211
      780 HN=SIGN(HEPS,HN) FNOL3 212
      790 IF (JAM .EQ. 0) GO TO 100 FNOL3 213
          DO 795 I=1,N FNOL3 214
          YD(I)=YK(I) FNOL3 215
215  795 CONTINUE FNOL3 216
          XD=XDS FNOL3 217
          JTEST=1 FNOL3 218
          GO TO 100 FNOL3 219
C-----THE TERMINATION LOOP FNOL3 220
      800 NRET=0 FNOL3 221

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SUBROUTINE FNOL3 TRACE

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JTEST1=1
805 IF (NRET .NE. 0) GO TO 806
H=XD-XS
GO TO 50
225 806 IF (F*FP.LT.0.) GOTO 810
IF (F*FP2.LT.0.) GOTO 820
GO TO 800
810 FP2 =FP
HP =H
230 820 FP =FP2
HP =H + HP
830 NRET=NRET+1
H=HP*F/(FP-F)
235 830 JTEST=4
GOTO 300
END
```

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FNOL3 222
FNOL3 223
FNOL3 224
FNOL3 225
FNOL3 226
FNOL3 227
FNOL3 228
FNOL3 229
FNOL3 230
FNOL3 231
FNOL3 232
FNOL3 233
FNOL3 234
FNOL3 235
FNOL3 236
FNOL3 237
FNOL3 238
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SUBROUTINE TERM2 TRACE

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```
      SUBROUTINE TERM2(X,Y,D,F)
      DIMENSION Y(1),D(1)
      COMMON/GAMM/WRAD,ZM,FLEN,CR,TL
      COMMON/BG/N,DRT,PRAD,PRTU,TOP
      5      F=TL-X
      IF(DRT.EQ.0.0)RETURN
      R1=(Y(1)**2+Y(2)**2)**.5
      R2=(Y(3)**2+Y(4)**2)**.5
      15     IF(ABS(R1-R2).GT.DRT)F=0.0
      10      RETURN
      END
```

	TERM2	2
	TERM2	3
	TERM2	4
	TERM2	5
	TERM2	6
	TERM2	7
	TERM2	8
	TERM2	9
	TERM2	10
	TERM2	11
	TERM2	12

SUBROUTINE	TERM3	TRACE	CDC 6400 FTN V3.0-P316 OPT=0	01/07/74	14.28.01.	PAGE	1
		SUBROUTINE TERM3(X,Y,D,F)			TERM3	2	
		COMMON/BG/N,DRT,PRAD,PRTU,TOP			TERM3	3	
		COMMON/GAMM/WRAD,ZM,FLEN,CR,TL,TSL			TERM3	4	
		DIMENSION Y(1),D(1)			TERM3	5	
5		DELT=0.0			TERM3	6	
		CALL PROFILE(R,DR,X,DELT,2)			TERM3	7	
		R2=R*R			TERM3	8	
		IF(Y(1)**2+Y(2)**2.GT.R2.AND.Y(3)**2+Y(4)**2.GT.R2)GO TO 2			TERM3	9	
		TL=X			TERM3	10	
10		WRITE(6,2000)			TERM3	11	
	2	XFIN=TS			TERM3	12	
		IF(TL.LT.XFIN)XFIN=TL			TERM3	13	
		F=XFIN-X			TERM3	14	
		RETURN			TERM3	15	
15	2000	FORMAT(51H0EXECUTION TERMINATED DUE TO NUMERICAL INSTABILITY.)			TERM3	16	
		END			TERM3	17	

SUBROUTINE DERIV TRACE CDC 6400 FTN V3.0-P316 OPT=0 01/07/74 14.28.01. PAGE 1

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        SUBROUTINE DERIV(X,Y,D)                               DERIV  2
        COMPLEXA,B,T,E,XF,G,XZ,XF1,XF2,XF3,DZOP,XJ,XJC,C1,C2,C3,C4,C5,  DERIV  3
        1C6,K11,K12,K13,K14,K21,K22,K23,K24,K15,K16,K25,K26,DTH,DZ0,TH2,  DERIV  4
        2KN1,KN2,WN,XK,XKK,XWN,CIR,XC,L1,L2,X1,X2,X0,XOP,P,H,C7,K19,K29,XNP  DERIV  5
5       INTEGER Z0,ZOP,Z1,Z2,ZJ                           DERIV  6
        COMMON/GARB/XLR,R,DR,XC,TAL,FVORT                  DERIV  7
        COMMON/BG/N,DRT,PRAD,PRTU,TOP                      DERIV  8
        COMMON/VEL/U,IE                                     DERIV  9
        DIMENSION Y(1),D(1),WN(15),DD(4,4),Q(4,1),XL(15),XC(15),XZ(2),  DERIV 10
10      1A(2,2),B(2),T(2,15),E(2,15),XF(2,15),G(2),P(2),H(2,2),XNP(15),  DERIV 11
        3XLR(15)                                         DERIV 12
        CALL PROFILE(R,DR,X,DELTA, 2)                      DERIV 13
        XK=CMPLX(0.0,1.0)                                 DERIV 14
        XKK=XK*R                                         DERIV 15
15      IC=(N-1)/2                                     DERIV 16
        DO 22 KQ=1,IC                                    DERIV 17
22      XL(KQ)=XLR(KQ)/R                           DERIV 18
        C...OBTAIN THE STAGNATION POINTS                 DERIV 19
        THETA=Y(N)*(1.-TOP)                            DERIV 20
20      S1=R*COS(THETA)                            DERIV 21
        S2=R*SIN(THETA)                            DERIV 22
        XZ(1)= CMPLX(S1,S2)                           DERIV 23
        S1=-S1                                         DERIV 24
        XZ(2)=CMPLX(S1,S2)                           DERIV 25
25      IF(IC.EQ.2.OR.IC.NE.2*(IC/2))GO TO 6        DERIV 26
        DTH=XZ(1)                                     DERIV 27
        XZ(1)=XZ(2)                                     DERIV 28
        XZ(2)=DTH                                     DERIV 29
        6 CONTINUE                                     DERIV 30
30      C.....TRANSFORMS VORTEX POSITIONS TO COMPLEX COORDINATES  DERIV 31
        DO 5 I=1,IC                                    DERIV 32
5       XC(I)=CMPLX(Y(2*I-1),Y(2*I))                DERIV 33
        C.....CALCULATE FUNCTIONAL VALUES               DERIV 34
        DO 4 L=1,2                                     DERIV 35
35      X1=XZ(L)                                     DERIV 36
        B(L)=(R*R/X1/X1+1.)/R+FVORT/X1                DERIV 37
        G(L)=-2.*R/X1/X1/X1-FVORT/X1/X1               DERIV 38
        P(L)=DR*(1./X1/X1-1./R/R)                     DERIV 39
        DO 4 K=1,IC                                    DERIV 40
40      XJ=XC(K)                                     DERIV 41
        XJC=CONJG(XJ)                                 DERIV 42
        XF1=1./(X1-XJ)-1./(X1-R*R/XJC)                DERIV 43
        XF2=1./(X1-R*R/XJC)**2-1./(X1-XJ)**2          DERIV 44
        XF3=-DR*2.*R/((X1-R*R/XJC)**2)/XJC          DERIV 45
45      XF(L,K)=(R/(X1-R*R/XJC)/XJC)**2            DERIV 46
        E(L,K)=1./(X1-XJ)**2                           DERIV 47
        IF(K.GT.2)GO TO 3                            DERIV 48
        A(L,K)=XF1                                     DERIV 49
        T(L,K)=XF2                                     DERIV 50
50      H(L,K)=XF3                                     DERIV 51
        GO TO 4                                     DERIV 52
        3 B(L)=B(L)+XL(K)*XF1                         DERIV 53
        G(L)=G(L)+XL(K)*XF2                         DERIV 54
        P(L)=P(L)+XL(K)*XF3                         DERIV 55
55      4 CONTINUE                                     DERIV 56

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SUBROUTINE	DERIV	TRACE	CDC 6400 FTN V3.0-P316 OPT=0	01/07/74	14.28.01.	PAGE
			C.....CALCULATE THE TWO UNKNOWN CIRCULATIONS			
	Z0=1		Z0=1	DERIV	57	
	ZOP=2		ZOP=2	DERIV	58	
	Z1=1		Z1=1	DERIV	59	
60	Z2=2		Z2=2	DERIV	60	
	A1=REAL(-B(Z0)*A(ZOP,Z2)+B(ZOP)*A(Z0,Z2))		A1=REAL(-B(Z0)*A(ZOP,Z2)+B(ZOP)*A(Z0,Z2))	DERIV	61	
	B1=REAL(-B(ZOP)*A(Z0,Z1)+B(Z0)*A(ZOP,Z1))		B1=REAL(-B(ZOP)*A(Z0,Z1)+B(Z0)*A(ZOP,Z1))	DERIV	62	
	C=REAL(A(Z0,Z1)*A(ZOP,Z2)-A(Z0,Z2)*A(ZOP,Z1))		C=REAL(A(Z0,Z1)*A(ZOP,Z2)-A(Z0,Z2)*A(ZOP,Z1))	DERIV	63	
	XL(1)=A1/C		XL(1)=A1/C	DERIV	64	
65	XL(2)=B1/C		XL(2)=B1/C	DERIV	65	
	XLR(1)=XL(1)*R		XLR(1)=XL(1)*R	DERIV	66	
	XLR(2)=XL(2)*R		XLR(2)=XL(2)*R	DERIV	67	
	IF(U.LT.0)RETURN		IF(U.LT.0)RETURN	DERIV	68	
70	C.....CALCULATE THE FLUID VELOCITY AT EACH VORTEX LOCATION		C.....CALCULATE THE FLUID VELOCITY AT EACH VORTEX LOCATION	DERIV	69	
	DO 20 K=1,IC		DO 20 K=1,IC	DERIV	70	
	X1=XC(K)		X1=XC(K)	DERIV	71	
	XWN=R*DR/X1/TAL-XK*(1.+R*R/X1/X1)-XK*FVORT/X1		XWN=R*DR/X1/TAL-XK*(1.+R*R/X1/X1)-XK*FVORT/X1	DERIV	72	
	DO 10 L=1,IC		DO 10 L=1,IC	DERIV	73	
	XJ=XC(L)		XJ=XC(L)	DERIV	74	
75	IF(L.NE.K)GO TO 8		IF(L.NE.K)GO TO 8	DERIV	75	
	XNP(L)=XKK*XL(L)/(X1-R*R/CONJG(XJ))		XNP(L)=XKK*XL(L)/(X1-R*R/CONJG(XJ))	DERIV	76	
	GO TO 10		GO TO 10	DERIV	77	
	8 XNP(L)=-XKK*XL(L)*(1./(X1-XJ)-1./(X1-R*R/CONJG(XJ)))		8 XNP(L)=-XKK*XL(L)*(1./(X1-XJ)-1./(X1-R*R/CONJG(XJ)))	DERIV	78	
80	10 CONTINUE		10 CONTINUE	DERIV	79	
	WN(K)=XWN+XNP(K)		WN(K)=XWN+XNP(K)	DERIV	80	
	DO 15 MBS=1,IC		DO 15 MBS=1,IC	DERIV	81	
	IF(MBS.EQ.K)GO TO 15		IF(MBS.EQ.K)GO TO 15	DERIV	82	
	WN(K)=WN(K)+XNP(MBS)		WN(K)=WN(K)+XNP(MBS)	DERIV	83	
	15 CONTINUE		15 CONTINUE	DERIV	84	
85	20 WN(K)=CONJG(WN(K))		20 WN(K)=CONJG(WN(K))	DERIV	85	
	C.....DETERMINE THE DERIVATIVE OF THETA		C.....DETERMINE THE DERIVATIVE OF THETA	DERIV	86	
	DELTA=ATAN(DR)		DELTA=ATAN(DR)	DERIV	87	
	THETAI=THETI(DELTA)		THETAI=THETI(DELTA)	DERIV	88	
	D(N)=(.7*TAL/R)*(THETAI-Y(N))		D(N)=(.7*TAL/R)*(THETAI-Y(N))	DERIV	89	
90	C.....CALCULATE DERIVATIVES FOR ALL FREE VORTICES		C.....CALCULATE DERIVATIVES FOR ALL FREE VORTICES	DERIV	90	
	IF(IC.LT.3)GO TO 40		IF(IC.LT.3)GO TO 40	DERIV	91	
	DO 30 IL=3,IC		DO 30 IL=3,IC	DERIV	92	
	RWN=REAL(WN(IL))		RWN=REAL(WN(IL))	DERIV	93	
	XIWN=AIMAG(WN(IL))		XIWN=AIMAG(WN(IL))	DERIV	94	
95	D(2*IL)=XIWN*TAL		D(2*IL)=XIWN*TAL	DERIV	95	
	30 D(2*IL-1)=RWN*TAL		30 D(2*IL-1)=RWN*TAL	DERIV	96	
	40 CONTINUE		40 CONTINUE	DERIV	97	
	C.....CALCULATE DD MATRIX FOR ATTACHED VORTICES		C.....CALCULATE DD MATRIX FOR ATTACHED VORTICES	DERIV	98	
100	X2=XC(2)		X2=XC(2)	DERIV	99	
	X1=XC(1)		X1=XC(1)	DERIV	100	
	X0=XZ(1)		X0=XZ(1)	DERIV	101	
	XOP=XZ(2)		XOP=XZ(2)	DERIV	102	
	C1=-(A(ZOP,Z2)*E(Z0,Z1)-A(Z0,Z2)*E(ZOP,Z1))/C		C1=-(A(ZOP,Z2)*E(Z0,Z1)-A(Z0,Z2)*E(ZOP,Z1))/C	DERIV	103	
	C2=-(A(ZOP,Z2)*XF(Z0,Z1)-A(Z0,Z2)*XF(ZOP,Z1))/C		C2=-(A(ZOP,Z2)*XF(Z0,Z1)-A(Z0,Z2)*XF(ZOP,Z1))/C	DERIV	104	
	C3=-(A(Z0,Z1)*E(ZOP,Z2)-A(ZOP,Z1)*E(Z0,Z2))/C		C3=-(A(Z0,Z1)*E(ZOP,Z2)-A(ZOP,Z1)*E(Z0,Z2))/C	DERIV	105	
	C4=-(A(Z0,Z1)*XF(ZOP,Z2)-A(ZOP,Z1)*XF(Z0,Z2))/C		C4=-(A(Z0,Z1)*XF(ZOP,Z2)-A(ZOP,Z1)*XF(Z0,Z2))/C	DERIV	106	
	K11=C1*(X1-X0)+1.		K11=C1*(X1-X0)+1.	DERIV	107	
	K12=C2*(X1-X0)		K12=C2*(X1-X0)	DERIV	108	
	K13=(C3-(B(Z0)*E(ZOP,Z2)-B(ZOP)*E(Z0,Z2))/A1)*(X1-X0)		K13=(C3-(B(Z0)*E(ZOP,Z2)-B(ZOP)*E(Z0,Z2))/A1)*(X1-X0)	DERIV	109	
110	K14=(C4-(B(Z0)*XF(ZOP,Z2)-B(ZOP)*XF(Z0,Z2))/A1)*(X1-X0)		K14=(C4-(B(Z0)*XF(ZOP,Z2)-B(ZOP)*XF(Z0,Z2))/A1)*(X1-X0)	DERIV	110	
				DERIV	111	

SUBROUTINE DERIV TRACE CDC 6400 FTN V3.0-P316 OPT=0 01/07/74 14.28.01.

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K21=(C1-(B(ZOP)*E(Z0,Z1)-B(Z0)*E(ZOP,Z1))/B1)*(X2-XOP) DERIV 112
K22=(C2-(B(ZOP)*XF(Z0,Z1)-B(Z0)*XF(ZOP,Z1))/B1)*(X2-XOP) DERIV 113
K23=C3*(X2-XOP)+1. DERIV 114
K24=C4*(X2-XOP) DERIV 115
115 DD(1,1)=REAL(K11+K12) DERIV 116
DD(1,2)=AIMAG(K12-K11) DERIV 117
DD(1,3)=REAL(K13+K14) DERIV 118
DD(1,4)=AIMAG(K14-K13) DERIV 119
DD(2,1)=AIMAG(K11+K12) DERIV 120
120 DD(2,2)=REAL(K11-K12) DERIV 121
DD(2,3)=AIMAG(K13+K14) DERIV 122
DD(2,4)=REAL(K13-K14) DERIV 123
DD(3,1)=REAL(K21+K22) DERIV 124
DD(3,2)=AIMAG(K22-K21) DERIV 125
125 DD(3,3)=REAL(K23+K24) DERIV 126
DD(3,4)=AIMAG(K24-K23) DERIV 127
DD(4,1)=AIMAG(K21+K22) DERIV 128
DD(4,2)=REAL(K21-K22) DERIV 129
130 DD(4,3)=AIMAG(K23+K24) DERIV 130
DD(4,4)=REAL(K23-K24) DERIV 131
C.....CALCULATE THE Q VECTORS DERIV 132
C5=(-A(ZOP,Z2)*T(Z0,Z1)+A(ZOP,Z1)*T(Z0,Z2))/C DERIV 133
C6=(-A(Z0,Z1)*T(ZOP,Z2)-A(Z0,Z2)*T(ZOP,Z1))/C DERIV 134
C7=(A(ZOP,Z1)*H(Z0,Z2)+A(Z0,Z2)*H(ZOP,Z1)-A(ZOP,Z2)*H(Z0,Z1) DERIV 135
1-A(Z0,Z1)*H(ZOP,Z2))/C DERIV 136
135 K15=(C5-(A(ZOP,Z2)*G(Z0)-B(ZOP)*T(Z0,Z2))/A1)*(X1-X0) DERIV 137
K16=(C6+(T(ZOP,Z2)*B(Z0)-A(Z0,Z2)*G(ZOP))/A1)*(X1-X0) DERIV 138
K19=(C7+(A(Z0,Z2)*P(ZOP)+B(ZOP)*H(Z0,Z2)-A(ZOP,Z2)*P(Z0)-B(Z0)*
1H(ZOP,Z2))/A1)*(X1-X0) DERIV 139
140 K25=(C5+(A(ZOP,Z1)*G(Z0)-T(Z0,Z1)*B(ZOP))/B1)*(X2-XOP) DERIV 140
K26=(C6+(A(Z0,Z1)*G(ZOP)-B(Z0)*T(ZOP,Z1))/B1)*(X2-XOP) DERIV 141
K29=(C7+(A(ZOP,Z1)*P(Z0)+B(Z0)*H(ZOP,Z1)-A(Z0,Z1)*P(ZOP)-B(ZOP)*
1H(Z0,Z1))/B1)*(X2-XOP) DERIV 142
DT2=D(N)*R*(1.-TOP) DERIV 143
145 DTH=CMPLX(DR,DT2) DERIV 144
DT2=0.0,THETA) DERIV 145
DZ0=CEXP(TH2)*DTH DERIV 146
DTH=CMPLX(-DR,DT2) DERIV 147
DZOP=-CEXP(-TH2)*DTH DERIV 148
150 IF(IC.EQ.2.0R.IC.NE.2*(IC/2))GO TO 55 DERIV 149
TH2=DZ0 DERIV 150
DZ0=DZOP DERIV 151
DZOP=TH2 DERIV 152
55 CONTINUE DERIV 152
155 L1=-K15*DZ0-K16*DZOP+TAL*WN(1)-DR*(X1-X0)/R-K19 DERIV 153
L2=-K25*DZ0-K26*DZOP+TAL*WN(2)-DR*(X2-XOP)/R-K29 DERIV 154
IF(IC.LT.3)GO TO 50 DERIV 155
DO 60 M=3,IC DERIV 156
DTH=CMPLX(D(2*M-1),D(2*M)) DERIV 157
160 TH2=CONJG(DTH) DERIV 158
ZJ=M DERIV 159
CIR=XL(M) DERIV 160
KN1=(X1-X0)*(-A(ZOP,Z2)*E(Z0,ZJ)+A(Z0,Z2)*E(ZOP,ZJ))/A1*CIR DERIV 161
KN2=(X1-X0)*(-A(ZOP,Z2)*XF(Z0,ZJ)+A(Z0,Z2)*XF(ZOP,ZJ))/A1*CIR DERIV 162
L1=L1-KN1*DTH-KN2*TH2 DERIV 163
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SUBROUTINE	DERIV	TRACE	CDC 6400 FTN V3.0-P316 OPT=0	01/07/74	14.28.01.	PAGE	4
			KN1=(-E(ZOP,ZJ)*A(ZO,Z1)+A(ZOP,Z1)*E(ZO,ZJ))*(X2-XOP)*CIR/B1	DERIV	167		
			KN2=(-XF(ZOP,ZJ)*A(ZO,Z1)+XF(ZO,ZJ)*A(ZOP,Z1))*(X2-XOP)*CIR/B1	DERIV	168		
	60	L2=L2-DTH*KN1-TH2*KN2		DERIV	169		
170	50	Q(1,1)=REAL(L1)		DERIV	170		
		Q(2,1)=AIMAG(L1)		DERIV	171		
		Q(3,1)=REAL(L2)		DERIV	172		
		Q(4,1)=AIMAG(L2)		DERIV	173		
		II=4		DERIV	174		
		IX=1		DERIV	175		
175		CALL MATRIX(DD,II,II,Q,IX,DETERM,1D)		DERIV	176		
		D(1)=Q(1,1)		DERIV	177		
		D(2)=Q(2,1)		DERIV	178		
		D(3)=Q(3,1)		DERIV	179		
		D(4)=Q(4,1)		DERIV	180		
180		RETURN		DERIV	181		
		END		DERIV	182		

SUBROUTINE OUT TRACE CDC 6400 FTN V3.0-P316 OPT=0 01/07/74 14.28.01. PAGE 1

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    SUBROUTINE OUT(X,Y,D,ERROR,N,L,H) OUT 2
    COMPLEX XC,F1 OUT 3
    DIMENSION XC(15),XL(15),D(1),ERROR(1),Y(1),XLR(15) OUT 4
    COMMON/STA/A,NGB OUT 5
    COMMON/XLIF/MC,M,CDMN,CDML,XS,CYM,CDCL1,CDCN1,CY1 OUT 6
    COMMON/GARB/XLR,R,DR,XC,TAL OUT 7
    COMMON/VEL/U,IE OUT 8
    COMMON/GAMM/R0,ZM,FLEN,CR,TL OUT 9
    COMMON/BG/IBLANK,DRT OUT 10
    CONVR=57.2957795 OUT 11
    IF(DRT.NE.0.0) GO TO 8 OUT 12
    Y(1)=(Y(1)-Y(3))/2. OUT 13
    Y(3)=-Y(1) OUT 14
    Y(2)=(Y(2)+Y(4))/2. OUT 15
    Y(4)=Y(2) OUT 16
    8 CONTINUE OUT 17
    U=-U OUT 18
    CALL DERIV(X,Y,D) OUT 19
    U=-U OUT 20
    20 C.....CALCULATE THE FORCE COEFFICIENTS OUT 21
    CDCL=2.*SIN(A)*COS(A)*R*R/R0/R0 OUT 22
    F1=CMPLX(0.0,0.0) OUT 23
    IC=(N-1)/2 OUT 24
    DO 22 KQ=1,IC OUT 25
    22 XL(KQ)=XLR(KQ)/R OUT 26
    DO 10 I=1,IC OUT 27
    10 F1=XL(I)*(XC(I)-R*R/CONJG(XC(I)))+F1 OUT 28
    F1=4.*SIN(A)*COS(A)*F1*R/R0/R0 OUT 29
    CDCN=REAL(F1) OUT 30
    CY=AIMAG(F1) OUT 31
    CD=0.0 OUT 32
    IF(X.NE.XS) CD=(CDCN-CDCN1)/(X-XS) OUT 33
    CD=CD*3.1416*R/2./SIN(A)/SIN(A) OUT 34
    C.....CALCULATE MOMENT COEFFICIENTS OUT 35
    CDML=(X+XS)*(CDCL-CDCL1)/4./R0 +CDML OUT 36
    CDMN=(XS+X)*(CDCN-CDCN1)/4./R0 +CDMN OUT 37
    CYM=(X+XS)*(CY-CY1)/4./R0 +CYM OUT 38
    C.....REINITIALIZE FORCE COEFFICIENTS OUT 39
    CDCL1=CDCL OUT 40
    CDCN1=CDCN OUT 41
    CY1=CY OUT 42
    IF(X.EQ.XS) CIT=0 OUT 43
    XS=X OUT 44
    M=M-1 OUT 45
    45 IF(X.EQ.TL) GO TO 16 OUT 46
    CIT=CIT+1 OUT 47
    IF(CIT.LT.NGB) GO TO 16 OUT 48
    IF(MC.EQ.0) GO TO 15 OUT 49
    IF(M.EQ.0) GO TO 16 OUT 50
    50 RETURN OUT 51
    15 IF(M.EQ.0) GO TO 17 OUT 52
    IF(X-Y(N+1).LT.Y(N+2)) RETURN OUT 53
    Y(N+1)=X OUT 54
    GO TO 17 OUT 55
    55 M=MC OUT 56
  
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SUBROUTINE	OUT	TRACE	CDC 6400 FTN V3.0-P316 OPT=0	01/07/74	14.28.01.	PAGE	2
	17	DDEG=CONVR*Y(N)		OUT	57		
		HPP=0.0		OUT	58		
		IF (ABS(H).LT.1.)HPP=H		OUT	59		
		WRITE(6,2000)X,R,DDEG,HPP		OUT	60		
60		WRITE(6,2004)		OUT	61		
	DO 20	KT=1,IC		OUT	62		
		K=1+IC-KT		OUT	63		
		IF (IC.EQ.2)K=KT		OUT	64		
		XR2=(Y(2*K-1)**2+Y(2*K)**2)**.5		OUT	65		
65		ANG2=CONVR*ATAN(Y(2*K-1)/Y(2*K))		OUT	66		
		IF (Y(2*K).LT.0.0)ANG2=Y(2*K-1)/ABS(Y(2*K-1))*(ABS(ANG2)+90.0)		OUT	67		
	20	WRITE(6,2001)KT,XL(K),XR2,ANG2,Y(2*K-1),Y(2*K)		OUT	68		
		WRITE(6,2002)CDCL,CDML,CDCN,CDMN,CD,CY,CYM		OUT	69		
		IF (IE.EQ.0)RETURN		OUT	70		
70		WRITE(6,2003)(ERROR(II),II=1,N)		OUT	71		
		RETURN		OUT	72		
	2000	FORMAT(5H0X = ,F12.7,4X,4HR = ,F12.7,4X,9HTHETA0 = ,F10.5,4X,12HST		OUT	73		
		1EP SIZE = ,E14.8)		OUT	74		
	2001	FORMAT(1H ,20X,I3,6X,5(F14.7,3X))		OUT	75		
75	2002	FORMAT(1H ,6X,2HCL,12X,4H CML,11X,2HCN,12X,3HCMN,11X,2HCD,12X,2HCY		OUT	76		
		1,12X,3HCYM,/7F14.7)		OUT	77		
	2003	FORMAT(8H ERRORS ,8E14.6)		OUT	78		
	2004	FORMAT(1H ,15X,14HVORTEX NUMBER ,7X,3HLAM,13X,6HRADIUS,11X,5HANGLE		OUT	79		
		1,13X,1HY,16X,1HZ)		OUT	80		
80		END		OUT	81		

SUBROUTINE	MATRIX	TRACE	CDC 6400 FTN V3.0-P316 OPT=0	01/07/74	14.28.01.	PAGE
			SUBROUTINE MATRIX(A,N1,I1,B,M1,DETERM, ID)		MATRIX	2
	C		MATRIX INVERSION WITH ACCOMPANYING SOLUTION OF LINEAR EQUATIONS		MATRIX	3
	C		TEST FOR LOST OF DIGITS DUE TO SUBTRACTION		MATRIX	4
5	C		DIMENSION A(I1,I1),B(I1,M1),INDEX(50,3)		MATRIX	5
			EQUIVALENCE (IROW,JROW), (ICOLUMN,JCOLUMN), (AMAX, T, SWAP)		MATRIX	6
			DATA EPS4/1.E-7/		MATRIX	7
10	C		ID=1		MATRIX	8
			M=M1		MATRIX	9
			N=N1		MATRIX	10
			10 DETERM=1.0		MATRIX	11
15			15 DO 20 J=1,N		MATRIX	12
			20 INDEX(J,3) = 0		MATRIX	13
			30 DO 550 I=1,N		MATRIX	14
	C		SEARCH FOR PIVOT ELEMENT		MATRIX	15
20	C		40 AMAX=0.0		MATRIX	16
			45 DO 105 J=1,N		MATRIX	17
			IF (INDEX(J,3) .EQ. 1) GO TO 105		MATRIX	18
			60 DO 100 K=1,N		MATRIX	19
			IF (INDEX(K,3)-1) 80, 100, 100		MATRIX	20
25			80 IF (AMAX .GE. ABS(A(J,K))) GO TO 100		MATRIX	21
			85 IROW=J		MATRIX	22
			90 ICOLUMN=K		MATRIX	23
			AMAX = ABS(A(J,K))		MATRIX	24
			100 CONTINUE		MATRIX	25
30			105 CONTINUE		MATRIX	26
			INDEX(ICOLUMN,3) = INDEX(ICOLUMN,3) +1		MATRIX	27
			260 INDEX(I,1)=IROW		MATRIX	28
			270 INDEX(I,2)=ICOLUMN		MATRIX	29
			IF (BMAX.EQ.AMAX) ID=2		MATRIX	30
35	C		INTERCHANGE ROWS TO PUT PIVOT ELEMENT ON DIAGONAL		MATRIX	31
	C		130 IF (IROW .EQ. ICOLUMN) GO TO 310		MATRIX	32
	C		140 DETERM=-DETERM		MATRIX	33
40			150 DO 200 L=1,N		MATRIX	34
			160 SWAP=A(IROW,L)		MATRIX	35
			170 A(IROW,L)=A(ICOLUMN,L)		MATRIX	36
			200 A(ICOLUMN,L)=SWAP		MATRIX	37
			IF(M) 310, 310, 210		MATRIX	38
45			210 DO 250 L=1, M		MATRIX	39
			220 SWAP=B(IROW,L)		MATRIX	40
			230 B(IROW,L)=B(ICOLUMN,L)		MATRIX	41
			250 B(ICOLUMN,L)=SWAP		MATRIX	42
	C		DIVIDE PIVOT ROW BY PIVOT ELEMENT		MATRIX	43
50	C		310 PIVOT =A(ICOLUMN,ICOLUMN)		MATRIX	44
			DETERM=DETERM*PIVOT		MATRIX	45
			330 A(ICOLUMN,ICOLUMN)=1.0		MATRIX	46
55			340 DO 352 L=1,N		MATRIX	47

SUBROUTINE	MATRIX	TRACE	CDC 6400 FTN V3.0-P316 OPT=0	01/07/74	14.28.01.	PAGE
						2
		IF (PIVOT .NE. 0.) GO TO 350	MATRIX	57		
		A(ICOLUMN,L) = 0.	MATRIX	58		
		GO TO 352	MATRIX	59		
60		350 A(ICOLUMN,L)=A(ICOLUMN,L)/PIVOT	MATRIX	60		
		352 CONTINUE	MATRIX	61		
		355 IF(M) 380, 380, 360	MATRIX	62		
		360 DO 375 L=1,M	MATRIX	63		
		IF (PIVOT .NE. 0.) GO TO 370	MATRIX	64		
		B(ICOLUMN,L) = 0.	MATRIX	65		
65		GO TO 375	MATRIX	66		
		370 B(ICOLUMN,L)=B(ICOLUMN,L)/PIVOT	MATRIX	67		
		375 CONTINUE	MATRIX	68		
	C		MATRIX	69		
	C	REDUCE NON-PIVOT ROWS	MATRIX	70		
70	C		MATRIX	71		
		BMAX=0.0	MATRIX	72		
		380 DO 550 L1=1,N	MATRIX	73		
		390 IF (L1 .EQ. ICOLUMN) GO TO 550	MATRIX	74		
		400 T=A(L1,ICOLUMN)	MATRIX	75		
75		420 A(L1,ICOLUMN)=0.0	MATRIX	76		
		430 DO 450 L=1,N	MATRIX	77		
		SUB =A(ICOLUMN,L)*T	MATRIX	78		
		A(L1,L)=A(L1,L)-SUB	MATRIX	79		
	C		MATRIX	80		
80	C		MATRIX	81		
		IF(INDEX(L1,3).EQ.1) GO TO 450	MATRIX	82		
		IF(ABS(A(L1,L))-EPS4* ABS(SUB)) 449,449,450	MATRIX	83		
		449 BMAX=AMAX1(BMAX,ABS(A(L1,L)))	MATRIX	84		
		450 CONTINUE	MATRIX	85		
85		455 IF(M) 550, 550, 460	MATRIX	86		
		460 DO 500 L=1,M	MATRIX	87		
		500 B(L1,L)=B(L1,L)-B(ICOLUMN,L)*T	MATRIX	88		
		550 CONTINUE	MATRIX	89		
	C		MATRIX	90		
90	C	INTERCHANGE COLUMNS	MATRIX	91		
	C		MATRIX	92		
		600 DO 710 I=1,N	MATRIX	93		
		610 L=N+1-I	MATRIX	94		
		620 IF (INDEX(L,1) .EQ. INDEX(L,2)) GO TO 710	MATRIX	95		
95		630 JROW=INDEX(L,1)	MATRIX	96		
		640 JCOLUMN=INDEX(L,2)	MATRIX	97		
		650 DO 705 K=1,N	MATRIX	98		
		660 SWAP=A(K,JROW)	MATRIX	99		
		670 A(K,JROW)=A(K,JCOLUMN)	MATRIX	100		
100		700 A(K,JCOLUMN)=SWAP	MATRIX	101		
		705 CONTINUE	MATRIX	102		
		710 CONTINUE	MATRIX	103		
		DO 730 K = 1,N	MATRIX	104		
		IF (INDEX(K,3) .NE. 1) GO TO 715	MATRIX	105		
105		730 CONTINUE	MATRIX	106		
		RETURN	MATRIX	107		
		715 ID=2	MATRIX	108		
		RETURN	MATRIX	109		
		END	MATRIX	110		

ANGLE OF ATTACK 45.00 DEGREES

OGIVE DIMENSIONS

MAXIMUM RADIUS	OGIVE LENGTH	Z MATCHING POINT	BODY LENGTH
.50000	3.100	3.10000	11.00000

*****FREE PARAMETERS*****

INITIAL VORTEX STRENGTH	.04000	
INITIAL RADIAL PERTURBATION	.00600	
RADIAL ASYMMETRY AT WHICH FIRST VORTEX IS SHED	.10000	
PER CENT DECREASE IN CIRCULATION OF SHED VORTICES	-2.50000	
LOCATION FACTOR FOR NEWLY INTRODUCED VORTICES	.07500	
STROUHAL NUMBER	.16660	
CORRESPONDING SHEDDING DISTANCE	3.00120	
SEPARATION ANGLE CONSTANTS	.73300	
-0.00000	POINT OF TRANSITION	1.00000

INTEGRATION OPTIONS

INITIAL STEP SIZE	.0002500	PRINT FREQUENCY	-0	ERROR BOUND	6	PRINT INCREMENT	1.0000000
				ERROR BOUND AFTER VORTEX SHEDDING	6		

X =	.1000000	R =	.0325289	THETAO =	41.99781	STEP SIZE =	.25000000E-03
VORTEX NUMBER	LAM	RADIUS	ANGLE	Y	Z		
1	.0355897	.0331857	46.6493614	.0241315	.0227807		
2	-.0578202	.0327899	-46.6493614	-.0238437	.0225090		
CL	CML	CN	CMN	CD	CY	CYM	
.0042325	.0002809	.0000144	.0000014	0.0000000	.0000029	.0000003	
ERRORS	0.	0.	0.	0.	0.		

X =	1.1030758	R =	.2956664	THETAO =	41.99781	STEP SIZE =	.21046273E-01
VORTEX NUMBER	LAM	RADIUS	ANGLE	Y	Z		
1	.2438641	.3327822	39.0463583	.2096358	.2584508		
2	-.2082865	.3244723	-40.5432500	-.2109141	.2465715		
CL	CML	CN	CMN	CD	CY	CYM	
.3496745	.2415966	.0430999	.0372341	.1302038	.0107881	.0094284	
ERRORS	.607359E-09	.218640E-09	.236340E-08	-.652509E-09	0.		

X =	2.1196964	R =	.4511470	THETAO =	41.99781	STEP SIZE =	.35862999E-01
VORTEX NUMBER	LAM	RADIUS	ANGLE	Y	Z		
1	.4714942	.6214725	30.9960132	.3200449	.5327282		
2	-.4356448	.5754213	-32.8963095	-.3125230	.4831553		
CL	CML	CN	CMN	CD	CY	CYM	
.8141343	.9785404	.4469488	.7358462	.9504091	.1361153	.2284237	
ERRORS	-.665432E-09	-.820530E-09	.471271E-10	-.646229E-09	0.		

X =	2.8623315	R =	.4971352	THETAO =	41.99781	STEP SIZE =	.25000000E-03
VORTEX NUMBER	LAM	RADIUS	ANGLE	Y	Z		
1	.6638951	.8170361	26.8023124	.3684126	.7292599		
2	-.6455558	.7118971	-27.9646625	-.3338277	.6287738		
CL	CML	CN	CMN	CD	CY	CYM	
.9885735	1.4009501	1.0517271	2.2520225	0.0000000	.3855470	.8595407	
ERRORS	-.105708E-09	.329976E-08	-.291649E-08	-.409414E-08	0.		

X =	2.8623315	R =	.4971352	THETAO =	41.99781	STEP SIZE =	.10000000E-02
VORTEX NUMBER	LAM	RADIUS	ANGLE	Y	Z		
1	.6472977	.8170361	26.8023124	.3684126	.7292599		
2	-.6421862	.7118971	-27.9646625	-.3338277	.6287738		
3	.0022805	.5227481	44.9565092	.3693581	.3699192		
CL	CML	CN	CMN	CD	CY	CYM	
.9885735	1.4009501	1.0344401	2.2025413	0.0000000	.3598689	.7860415	
ERRORS	0.	0.	0.	0.	0.	0.	

X = 3.1358214 R = .5000000 THETA0 = 41.99781 STEP SIZE = .17107098E-01
 VORTEX NUMBER LAM RADIUS ANGLE Y Z
 1 .6435889 .9293537 23.6852872 .3733331 .8510703
 2 -.7280791 .7461542 -26.2962073 -.3305552 .6689390
 3 .0139121 .5284685 46.1529236 .3811269 .3660892
 CL CML CN CMN CD CY CYM
 1.0000000 1.4345633 1.2155258 2.7449078 .9638301 .4855503 1.1628711
 ERRORS .166824E-07 .117700E-07 -.490723E-08 .449925E-08 .344445E-08 .805208E-08 0.
 X = 4.1446835 R = .5000000 THETA0 = 41.99781 STEP SIZE = .35007074E-01
 VORTEX NUMBER LAM RADIUS ANGLE Y Z
 1 .6435889 1.4027342 17.7113434 .4267421 1.3362464
 2 -.1.0237597 .8997440 -20.3545102 -.3129559 .8435625
 3 .1757541 .5636118 39.5839052 .3591377 .4343713
 CL CML CN CMN CD CY CYM
 1.0000000 1.4345633 1.8985781 5.2548438 1.3165386 .6803161 1.8155487
 ERRORS .371962E-08 .239884E-11 .809811E-10 -.652020E-09 -.393803E-09 -.422120E-09 0.
 X = 5.1549628 R = .5000000 THETA0 = 41.99781 STEP SIZE = .29757162E-01
 VORTEX NUMBER LAM RADIUS ANGLE Y Z
 1 .6435889 1.9788370 16.2281326 .5530109 1.8999934
 2 -.1.2720087 1.2138898 -16.4435008 -.3436154 1.1642409
 3 .3921419 .5724520 36.5638004 .3410197 .4597900
 CL CML CN CMN CD CY CYM
 1.0000000 1.4345633 2.9112933 9.9895735 1.8015862 -.1686913 -2.2403960
 ERRORS -.311141E-07 -.123495E-07 .649968E-10 -.286473E-09 -.973207E-12 -.383286E-10 0.
 X = 5.8635320 R = .5000000 THETA0 = 41.99781 STEP SIZE = 0.
 VORTEX NUMBER LAM RADIUS ANGLE Y Z
 1 .6435889 2.4158047 15.6478473 .6516006 2.3262693
 2 -.1.4302811 1.5175554 -15.0338253 -.3936376 1.4656138
 3 .5239926 .5833192 35.0541973 .3350300 .4775104
 CL CML CN CMN CD CY CYM
 1.0000000 1.4345633 3.7994491 14.8912346 0.0000000 -.1.4771218 -.94785667
 ERRORS -.165932E-06 -.758859E-07 .640592E-09 .155940E-08 .114145E-08 .486370E-09 0.
 X = 5.8635320 R = .5000000 THETA0 = 41.99781 STEP SIZE = .10000000E-02
 VORTEX NUMBER LAM RADIUS ANGLE Y Z
 1 .6435889 2.4158047 15.6478473 .6516006 2.3262693
 2 -.1.3945240 1.5175554 -15.0338253 -.3936376 1.4656138
 3 .5145310 .5833192 35.0541973 .3350300 .4775104
 4 -.0028980 .5263422 -45.0209574 -.3723162 .3720440
 CL CML CN CMN CD CY CYM
 1.0000000 1.4345633 3.7463173 14.5796946 0.0000000 -.1.2954685 -.8.4134367
 ERRORS 0. 0. 0. 0. 0. 0. 0.
 ERRORS 0.
 X = 6.1615320 R = .5000000 THETA0 = 41.99781 STEP SIZE = .16000000E-01
 VORTEX NUMBER LAM RADIUS ANGLE Y Z
 1 .6435889 2.6042565 15.3967217 .6914326 2.5107913
 2 -.1.3945240 1.7091872 -13.9904236 -.4132126 1.6584861
 3 .5714561 .5970385 33.7584075 .3317697 .4963707
 4 -.0076643 .5335495 -46.3339277 -.3859569 .3683915
 CL CML CN CMN CD CY CYM
 1.0000000 1.4345633 4.0499953 16.4051674 1.5663684 -.1.8966131 -.12.0277866
 ERRORS .353227E-08 .589238E-09 -.121830E-09 -.504559E-10 .452755E-11 .363392E-11 -.873714E-12 .946037E-12
 ERRORS 0.
 X = 7.1721751 R = .5000000 THETA0 = 41.99781 STEP SIZE = .32094397E-01
 VORTEX NUMBER LAM RADIUS ANGLE Y Z
 1 .6435889 3.2348758 14.4416539 .8067585 3.1326605
 2 -.1.3945240 2.4021354 -11.1526320 -.4646290 2.3567720
 3 .8646256 .7177918 25.8484516 .3129517 .6459770
 4 -.0825237 .5600676 -41.4797306 -.3709636 .4195971
 CL CML CN CMN CD CY CYM
 1.0000000 1.4345633 5.0887432 23.3413337 1.7434051 -.3.5828642 -.23.1921245

ERRORS .132244E-08 -.211717E-08 .292732E-09 -.290729E-10 -.248214E-10 .382313E-10 -.172845E-10 .215497E-10
 ERRORS 0.

X = 8.1991958 R = .5000000 THETA0 = 41.99781 STEP SIZE = .32094397E-01
 VORTEX NUMBER LAM RADIUS ANGLE Y Z
 1 .6435889 3.8595730 13.1682765 .8792562 3.7580862
 2 -1.3945240 3.1316880 -9.5957769 -.5220401 3.0878705
 3 1.1935511 1.0327161 18.2391696 .3232239 .9808307
 4 -.2689768 .5561387 -38.7058756 -.3477662 .4339919
 CL CML CN CMN CD CY CYM
 1.0000000 1.4345633 6.3164263 32.7899077 1.9881864 -3.7775881 -24.5257429
 ERRORS .731932E-09 -.986816E-10 -.325549E-09 -.216827E-09 .257688E-10 .622232E-11 .236255E-11 -.817124E-12
 ERRORS 0.

X = 8.8647325 R = .5000000 THETA0 = 41.99781 STEP SIZE = 0.
 VORTEX NUMBER LAM RADIUS ANGLE Y Z
 1 .6435889 4.2580371 12.1574910 .8967408 4.1625396
 2 -1.3945240 3.6060791 -9.1234539 -.5717880 3.5604585
 3 1.3749567 1.3339893 15.6128605 .3590245 1.2847680
 4 -.3515362 .5549913 -38.6311965 -.3464839 .4335485
 CL CML CN CMN CD CY CYM
 1.0000000 1.4345633 7.1938063 40.2793319 0.0000000 -2.9521614 -17.4474940
 ERRORS .888596E-06 -.352818E-06 -.744009E-09 -.187193E-08 .596501E-11 -.720727E-10 -.225967E-10 -.219927E-10
 ERRORS 0.

X = 8.8647325 R = .5000000 THETA0 = 41.99781 STEP SIZE = .10000000E-02
 VORTEX NUMBER LAM RADIUS ANGLE Y Z
 1 .6435889 4.2580371 12.1574910 .8967408 4.1625396
 2 -1.3945240 3.6060791 -9.1234539 -.5717880 3.5604585
 3 1.3405828 1.3339893 15.6128605 .3590245 1.2847680
 4 -.3449832 .5549913 -38.6311965 -.3464839 .4335485
 5 .0030946 .5254779 44.9261199 .3710896 .3720478
 CL CML CN CMN CD CY CYM
 1.0000000 1.4345633 7.1501010 39.8918962 0.0000000 -3.1014182 -18.7706156
 ERRORS 0. 0. 0. 0. 0. 0. 0.
 ERRORS 0. 0. 0.

X = 9.2130058 R = .5000000 THETA0 = 41.99781 STEP SIZE = .14707534E-01
 VORTEX NUMBER LAM RADIUS ANGLE Y Z
 1 .6435889 4.4654066 11.5864241 .8968582 4.3744144
 2 -1.3945240 3.8539932 -8.9834568 -.6017983 3.8067181
 3 1.3405828 1.5724424 13.6477708 .3710216 1.5280439
 4 -.3858415 .5614540 -37.9124832 -.3449894 .4429592
 5 .0106460 .5376864 46.1574387 .3878041 .3724441
 CL CML CN CMN CD CY CYM
 1.0000000 1.4345633 7.4810645 42.8818751 1.3880138 -2.5306527 -13.6101726
 ERRORS -.165317E-08 .131929E-08 .843059E-11 -.463477E-11 -.817403E-11 .558131E-11 .408816E-12 -.119777E-12
 ERRORS .639488E-13 -.178651E-12 0.

X = 10.2611996 R = .5000000 THETA0 = 41.99781 STEP SIZE = .63391534E-01
 VORTEX NUMBER LAM RADIUS ANGLE Y Z
 1 .6435889 5.0895694 9.7266230 .8598694 5.0164072
 2 -1.3945240 4.5943797 -8.8936108 -.7102921 4.5391420
 3 1.3405828 2.3492428 9.1053566 .3717686 2.3196400
 4 -.6270742 .6291921 -31.5159190 -.3289010 .5363831
 5 .0965087 .6233788 40.7982071 .4073138 .4719074
 CL CML CN CMN CD CY CYM
 1.0000000 1.4345633 8.3707261 51.5497577 1.4464884 -.7860946 3.3568254
 ERRORS -.141296E-07 -.659408E-08 -.998941E-08 .161486E-07 -.408549E-09 .977546E-09 .201482E-09 -.208869E-09
 ERRORS .377907E-10 -.231874E-09 0.

X = 11.0000000 R = .5000000 THETA0 = 41.99851 STEP SIZE = 0.
 VORTEX NUMBER LAM RADIUS ANGLE Y Z
 1 .6435889 5.5301893 8.3632984 .8043622 5.4713797
 2 -1.3945240 5.1088224 -8.9891053 -.7982364 5.0460762
 3 1.3405828 2.9151405 7.0659801 .3585980 2.8930005

	4	- .8784763	.7550590	-24.7394574	- .3159866	.6857598
	5	.2409266	.6908357	36.0218935	.4062766	.5587426
CL	CML	CN	CMN	CD	CY	CYM
1.0000000	1.4345633	9.1400144	59.7364535	0.0000000	.0524773	12.2374875
ERRORS	.644419E-07	.419243E-06	.855520E-07	.533789E-06	.806359E-11	.303639E-11
ERRORS	.160275E-11	.120698E-11	.892321E-06			.688279E-12
						.253901E-11

*****FINAL STATISTICS*****
 NORMAL FORCE COEFFICIENT=CL+CN= 10.14001
 PITCHING MOMENT COEFFICIENT=CML+CMN= 61.17102
 CENTER OF NORMAL FORCE= 6.03264
 YAW FORCE COEFFICIENT= .05248
 YAW MOMENT COEFFICIENT= 12.23749
 CENTER OF YAW FORCE= 233.19566

APPENDIX C

EXPERIMENTAL INVESTIGATION

The experiments referred to in this report were carried out in the NOL Supersonic Tunnel No. 1. This facility is an open jet, blowdown tunnel, 16 inches square, and is capable of subsonic as well as supersonic speeds.

The test model had an overall fineness of 12, a nose fineness of 4 and a diameter of 1.1 inches. It was tested at the Mach numbers of 0.5, 0.7, 0.9 and 1.1, which corresponded to a Reynolds number variation of $2(10^5)$ to $5(10^5)$ based on the body diameter. The model was swept from 0 to 70 degrees angle of attack. Four runs were made at each Mach number with the nose roll orientation being varied 90 degrees from run to run. Select runs were made using a boundary-layer trip. This consisted of a grit strip 1/8-inch wide along the leeward meridian of the body. Since the wind tunnel is an open jet type, it was possible to use a subsonic nozzle throughout the tests without encountering problems with reflected shock waves.

Force measurements were made using a four-component strain-gage balance. This type of information was augmented by schlieren photographs which were taken at increments of five degrees in angle of attack. Oil-flow studies were made at each Mach number at angles of attack of 20, 35, 50 and 70 degrees.

Results from the force tests indicate that normal force increases slightly with increasing Mach number, while the yaw force decreased strongly under these conditions. The measured forces and moments agree well with the data taken by Pick (Ref. (20)) using the same model. Pick's results, though, only go up to an angle of attack of 45 degrees. The Strouhal number measured from the schlieren photographs appears to be lower than that measured by Pick, although there is considerable scatter in the data obtained in the present study. The boundary-layer trip significantly lowered the normal force and greatly reduced the yaw force at the lower Mach numbers.

In order for force measurements to be meaningful at low crossflow Mach numbers, it is necessary to know whether the boundary layer is laminar or turbulent. It is well known that for $M_c \lesssim 0.5$, the transition of the boundary is accompanied by a large decrease in normal force. This results from an increase in the crossflow separation angle. At larger crossflow Mach numbers the transition of the boundary layer no longer has a large effect on normal force (Ref. (3)).

In the tests conducted, the crossflow Reynolds number ($Re \sin\alpha$) were of the order of 10^5 . Since transition is expected near this value, it is difficult to predict the state of the boundary layer. However, measurements of the crossflow separation angle taken from the oil-flow studies indicate that the boundary layer was laminar. This conclusion is supported by the reduction of force observed when the boundary-layer trip was used.

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